

NATIONAL LIBRARY OF CANADA  
BIBLIOTHÈQUE NATIONALE DU CANADA  
THESIS SERVICES / SERVICE DES THÈSES CANADIENNES

NUMERICAL SOLUTIONS OF A  
NON-LINEAR CLASSICAL FIELD THEORY  
OF LEPTONS AND QUARKS

by

Philip Howard Sharman,  
B.A.A., Ryerson Polytechnical Institute, 1978,  
B.Sc., University of Calgary, 1986.

A thesis submitted in partial fulfillment  
of the requirements for the degree of  
Master of Science  
in the Department  
of Physics and Astronomy.

ACCEPTED  
FACULTY OF GRADUATE STUDIES

DATE

1990-03-05

DEAN

We accept this thesis as conforming  
to the required standard.

Dr. F. I. Cooperstock, Department of Physics and Astronomy, University of Victoria.

Dr. C. E. Picciotto, Department of Physics and Astronomy, University of Victoria.

Dr. G.G. Miller, Department of Mathematics, University of Victoria.

Dr. J. Carimati, Department of Applied Mathematics, University of Waterloo.

© Philip Howard Sharman, 1989,  
University of Victoria.

All rights reserved. This thesis may not be reproduced in whole or in part, by  
mimeograph or other means, without the permission of the author.



Supervisor: Dr. F. I. Cooperstock.

ABSTRACT

This paper studies a non-linear classical field theory of elementary

The author has granted an irrevocable non-exclusive licence allowing the National Library of Canada to reproduce, loan, distribute or sell copies of his/her thesis by any means and in any form or format, making this thesis available to interested persons.

L'auteur a accordé une licence irrévocable et non exclusive permettant à la Bibliothèque nationale du Canada de reproduire, prêter, distribuer ou vendre des copies de sa thèse de quelque manière et sous quelque forme que ce soit pour mettre des exemplaires de cette thèse à la disposition des personnes intéressées.

The author retains ownership of the copyright in his/her thesis. Neither the thesis nor substantial extracts from it may be printed or otherwise reproduced without his/her permission.

L'auteur conserve la propriété du droit d'auteur qui protège sa thèse. Ni la thèse ni des extraits substantiels de celle-ci ne doivent être imprimés ou autrement reproduits sans son autorisation.

ISBN 0-315-62350-0

Examiners:

*F. Cooperstock*

Dr. F. I. Cooperstock, Department of Physics and Astronomy, University of Victoria.

*Chris Picciotto*

Dr. G. E. Picciotto, Department of Physics and Astronomy, University of Victoria.

*Gary G. Müller*

Dr. G. G. Müller, Department of Mathematics, University of Victoria.

*John Carminati*

Dr. J. Carminati, Department of Applied Mathematics, University of Waterloo.

TABLE OF CONTENTS

ABSTRACT ..... Supervisor: Dr. F. I. Cooperstock. .... 11

TABLE OF CONTENTS ..... 11

LIST OF TABLES ..... 12

ABSTRACT

This paper studies a non-linear, classical field theory of elementary particles. This theory is gauge-invariant and singularity-free and proposes that these particles are made up purely of three fundamental fields whose non-linear interactions produce a particle-like region of localized energy. Numerical solutions of the resulting equations are presented which show that the model can be used to describe not only the lepton family, but also the two known quark families.

C) Rosen ..... 9

D) Rosen and Rosenstock ..... 15

E) Cooperstock and Rosen ..... 27

CHAPTER 3 NEW WORK: FITTING LEPTONS AND QUARKS ..... 39

CHAPTER 4 FUTURE POSSIBILITIES ..... 60

Examiners:

[Redacted Name] ..... 62

Dr. F. I. Cooperstock, Department of Physics and Astronomy, University of Victoria.

[Redacted Name] ..... 65

Dr. G. E. Picciotto, Department of Physics and Astronomy, University of Victoria.

[Redacted Name] ..... 67

Dr. G. G. Miller, Department of Mathematics, University of Victoria.

[Redacted Name] ..... 68

Dr. J. Carminati, Department of Applied Mathematics, University of Waterloo.

	<u>TABLE OF CONTENTS</u>	102
	2) ESC	116
	3) IMPROVE	121
ABSTRACT		ii
TABLE OF CONTENTS		iii
LIST OF TABLES		v
LIST OF FIGURES		vi
ACKNOWLEDGEMENTS		vii
CHAPTER 1	INTRODUCTION	1
CHAPTER 2	PREVIOUS WORK	7
	A) Standard Electromagnetism	8
	B) Mie, Born and Infeld	9
	C) Rosen	9
	D) Rosen and Rosenstock	15
	E) Cooperstock and Rosen	27
CHAPTER 3	NEW WORK: FITTING LEPTONS AND QUARKS	39
CHAPTER 4	FUTURE POSSIBILITIES	60
BIBLIOGRAPHY		62
APPENDIX A	CONVENTIONS USED	64
APPENDIX B	THE DIMENSIONS OF VARIOUS SYMBOLS	65
APPENDIX C	GEOMETRIC UNITS	66
APPENDIX D	COMPUTER PROGRAM NOTES	67
APPENDIX E	COMPUTER PROGRAM LISTINGS	
	1) GETBCD	81

2) ESCALINT3 .....	102
3) IMPROVE .....	116
4) TRANSFER .....	121
5) PARIMP .....	123
6) REVISE_FS .....	135
7) PAR_ALPHA .....	137
8) TWEAK .....	144

TABLE 1:

The two quark equation

TABLE 2:

The EMT group as found by

TABLE 3:

Different estimates for the masses of quarks and leptons ( $\text{MeV}/c^2$ ) from Griffiths [8], Veltman [13], and *Physics Letters* [14].

TABLE 4:

The solution for the EMT group used as a basis for the two quark groups. 43

TABLE 5:

Solutions for the UCT and DSB groups which achieve the small mass ratios of the quarks' effective masses in baryons. 49

TABLE 6:

Solutions for the UCT and DSB groups which achieve the large mass ratios of the quarks' bare masses. 50

## LIST OF TABLES

TABLE 1:	The $y(0)$ values for the first fifty eigensolutions for the Rosenstock equation.	20
TABLE 2:	The solution for the EMT group as found by Cooperstock and Rosen.	37
TABLE 3:	Different values of estimates for the masses of quarks and the resulting mass ratios (in $\text{MeV}/c^2$ ) from Griffiths [8], Veltman [13], and <i>Physics Letters</i> [14].	40
TABLE 4:	The solution for the EMT group used as a basis for the two quark groups.	43
TABLE 5:	Solutions for the UCT and DSB groups which achieve the small mass ratios of the quarks' effective masses in baryons.	49
TABLE 6:	Solutions for the UCT and DSB groups which achieve the large mass ratios of the quarks' bare masses.	50
FIGURE 10:	The basic tau solution.	53
FIGURE 11:	The up solution.	54
FIGURE 12:	The charm solution.	55
FIGURE 13:	The top solution.	56
FIGURE 14:	The down solution.	57
FIGURE 15:	The strange solution.	58
FIGURE 16:	The bottom solution.	59
FIGURE 17:	The data flow between various programs.	68

## LIST OF FIGURES

FIGURE 1:	A comparison of the different approaches to the field/particle question.	4
FIGURE 2:	The first three Rosenstock eigensolutions.	21
FIGURE 3:	The second three Rosenstock eigensolutions.	22
FIGURE 4:	The Rosenstock equation with $y(0)=4.0$ .	24
FIGURE 5:	The Rosenstock equation with $y(0)=4.337387680187422$ .	25
FIGURE 6:	The Rosenstock equation With $y(0)=5.0$ .	26
FIGURE 7:	Data for the Rosen equations.	42
FIGURE 8:	The basic electron solution.	51
FIGURE 9:	The basic muon solution.	52
FIGURE 10:	The basic tau solution.	53
FIGURE 11:	The up solution.	54
FIGURE 12:	The charm solution.	55
FIGURE 13:	The top solution.	56
FIGURE 14:	The down solution.	57
FIGURE 15:	The strange solution.	58
FIGURE 16:	The bottom solution.	59
FIGURE 17:	The data flow between various programs.	68

## ACKNOWLEDGEMENTS

I would like to thank my supervisor, Dr. F. I. Cooperstock, for his support and encouragement.

I would also like to thank Diane Fox, Paul Lim, Joel Newman, and Bryce Young for much help, and the Natural Sciences and Engineering Research Council for financial support under grant #A5340.

## CHAPTER 1: INTRODUCTION

Physics now has well-developed theories of quantum and particle physics; they accurately predict many things about particles' interactions, but do not explain these particles in terms of more fundamental constructs nor how or why they have their observed properties. They take the fundamental particles, including leptons and quarks, to be "given". Nor do they even attempt to make such an explanation; such gaps are not deficiencies of quantum or nuclear physics, they simply reflect the fact that such questions are extra-theoretical for those disciplines. There are some current attempts to describe elementary particles in terms of simpler constituents ("rishons", "preons", etc.)<sup>1</sup> but these sub-units are not themselves explained or described in any more fundamental sense than the particles they supposedly make up.

---

<sup>1</sup> See Das et.al. [1] for more details.

Their hypothetical existence explains the patterns, but not the fundamental existence of matter.

The work presented here is the latest in a long series of attempts by physicists to determine the relationship between matter and fields. This particle/field duality has been a feature of physics for centuries.

# NUMERICAL SOLUTIONS OF A NON-LINEAR CLASSICAL FIELD THEORY OF LEPTONS AND QUARKS

## CHAPTER 1: INTRODUCTION

Physics now has well-developed theories of quantum and particle physics; they accurately predict many things about particles' interactions, but do not explain these particles in terms of more fundamental constructs nor how or why they have their observed properties. They take the fundamental particles, including leptons and quarks, to be "given". Nor do they even attempt to make such an explanation; such gaps are not deficiencies of quantum or nuclear physics, they simply reflect the fact that such questions are extra-theoretical for those disciplines. There are some current attempts to describe elementary particles in terms of simpler constituents ("rishons", "preons", etc.)<sup>1</sup> but these sub-units are not themselves explained or described in any more fundamental sense than the particles they supposedly make up.

---

<sup>1</sup> See Das et.al. [1] for more details.

<sup>2</sup> See, [2], page viii.

Their hypothetical existence explains the *patterns*, but not the fundamental *existence* of matter.

The work presented here is the latest in a long series of attempts by physicists to determine the relationship between *matter* and *fields*. This particle/field duality has been a feature of physics for centuries.

In the history of fundamental physics no other basic concepts have assumed a more important role than that of *fields* and *particles*. In fact it seems to be a strange characteristic of the human mind that it is forced to describe the physical properties of matter either as fields or as particles. The whole history of physics appears as a struggle to either clarify or escape from this *either or* dichotomy.<sup>2</sup>

On the one hand we have *matter* made up of particles, and on the other we have the concept of *fields*. In the course of scientific development there has been frequent controversy about the ontological priority of fields versus matter. Newtonian mechanics was eminently successful at using the particle and action-at-a-distance concepts. But then came James Clerk Maxwell's electromagnetism; he was a pioneer in maintaining that the electromagnetic field concept was a fundamentally important one. In his "A Dynamical Theory of the Electromagnetic Field", he ushers in the era of classical field theory; now the field becomes a real entity where energy can reside.

In speaking of the Energy of the field ... I wish to be understood literally. All energy is the same as mechanical energy, whether it exists in the form of motion or in that of elasticity, or in any other form. The energy in electromagnetic phenomena is mechanical energy. The only question is, where

---

<sup>2</sup> Sen, [2], page viii.

does it reside? On the old theories it resides in the electrified bodies, conducting circuits, and magnets, in the form of an unknown quality called potential energy, or the power of producing certain effects at a distance. On our theory it resides in the electromagnetic field, in the space surrounding the electrified and magnetic bodies, as well as in those bodies themselves....<sup>3</sup>

We can distinguish three types of approaches.<sup>4</sup> First, there are the *dualistic* theories, such as classical electromagnetism, which assume that the field-sources (i.e, the particles) are ontologically separate from the fields themselves. Secondly there are the non-dualistic theories which attempt to describe physics in terms of either pure fields or only particles. And thirdly there are the *unified non-dualistic* theories attempting to unify *all* fields or particles. Figure 1 illustrates some of these kinds of theories.

The work presented here is firmly in the non-dualistic, field-theory camp. It follows the vision of Einstein and Nathan Rosen who wrote:

A complete field theory knows only fields and not the concepts of particle and motion. For these must not exist independently of the field but are to be treated as part of it. On the basis of the description of a particle without singularity one has the possibility of a logically more satisfactory treatment of the combined problem: The problem of the field and that of motion coincide.<sup>5</sup>

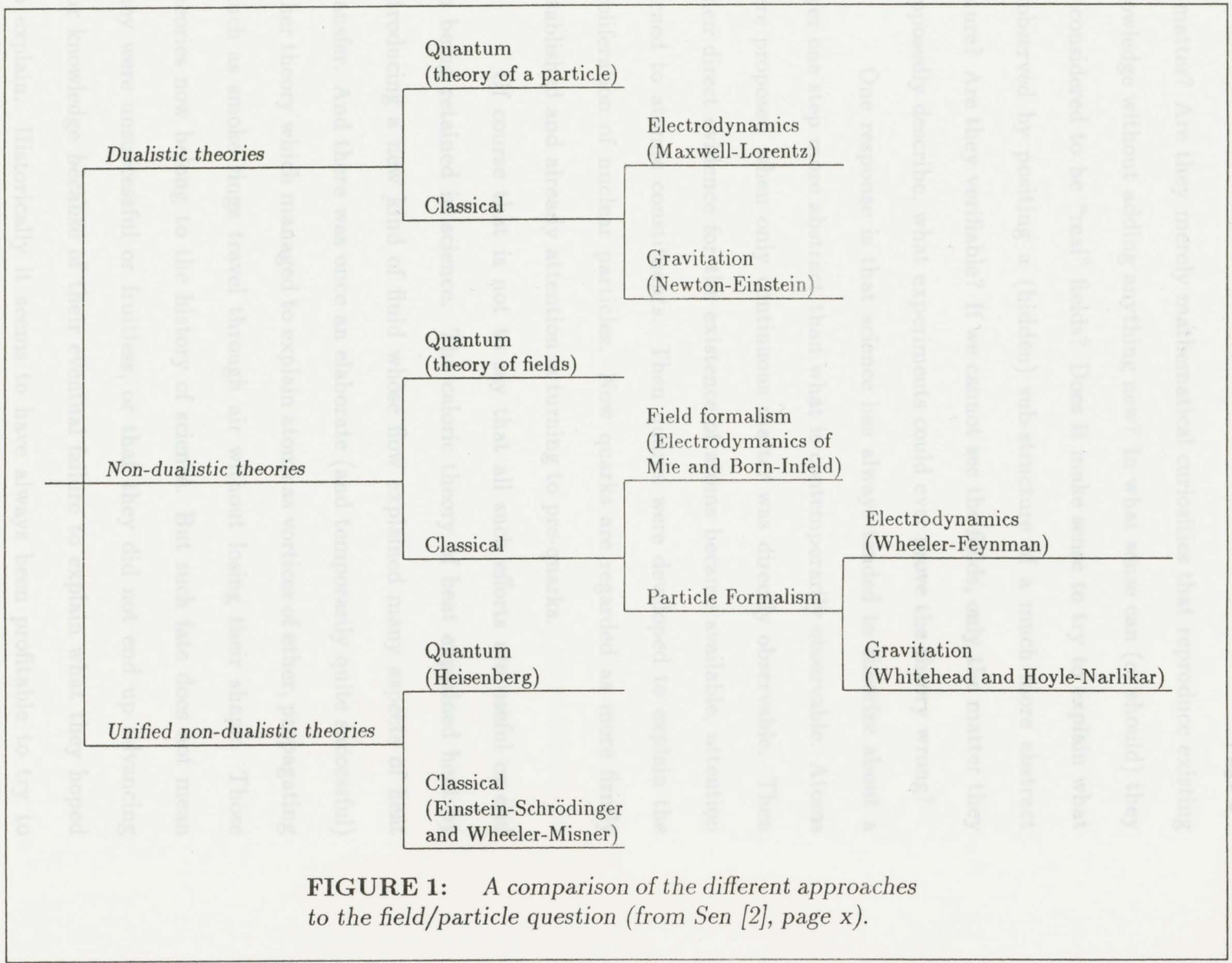
Naturally, such fields are abstract, remote, and (presently) immune from direct observation. So in what sense do they “explain” the existence

<sup>3</sup> James Clerk Maxwell, from “A Dynamical Theory of the Electromagnetic Field”, quoted in Williams, [3], page 134.

<sup>4</sup> (This classification scheme, and much of this section, is due to Sen [2], pp. viii-x.)

<sup>5</sup> Einstein and Rosen, [4], page 76.

FIGURE 1: A comparison of the different approaches to the field/particle question (from Sen [2], page x).



**FIGURE 1:** A comparison of the different approaches to the field/particle question (from Sen [2], page x).

of matter? Are they merely mathematical curiosities that reproduce existing knowledge without adding anything new? In what sense can (or should) they be considered to be "real" fields? Does it make sense to try to explain what is observed by positing a (hidden) sub-structure of a much more abstract nature? Are they verifiable? If we cannot see the fields, only the matter they supposedly describe, what experiments could ever prove the theory wrong?

One response is that science has always tended to theorize about a layer one step more abstract than what is contemporarily observable. Atoms were proposed when only continuous matter was directly observable. Then when direct evidence for the existence of atoms became available, attention turned to atomic constituents. Then quarks were developed to explain the proliferation of nuclear particles. Now quarks are regarded as more firmly established and already attention is turning to pre-quarks.

Of course that is not to say that all such efforts are useful or end up being retained in science. The caloric theory of heat explained heat by introducing a new kind of fluid whose flow explained many aspects of heat transfer. And there was once an elaborate (and temporarily quite successful) ether theory which managed to explain atoms as vortices of ether, propagating much as smoke rings travel through air without losing their shape. Those theories now belong to the history of science. But such fate does not mean they were unsuccessful or fruitless, or that they did not end up advancing our knowledge *because* of their eventual failure to explain what they hoped to explain. Historically it seems to have always been profitable to try to

guess what lies hidden beyond our current powers of observation. Whether the work presented here will ultimately be successful is presently unknown, but it is worth pursuing. And even though presumably the full explanation of matter will have to be quantized, it is hoped that pursuing the strictly *classical* approach as far as possible, as is done here, will lay the correct groundwork for the future.

## CHAPTER 2: PREVIOUS WORK

One attractive feature of the present work is that it needs a small number of assumptions and *ad hoc* hypotheses. Apart from numeric values (for mass, charge, and for various coupling constants whose magnitude is unknown), the main assumption used here is that the behaviour of matter can be described by an appropriate Lagrangian density which satisfies an action-extremizing principle. The correct Lagrangian must still be deduced (or guessed), but once that is done almost everything else flows automatically. In particular, as we shall see, the non-linearities of the fields involved here automatically produce their own field equations. Thus the beauty of the pure field theory approach is that it gives both the equations of the field development *and* the equations of motion of matter which is composed from the fields.

We must “plug in” from the outside values for the charge and mass of the particles described, but then the equations produce an independent prediction of the size of the particle. And as we have available experimentally-determined upper limits for the size of leptons and quarks, this provides an important independent check of the theory.

\* \* \*

<sup>1</sup>  $\partial_\mu$  and other conventions used here are defined in Appendix A.

## A) STANDARD ELECTROMAGNETISM

It will be useful to look first at how this works in the case of standard classical (dualistic) electromagnetism to show both similarities and differences with the later, non-dualistic, pure field-theory approach. For example, for the simple case of charged particles in an electromagnetic field we have

CHAPTER 2: PREVIOUS WORK

Classical field theory starts by positing a Lagrangian density ( $L$ ) for a given system (where  $L$  is a function of the fields and their first partial derivatives), and assumes that the action

$$S = \int L dV dt$$

satisfies an extremal principle

$$\delta S = 0. \quad (1)$$

If we let  $B$  represent the fields, so that  $L = L(B, \partial_u B)$  then<sup>6</sup> in the same way that in classical mechanics Hamilton's principle yields Lagrange's equations, (1) yields the "Euler-Lagrange" equations:

$$\left( \frac{\partial L}{\partial B_{,a}} \right)_{;a} - \frac{\partial L}{\partial B} = 0. \quad (2)$$

Also,

$$T_{ij} = \left( \frac{\partial L}{\partial g^{ij}} \right) - \frac{L g_{ij}}{2} \quad (3)$$

where  $T_{ij}$  is the "energy-momentum" tensor and  $g_{ij}$  is the metric tensor. In particular,  $T_{00}$  is the energy density of the system.

<sup>6</sup>  $\partial_u$  and other conventions used here are defined in Appendix A.

## A) STANDARD ELECTROMAGNETISM

It will be useful to look first at how this works in the case of standard, classical (dualistic) electromagnetism to show both similarities and differences with the later, non-dualistic, pure field-theory approach. For example, for the simple case of charged particles in an electromagnetic field we have

$$L_{\text{EM}} = -\frac{F_{ab}F^{ab}}{8\pi} - j_a A^a$$

where

$$F_{kl} \equiv A_{l,k} - A_{k,l}$$

$$\mathbf{j} \equiv \rho(1, \vec{v})$$

$\rho \equiv$  the charge density

$\vec{v} \equiv$  velocity of the charged particles

$$\mathbf{A} \equiv (\varphi, \vec{A})$$

$\varphi \equiv$  the scalar electromagnetic potential

$\vec{A} \equiv$  the vector electromagnetic potential.

So now (2) yields

$$F^{ka}{}_{;a} = -4\pi j^k \quad (4)$$

(which is in fact the tensor formulation of two of Maxwell's equations). Note that here  $\mathbf{j}$ , the current four-vector, must be defined *outside* the field theory. Both it and the concept of a charged mass-point itself are *auxiliary* concepts, thus illustrating the fact the classical electromagnetism is indeed a dualistic theory.

## B) MIE, BORN AND INFELD

Classical electromagnetism with its Maxwell's equations has great power, but there are two fundamental difficulties. Firstly, the concept of a point-charge leads to difficulties of infinite self-energies. Secondly, the field equations break down at the point-charge itself because of the singularity there. So, as we have seen, the equations of motion have to be added separately from the field equations.

G. Mie [5] proposed a non-dualistic field theory where the particle's mass and charge are determined solely by the fields. But there are two criticisms of his work.<sup>7</sup> One is that Mie's Lagrangian does not provide unique solutions for each particle. Another is that his theory is not gauge-invariant; because his Lagrangian and his field equations both depend explicitly on  $\varphi$ , adding a constant external potential to  $\varphi$  changes the solution.

When it was determined that quantum electrodynamics also had difficulties with the notion of a point-electron, Born and Infeld [6] tried to develop a non-linear theory of electrodynamics. However Rosen [7] showed that there are still singularities present at the origin.

## C) ROSEN

Rosen [7] concluded that  $L$  should depend on the potentials, to avoid the singularities of Born and Infeld, but must somehow involve other variables which will make the Lagrangian gauge-invariant.

---

<sup>7</sup> Sen, [2], page 41.

He investigated a Lagrangian formed by combining an electromagnetic field with a complex, scalar matter-field  $\psi$ , and demanded that gauge invariance hold when the electromagnetic vector potential undergoes a local gauge transformation

$$A_u \rightarrow A'_u = A_u + \frac{\partial \lambda}{\partial x^u} \quad (5)$$

(where  $\lambda$  is an arbitrary function of the space-time coordinates). He found he could make  $L$  gauge-invariant by imposing the condition that under the gauge transformation (5) we also have

$$\psi \rightarrow \psi' = e^{i\epsilon\lambda(x)}\psi$$

(where  $\epsilon$  is a real constant of dimensions  $1/\text{charge}$ .<sup>8</sup>) Then  $\psi\psi^*$  is gauge-invariant since the phase factors cancel out. ( $F^{uv}$  is already gauge-invariant since the second derivatives of  $\lambda$  cancel out.) But if  $L$  contains factors like  $\partial_u\psi$  they will *not* be gauge-invariant since

$$\partial_u\psi \rightarrow e^{i\epsilon\lambda}(\partial_u + i\epsilon\lambda_{,u})\psi.$$

So, to cancel out these extra  $\lambda_{,u}$  terms,  $\partial_u$  must be replaced by the so-called "gauge-covariant derivative"<sup>9</sup>

$$D_u \equiv \partial_u - i\epsilon A_u.$$

Then

$$\begin{aligned} (D_u\psi) &\rightarrow [\partial_u - i\epsilon(A_u + \lambda_{,u})]e^{i\epsilon\lambda}\psi \\ &= e^{i\epsilon\lambda}(D_u\psi) \end{aligned}$$

<sup>8</sup> (The dimensions of a number of quantities are summarized for convenience in Appendix B.)

<sup>9</sup> Griffiths, [8], page 349.

and so factors like  $(D^a\psi)(D_a\psi)^*$  are gauge-invariant. Rosen combined the simpler possible scalars (omitting more complicated ones like  $F_{ab}F^{ab}\psi\psi^*$ ) to form

$$L_{\text{ROSEN}} = -\frac{F_{ab}F^{ab}}{8\pi} - (D^a\psi)(D_a\psi)^* + \sigma^2\psi\psi^*$$

where  $\sigma$  is a (real) constant. (The signs of the last two terms were chosen as shown as it was found that only this choice of sign structure gave particle-like solutions.) Variation with respect to  $A_u$  and  $\psi^*$  then gives two field equations

$$F^{ua}{}_{;a} = -4\pi J^u \quad (6)$$

$$\left( \text{where } J^u \equiv \frac{1}{2}i\epsilon [\psi^*(D^u\psi) - \psi(D^u\psi)^*] \right)$$

$$(D^a\psi)_{;a} - i\epsilon A_a(D^a\psi) + \sigma^2\psi = 0.$$

Notice that now the current four-vector  $\mathbf{J}$  is produced by the fields themselves; the theory is self-contained. It then follows via (3) that the energy-momentum tensor is

$$T_{uv} = \frac{1}{16\pi}(g_{uv}F_{ab}F^{ab} - 4F_{ua}F_v{}^a) + \frac{1}{2}\left[ g_{uv}\{(D^a\psi)(D_a\psi)^* - \sigma^2\psi\psi^*\} - (D_u\psi)(D_v\psi)^* - (D_u\psi)^*(D_v\psi) \right].$$

For the spherically-symmetric case we can use the gauge freedom of (5) to choose a gauge where  $\vec{A} = 0$  so we have

$$\mathbf{A} = (\varphi(r), \vec{0}).$$

Rosen now (implicitly) assumes that the  $\mathbf{J}$  of (6) should be identified with the usual four-vector current  $\mathbf{j}$  of standard electromagnetic theory. (This is reasonable since we want Maxwell's equations (4) to hold in both cases.) So

$$J^0 \equiv \rho(r),$$

and, for the *static* case where the charge does not alter with time,  $\vec{v} = 0$  so

$$J^\alpha \equiv v^\alpha = 0 \quad (\alpha = 1, 2, 3).$$

Rosen takes  $\psi$  to be of the form

$$\psi = \theta(r)e^{-i\epsilon\mu t}.$$

Then the field equations (6) become

$$\nabla^2\theta + [\epsilon^2(\varphi + \mu)^2 - \sigma^2]\theta = 0$$

$$\nabla^2\varphi + 4\pi\rho = 0$$

$$\left(\text{where } \rho = \epsilon^2(\varphi + \mu)\theta^2\right).$$

For solutions to represent a particle some boundary conditions must be imposed. To ensure that there are no discontinuities at the origin we require that

$$\frac{d\theta}{dr} = \frac{d\varphi}{dr} = 0 \quad \text{at } r = 0.$$

We also require that

$$\theta \rightarrow 0 \text{ exponentially as } r \rightarrow \infty$$

so that the particle is concentrated near the origin, and require that

$$\varphi \rightarrow \frac{\text{constant}}{r} \quad \text{as } r \rightarrow \infty$$

so the electric field behaves like a Coulomb field at large distances from the particle.

It is easier to work with dimensionless variables by defining

$$x \equiv r \sigma, \quad z \equiv \frac{\epsilon}{\sigma}(\varphi + \mu), \quad y \equiv \sqrt{4\pi} \epsilon \frac{\theta}{\sigma}. \quad (7)$$

Then the field equations are

$$y'' + \frac{2}{x}y' - y + yz^2 = 0 \quad (8)$$

$$z'' + \frac{2}{x}z' + y^2z = 0.$$

The total charge ( $q$ ) of the particle is

$$q = \int_{\text{all space}} \rho(r) dV \quad (9)$$

$$= 4\pi \int_0^\infty \rho r^2 dr$$

$$= 4\pi \int_0^\infty \epsilon^2(\varphi + \mu)^2 \theta^2 r^2 dr$$

$$= \frac{\alpha}{\epsilon}$$

$$\left(\text{where } \alpha \equiv \int_0^\infty y^2 z x^2 dx\right).$$

The energy density is

$$T_{00} = \frac{1}{8\pi} \left(\frac{d\varphi}{dr}\right)^2 - \frac{1}{2} \left[ \epsilon^2(\varphi + \mu)^2 \theta^2 - \left(\frac{d\theta}{dr}\right)^2 - \sigma^2 \theta^2 \right].$$

By integrating by parts and using the field equations we can find the total energy ( $E$ ) is

$$\begin{aligned}
 E &= \int_{\text{all space}} T_{00} dV & (10) \\
 &= 4\pi \int_0^\infty T_{00} r^2 dr \\
 &= -\frac{\sigma w}{2\epsilon^2}
 \end{aligned}$$

where

$$w \equiv \alpha\beta + \gamma, \quad \beta \equiv \frac{\epsilon\mu}{\sigma}, \quad \gamma \equiv \int_0^\infty y^2 z^2 x^2 dx.$$

We require  $\beta$  be less than one for the following reason. From (8), the field equation for  $y$ , we have

$$y'' + \frac{2}{x}y' + [z^2 - 1]y = 0.$$

And we also require that

$$\varphi \rightarrow \varphi_{\text{coulombic}} = \frac{q}{r} \quad \text{as } r \rightarrow \infty.$$

So, using (9),

$$\begin{aligned}
 z &= \frac{\epsilon(\varphi + \mu)}{\sigma} \rightarrow \frac{\epsilon q}{\sigma r} + \frac{\epsilon\mu}{\sigma} & (11) \\
 &= \frac{\alpha}{x} + \beta.
 \end{aligned}$$

As  $x \rightarrow \infty$ ,  $z \rightarrow \beta$  so if we define

$$C \equiv (z^2 - 1)$$

then

$$C \rightarrow (\beta^2 - 1) \text{ as } x \rightarrow \infty.$$

So as  $x \rightarrow \infty$  we have

$$y'' + \frac{2}{x}y' + Cy = 0.$$

If  $\beta^2 > 1$  then  $C > 0$  and

$$y(x) = A \frac{\sin(\sqrt{C}x)}{x} + B \frac{\cos(\sqrt{C}x)}{x} \quad (A, B = \text{constants}). \quad (12)$$

If  $\beta^2 < 1$  then  $C < 0$  and

$$y(x) = A \frac{e^{\sqrt{-C}x}}{x} + B \frac{e^{-\sqrt{-C}x}}{x}. \quad (13)$$

Of these two solutions, (12) has oscillatory behaviour whereas (13) has the desired exponentially-decaying solution, so we need  $\beta^2 < 1$ .

This model will produce particles, in the sense that there are localized solutions satisfying the boundary conditions, but all the solutions found have  $w > 0$ , and hence, via (10) the total energy is *negative*. And since mass is related to energy via  $E = mc^2$ , to produce positive-mass particles a more complicated Lagrangian must be used.

#### D) ROSEN AND ROSENSTOCK

Rosen and Rosenstock [9] investigated the case of the purely scalar field, without electromagnetism. They took

$$L_{RR} = (\partial^a \psi)(\partial_a \psi)^* - \sigma^2 \psi \psi^* + \frac{1}{2} g \psi^2 \psi^{*2}$$

(where  $g$  is a positive constant). For the static, spherically-symmetric case they take

$$\psi = \theta(r) e^{i\epsilon\mu t}$$

(where  $\epsilon\mu$  is a real constant). The field equation is then

$$\frac{d^2\theta}{dr^2} + \frac{2}{r} \frac{d\theta}{dr} - [\sigma^2 - (\epsilon\mu)^2]\theta + g\theta^3 = 0 \tag{14}$$

(where  $[\sigma^2 - (\epsilon\mu)^2]$  is assumed to be positive to prevent oscillatory solutions). The energy density is

$$T_{00} = \frac{1}{2} \left[ (\epsilon\mu)^2 \theta^2 + \sigma^2 \theta^2 + \left( \frac{d\theta}{dr} \right)^2 - \frac{1}{2} g \theta^4 \right].$$

Again, integrating by parts yields a total energy

$$E = 2\pi \int_0^\infty \left[ 2(\epsilon\mu)^2 \theta^2 + \frac{1}{2} g \theta^4 \right] r^2 dr$$

which is positive definite, so we have positive-energy particles. Rosen and Rosenstock identify  $\mathbf{J}$ , the current four-vector arising from the field equations, with  $\mathbf{j}$ , the current four-vector of standard electromagnetic theory, the same way as was done before (in Section C, page 11). They then identify a quantity they call "charge" via

$$\text{"charge"} \equiv \int_{\text{all space}} \rho dV = \int J^0 dV.$$

If  $\psi$  is real then the charge is zero, but if  $\psi = \theta e^{i\epsilon\mu t}$  then

$$\text{"charge"} = 2\epsilon^2 \mu \int \theta^2 dV.$$

But this scalar "charge" is *not* the usual electromagnetic charge. Whereas before  $\mathbf{J}$  was identified with  $\mathbf{j}$  through the appearance of Maxwell's equations, here there is not the same incentive. So this  $L$  is successful at producing *neutral* particles, but not really satisfactory for modelling *charged* particles.

However, it is still useful to look at the solutions to (14) since many of the properties of this field equation are also characteristics of the more complicated field equations which will appear later. Using the same dimensionless variables the Rosenstock equation (14) becomes

$$y'' + \frac{2}{x}y' - y + y^3 = 0. \quad (15)$$

We are interested in the case where we only have one free parameter:  $y(0)$ . (The other boundary condition is  $y'(0)=0$  which is imposed to ensure that  $\theta'(r=0)=0$  for smoothness at the origin.) This is an interesting equation because it looks deceptively simple yet its non-linearity produces some interesting behaviour. What are the solutions to this equation?

First, we can note that (15) has three constant solutions:

$$y = 0, \quad y = -1, \quad y = +1.$$

We are interested in solutions where the particle is concentrated in a localized region near the origin and has very little mass far away from the origin. The mass depends on  $T_{00}$ , which depends on  $y$ , so we need to find solutions where  $y \rightarrow 0$  as  $x \rightarrow \infty$ .

We can deal with this equation in two ways. Firstly, we can set  $y(0)$  to our desired start parameter and then integrate  $y$  numerically along the

$x$ -axis to find  $y(x)$  for any desired  $x > 0$  value (subject to the limitations of numerical accuracy to be discussed below). Secondly, we can treat the equation analytically by making some linear approximations for asymptotic solutions for large  $x$ .

We can start by looking at solutions where  $y$  approaches one of the constant solutions as  $x$  grows large.

First, if we look at the behaviour of  $y$  where  $y \rightarrow \pm 1$  as  $x \rightarrow 0$  we can set

$$y = \pm 1 + u(x) \quad (\text{where } u(x) \text{ is small})$$

then (15) approaches

$$x^2 u'' + 2x u' + (2x^2)u = 0$$

which can be solved (via spherical Bessel functions) to show

$$y(x) = A \frac{\sin(\sqrt{2}x)}{x} + B \frac{\cos(\sqrt{2}x)}{x} \quad (A, B = \text{constants}).$$

So  $y \rightarrow \pm 1$  with a decaying oscillation around the  $\pm 1$  value. These obviously are not suitable candidates for particle-solutions, but we will see below that numerically we often run into these oscillatory solutions because of numerical instability.

More interesting for us are the particle-like solutions where  $y \rightarrow 0$  as  $x \rightarrow \infty$ . In this case the  $(y)^3$  term in (15) is negligible for large  $x$  and we have

$$y'' + \frac{2}{x}y' - y = 0$$

which has the (exact) solution

$$y = A \frac{e^x}{x} + B \frac{e^{-x}}{x} \quad (A, B = \text{constants}). \quad (16)$$

Analytically we can set  $A=0$  and assume the  $e^{-x}/x$  term is our desired solution, but numerically this presents some problems as we will see below.

The important fact is that there are these particle-like solutions to (15) where  $y$  goes exponentially to 0 as  $x$  goes to  $\infty$ . But how many of these particle-like solutions are there? And what do they look like?

Finkelstein, LeLevier, and Ruderman [10] indicated that there are quantized "eigensolutions" which have the following characteristics: if  $y(0)$  is a very specific value (an eigenvalue) then  $y \rightarrow 0$ , but if  $y(0)$  is raised even slightly above or lowered slightly below this value then  $y \rightarrow \pm 1$  as  $x \rightarrow \infty$ . Furthermore there are an infinite number of such eigensolutions. And each eigensolution can be characterized by the number of "nodes" it has (i.e., how many times it crosses the  $x$ -axis). It turns out the  $n$ th eigensolution has exactly  $n$  nodes.

Table 1 shows the first fifty eigenvalues and Figures 2 and 3 show the  $n = 1$  to  $n = 6$  eigensolutions.

Numerically this sort of equation presents certain problems which occur both here and in the full model. For the solutions (16) we can analytically set  $A=0$  and thus ensure that  $y \rightarrow 0$  as  $x \rightarrow 0$ . But when integrating (16) numerically from  $y(0)$  we cannot numerically set  $A=0$  and so eventually round-off error and lack of precision accumulate to the point where the  $e^x/x$  term starts

$n$	$y(0)_n$	$n$	$y(0)_n$
1	4.33738	26	921.95804
2	14.10358	27	941.66012
3	29.13121	28	959.70671
4	49.36079	29	976.30937
5	74.77307	30	991.66159
6	105.38181	31	1005.92686
7	141.16164	32	1019.23873
8	181.91134	33	1031.72519
9	226.57026	34	1043.54078
10	279.37381	35	1054.87351
11	335.02740	36	1065.90501
12	393.64247	37	1076.70618
13	452.41196	38	1087.07776
14	508.20404	39	1096.53831
15	559.10233	40	1104.70036
16	604.71920	41	1111.56826
17	645.61402	42	1117.37910
18	683.52714	43	1122.39493
19	721.00998	44	1126.82993
20	757.24527	45	1130.83396
21	791.07109	46	1134.49822
22	822.21383	47	1137.88114
23	850.71010	48	1141.03927
24	876.70201	49	1144.03985
25	900.37976	50	1146.95335

**TABLE 1:** *The  $y(0)$  values for the first fifty eigen-solutions for the Rosenstock equation.*

to show up and the solution “blows up” as  $x$  gets large. So numerically we must take care to adjust the step size of integration and carry enough precision internally in the computer programs to suppress the  $e^x/x$  term until we have integrated far enough down the  $x$ -axis to extract the information we need (i.e.,

Figure 2: The first three Rosenstock eigensolutions

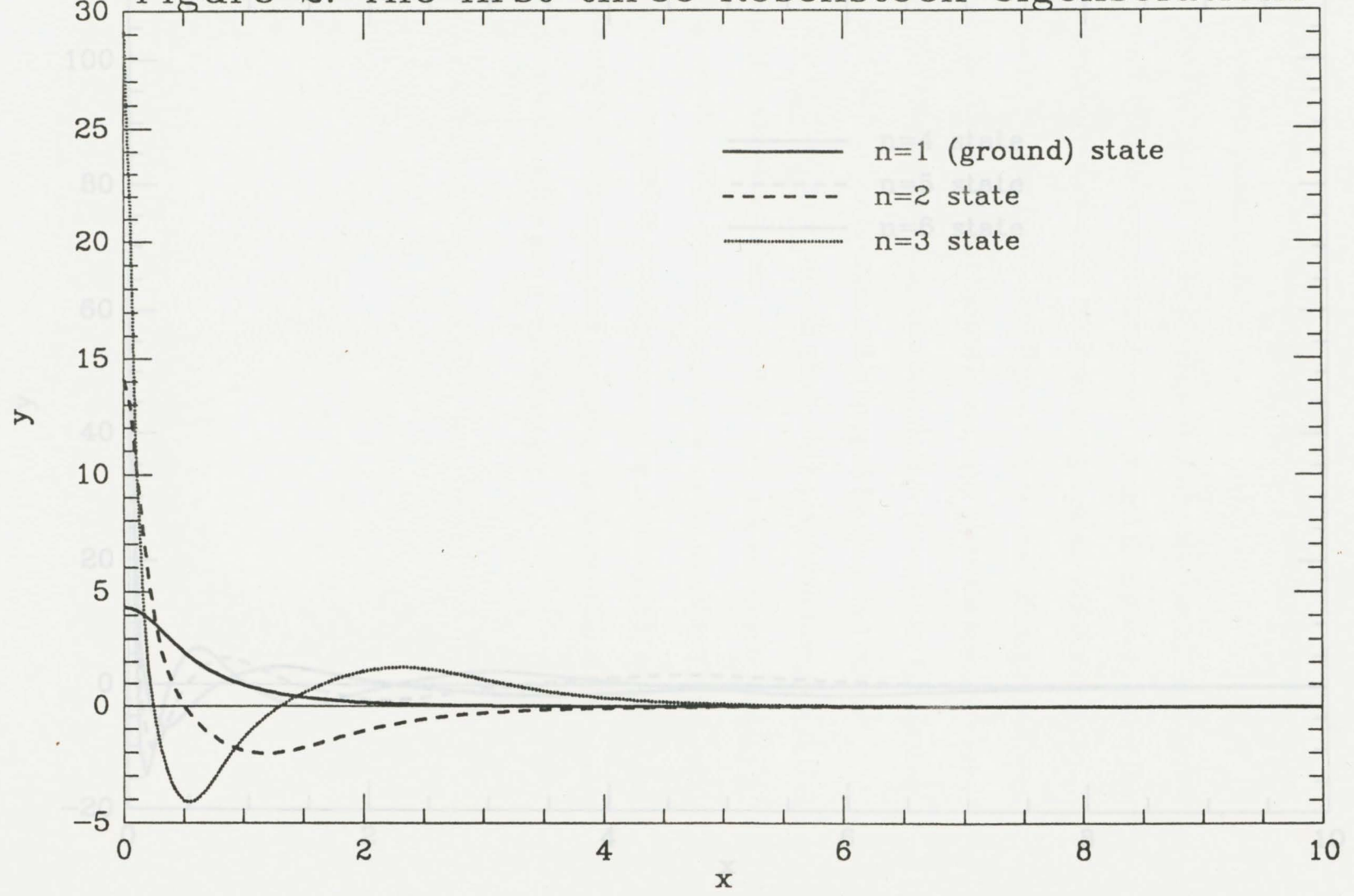
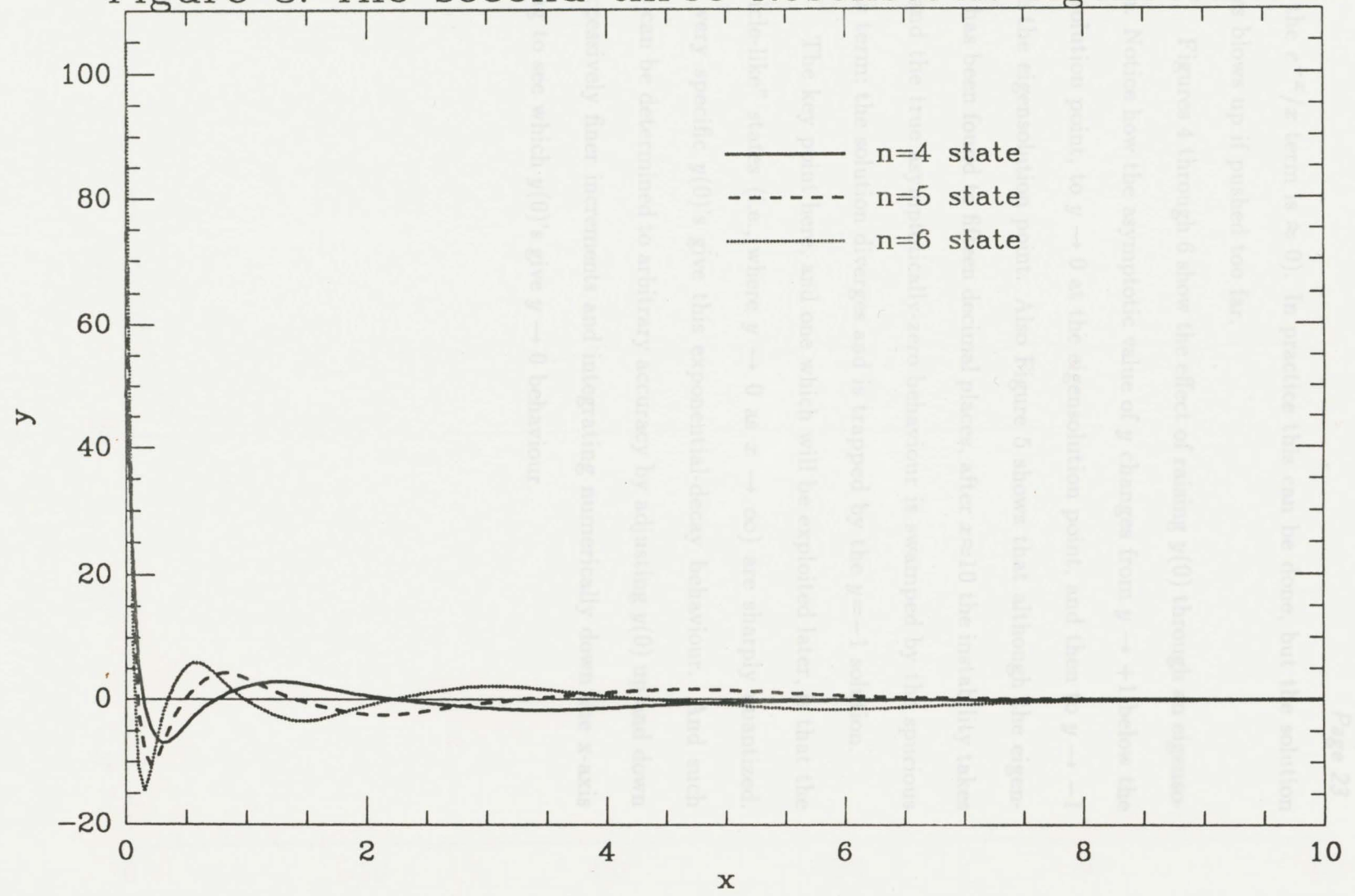


Figure 3: The second three Rosenstock eigensolutions



until the  $e^{-x}/x$  term is  $\approx 0$ ). In practice this can be done, but the solution always blows up if pushed too far.

Figures 4 through 6 show the effect of raising  $y(0)$  through an eigensolution. Notice how the asymptotic value of  $y$  changes from  $y \rightarrow +1$  below the eigensolution point, to  $y \rightarrow 0$  at the eigensolution point, and then to  $y \rightarrow -1$  above the eigensolution point. Also Figure 5 shows that although the eigenvalue has been found to fifteen decimal places, after  $x \approx 10$  the instability takes over and the true asymptotically-zero behaviour is swamped by the spurious  $Ae^x/x$  term; the solution diverges and is trapped by the  $y = -1$  solution.

The key point here, and one which will be exploited later, is that the “particle-like” states (i.e., where  $y \rightarrow 0$  as  $x \rightarrow \infty$ ) are sharply quantized. Only very specific  $y(0)$ 's give this exponential-decay behaviour. And such  $y(0)$ 's can be determined to arbitrary accuracy by adjusting  $y(0)$  up and down in successively finer increments and integrating numerically down the  $x$ -axis looking to see which  $y(0)$ 's give  $y \rightarrow 0$  behaviour.

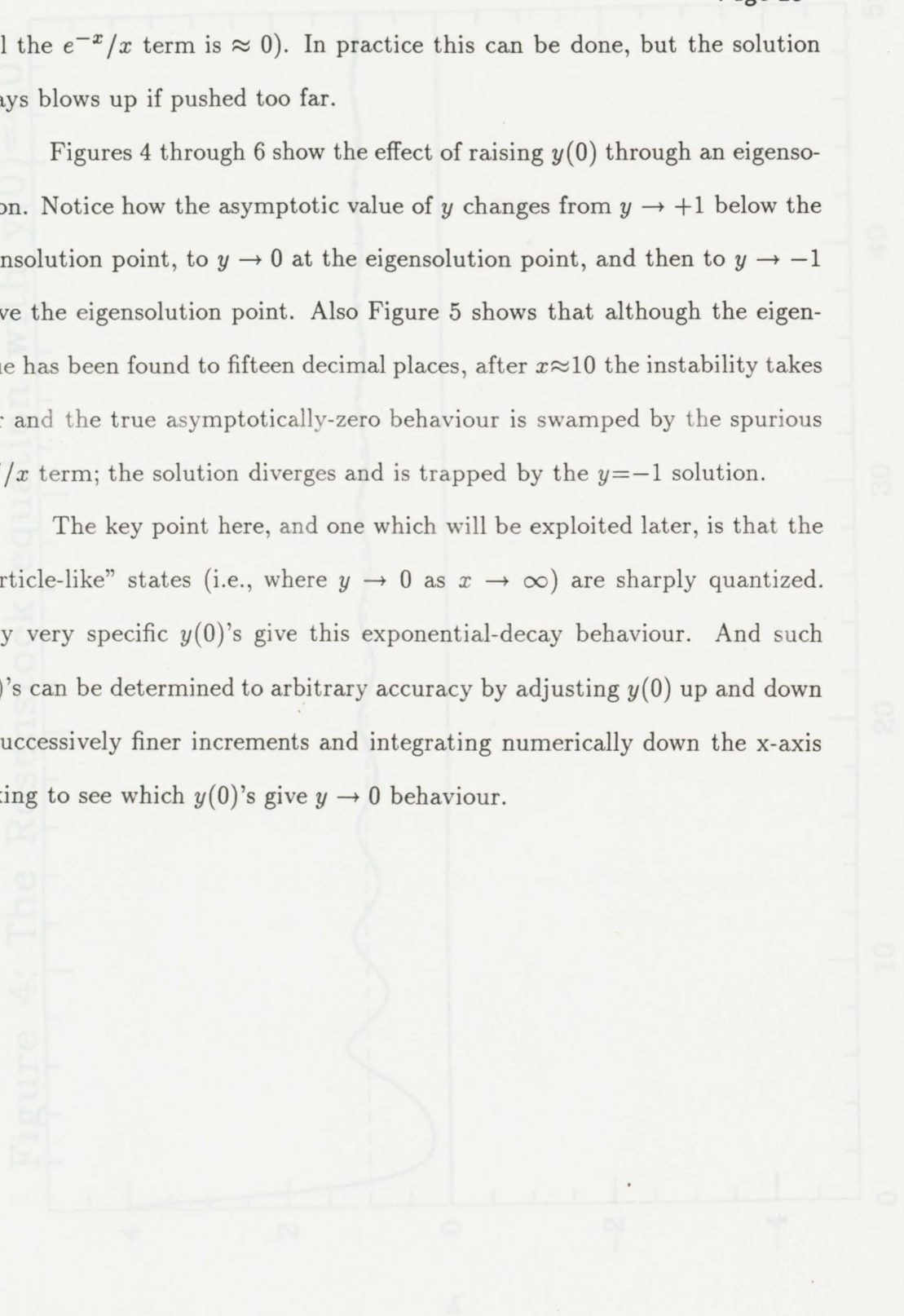


Figure 4: The Rosenstock equation with  $y(0)=4.0$

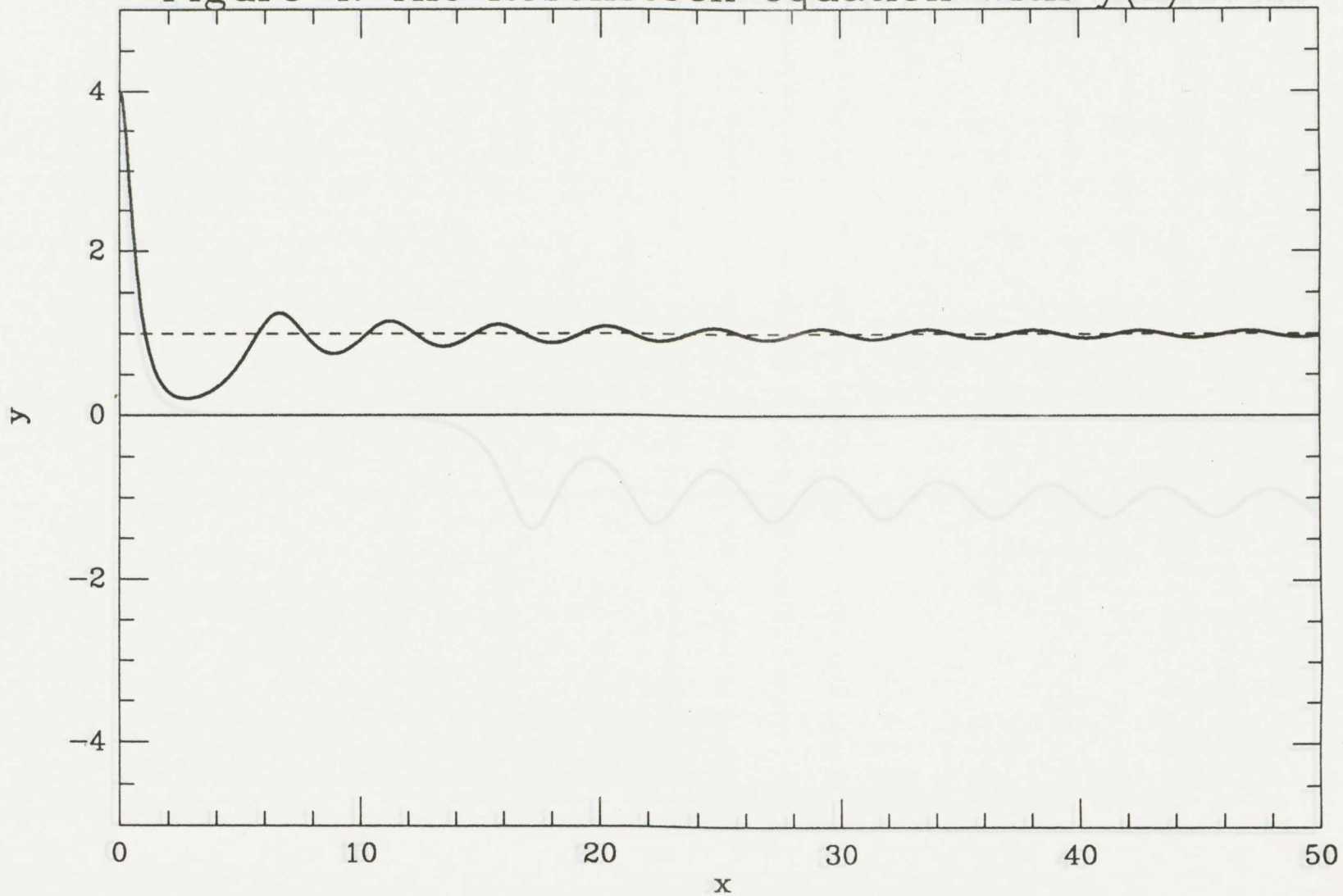


Figure 5: The Rosenstock equation,  $y(0)=4.337387680187422$

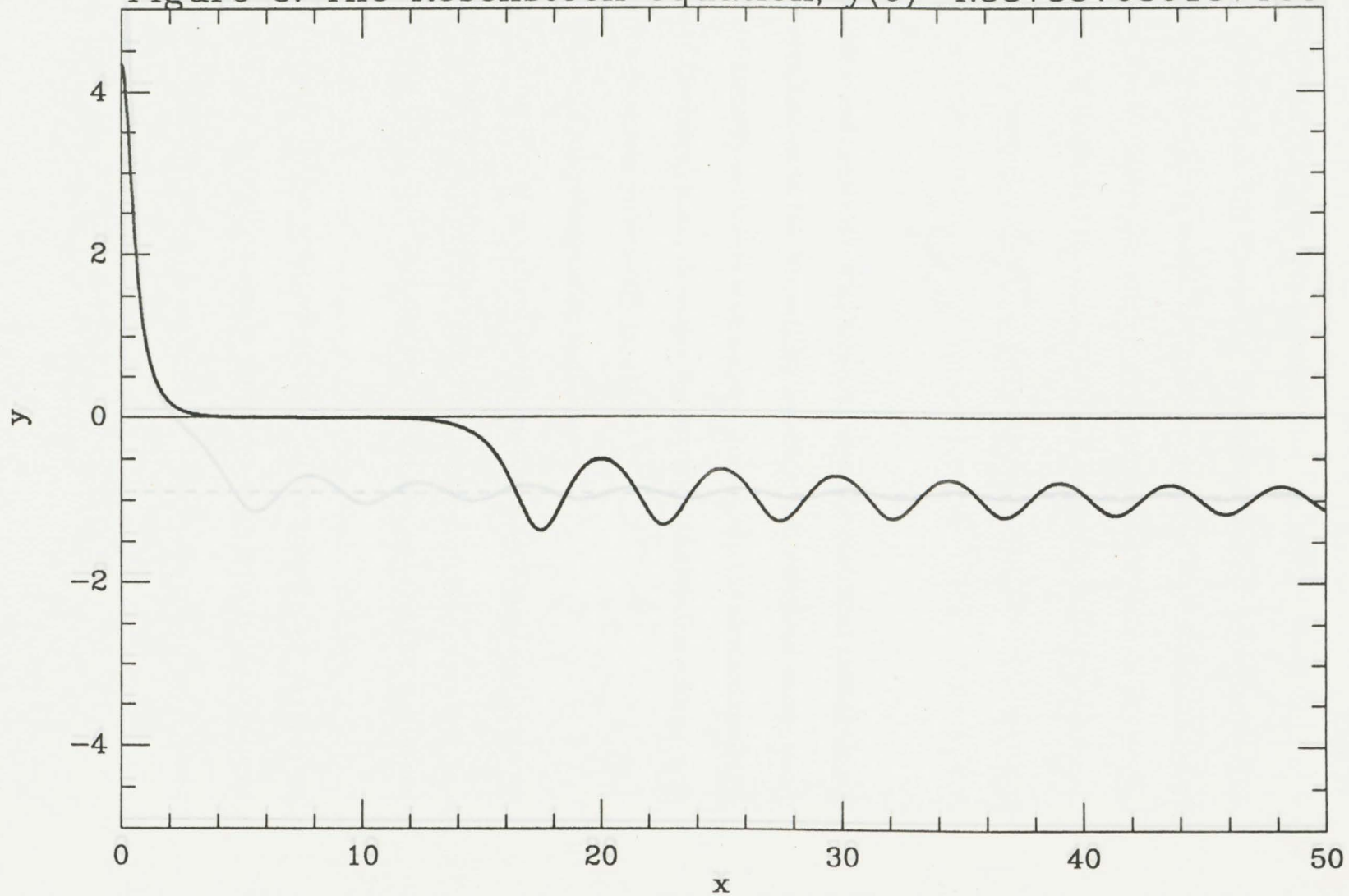
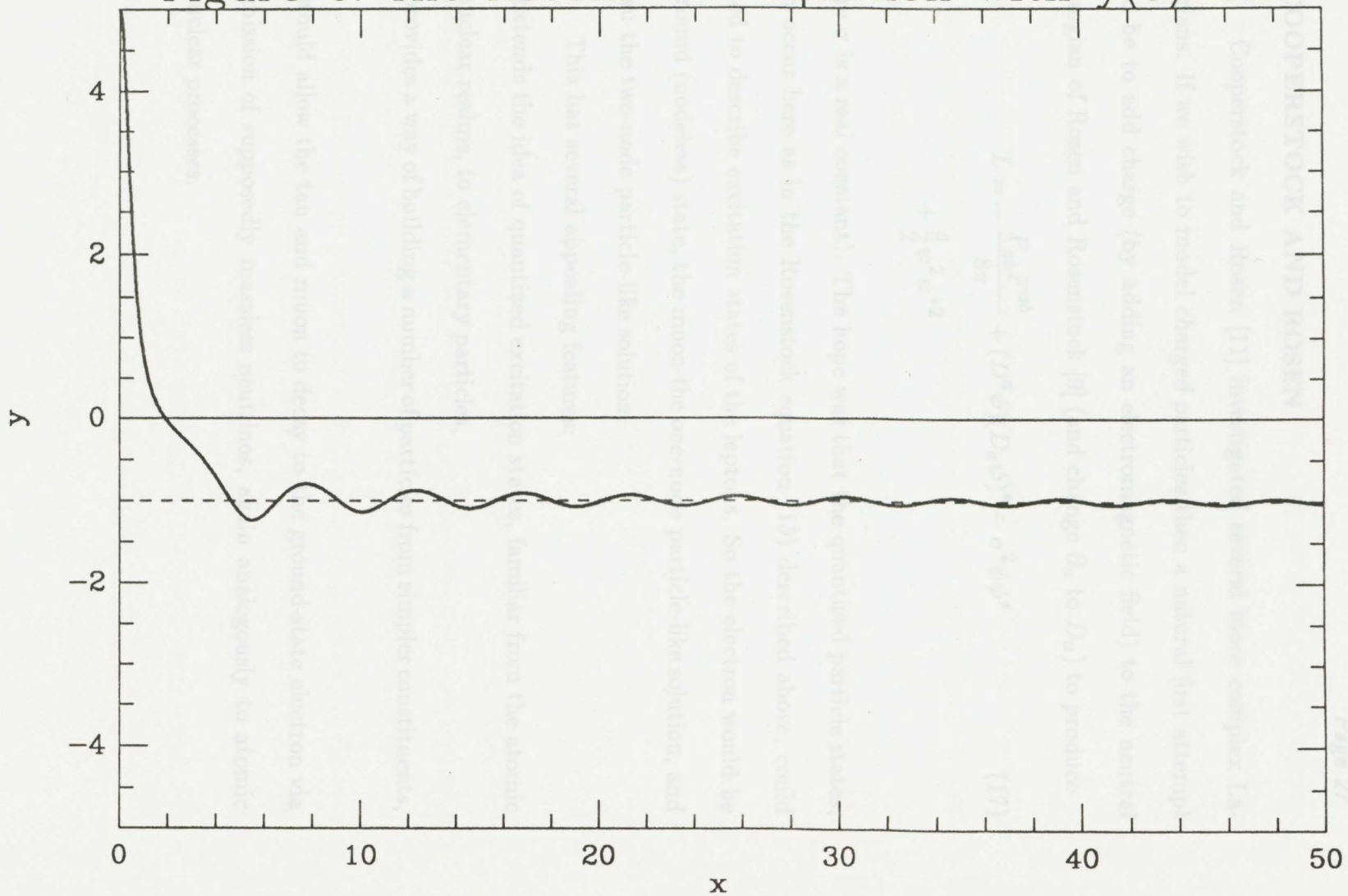


Figure 6: The Rosenstock equation with  $y(0)=5.0$



## E) COOPERSTOCK AND ROSEN

Cooperstock and Rosen [11] investigated several more complex Lagrangians. If we wish to model *charged* particles then a natural first attempt would be to add charge (by adding an electromagnetic field) to the neutral Lagrangian of Rosen and Rosenstock [9] (and change  $\partial_u$  to  $D_u$ ) to produce

$$L = -\frac{F_{ab}F^{ab}}{8\pi} + (D^a\psi)(D_a\psi)^* - \sigma^2\psi\psi^* + \frac{g}{2}\psi^2\psi^{*2} \quad (17)$$

(where  $\sigma$  is a real constant). The hope was that the quantized particle states, which occur here as in the Rosenstock equation (15) described above, could be used to describe excitation states of the leptons. So the electron would be the ground (nodeless) state, the muon the one-node particle-like solution, and the tau the two-node particle-like solution.

This has several appealing features:

- a) it extends the idea of quantized excitation states, familiar from the atomic and nuclear realms, to elementary particles,
- b) it provides a way of building a number of particles from simpler constituents, and
- c) it would allow the tau and muon to decay to the ground-state electron via the emission of supposedly massless neutrinos, again analogously to atomic and nuclear processes.

<sup>10</sup> Quigg, [12], page 3.

However, there is a problem with this Lagrangian. For all the models, the fields approach zero very quickly. Figure 2 shows that for the Rosen-Rosenstock equation (16), for example,  $y \approx 0$  when  $x > \sim 5$  so  $x$  is on the order of the particle size. And since

$$r = \frac{x}{\sigma}$$

$1/\sigma$  is also a dimension on the order of the particle size. And experimentally it is known that electrons are not bigger than  $\sim 10^{-16}$  cm.<sup>10</sup> But (17) gives  $1/\sigma$  of  $10^{-12}$  cm, which is too big.

Another Lagrangian which ultimately proved unsatisfactory, but which has some attractive features, is

$$L = -\frac{F_{ab}F^{ab}}{8\pi} - (D^a\psi)(D_a\psi)^* + \sigma^2\psi\psi^* \quad (18)$$

$$+ (\partial^a\psi_1)(\partial_a\psi_1)^* - \sigma^2\psi_1\psi_1^* + \frac{1}{2}g\psi_1^2\psi_1^{*2}$$

$$- f\psi\psi^*\psi_1\psi_1^*.$$

This combines the scalar field ( $\psi$ ) and electromagnetic field of Rosen [7] with the scalar field (now denoted  $\psi_1$ ) of Rosen and Rosenstock [9]. In addition, a simple coupling term  $f\psi\psi^*\psi_1\psi_1^*$  (where  $f$  is a constant) has been added.

With the correct parameters this Lagrangian produces an electron well within the upper size limit, but when  $\psi_1$  is excited, Cooperstock and Rosen found this Lagrangian cannot represent the muon since the particle size predicted is too large. Also, no single set of parameters was found which will give the correct tau mass when  $\psi_1$  is doubly excited.

<sup>10</sup> Quigg, [12], page 3.

However there are two features which will be used later. The negative coupling term ( $-f\psi\psi^*\psi_1\psi_1^*$ ) adds a negative term to the energy integral of the form

$$-\int_0^\infty y^2(y_1)^2 x^2 dx.$$

This binding energy term grows larger in absolute value as  $y_1$  is excited, thereby holding the particle constituents together, even when highly-excited. (It would be a problem for this theory if the fields were free to separate since one of the free fields would be the negative-energy  $\psi$  scalar field.) Also, via (9), the charge is proportional to

$$\alpha = \int_0^\infty y^2 z x^2 dx$$

which depends mainly on the charged scalar  $\psi$  (i.e.,  $y$ ) and only slightly, via the coupling, on the neutral scalar  $\psi_1$  (i.e.  $y_1$ ). So it is numerically easy to adjust the parameters to keep the charge conserved for all three states in a family.

Cooperstock and Rosen achieved their goal of modelling the complete lepton family by adding another neutral scalar field ( $\psi_2$ ) to the two previous fields: the charged field ( $\psi$ ) and the first neutral scalar field ( $\psi_1$ ). So

$$\begin{aligned} L_{CR} = & -\frac{F_{ab}F^{ab}}{8\pi} - (D^a\psi)(D_a\psi)^* + \sigma^2\psi\psi^* \\ & + (\partial^a\psi_1)(\partial_a\psi_1)^* - \sigma^2\psi_1\psi_1^* + \frac{1}{2}g_1\psi_1^2\psi_1^{*2} \\ & + (\partial^a\psi_2)(\partial_a\psi_2)^* - \sigma^2\psi_2\psi_2^* + \frac{1}{2}g_2\psi_2^2\psi_2^{*2} \\ & - f_1\psi\psi^*\psi_1\psi_1^* - f_2\psi\psi^*\psi_2\psi_2^* - f_3\psi_1\psi_1^*\psi_2\psi_2^*. \end{aligned}$$

This Lagrangian includes couplings between all three fields. (The last  $\psi_1\psi_2$  coupling term was not used in [11] as it proved unnecessary for that work.) It turns out this Lagrangian gives particles of the correct masses and size to model the lepton family (the electron, muon, and tau).

The field equations obtained by varying  $\psi^*$ ,  $\psi_1^*$ ,  $\psi_2^*$ , and  $A^u$  are

$$(D^a\psi)_{;a} - i\epsilon A_a(D^a\psi) + \sigma^2\psi - f_1(\psi_1\psi_1^*)\psi - f_2(\psi_2\psi_2^*)\psi = 0$$

$$(\partial^a\psi_1)_{;a} + \sigma^2\psi_1 - g_1(\psi_1)^2\psi_1^* + f_1(\psi\psi^*)\psi_1 + f_3(\psi_2\psi_2^*)\psi_1 = 0$$

$$(\partial^a\psi_2)_{;a} + \sigma^2\psi_2 - g_2(\psi_2)^2\psi_2^* + f_2(\psi\psi^*)\psi_2 + f_3(\psi_1\psi_1^*)\psi_2 = 0$$

$$F^{ua}{}_{;a} = -4\pi J^u$$

where

$$J^u \equiv \frac{1}{2}i\epsilon[\psi^*(D^u\psi) - \psi(D^u\psi)^*].$$

(Note, the subscripts 1 and 2 on  $\psi_1$  and  $\psi_2$  are not tensor indices; they serve only to distinguish the two neutral fields.)

Again looking at the case of static, spherical symmetry we have

$$A^0 = \varphi(r), \quad A^\alpha = 0 \quad (\alpha = 1, 2, 3)$$

$$\psi = \theta(r)e^{-i\epsilon\mu t} \quad (\text{charged field})$$

$$\psi_1 = \theta_1(r) \quad (\text{neutral field})$$

$$\psi_2 = \theta_2(r) \quad (\text{neutral field})$$

It is assumed (as before) that  $\psi_1$  and/or  $\psi_2$  is excited to produce the muon and tau states. But now that there are three fields there are different

possibilities. Using  $-$  to represent the unexcited state,  $\uparrow$  to represent the first excited state, and  $\uparrow\uparrow$  the second excited state, one formulation is to only excite the  $\psi_2$  (i.e., the  $y_2$ ) field:

	$y$	$y_1$	$y_2$
TAU:	$-$	$-$	$\uparrow\uparrow$
MUON:	$-$	$-$	$\uparrow$
ELECTRON:	$-$	$-$	$-$

But as with (18) this does not successfully model all three particles.

But by exciting alternatively  $\psi_1$  and  $\psi_2$  (i.e.,  $y_1$  and  $y_2$ ) in this scheme

	$y$	$y_1$	$y_2$
TAU:	$-$	$-$	$\uparrow$
MUON:	$-$	$\uparrow$	$-$
ELECTRON:	$-$	$-$	$-$

the correct masses are found for all three particles. But this scheme seems to suggest that the tau would naturally decay into an electron, whereas experimentally it is found to decay into both muons and electrons.

So Cooperstock and Rosen preferred the scheme

	$y$	$y_1$	$y_2$
TAU:	$-$	$\uparrow$	$\uparrow$
MUON:	$-$	$\uparrow$	$-$
ELECTRON:	$-$	$-$	$-$

where a tau can decay into a muon by de-exciting  $y_2$ , or decay into an electron by de-exciting both  $y_1$  and  $y_2$ . Just de-exciting  $y_1$  would presumably lead to a short-lived intermediate state. (However this is a purely *static* model and such questions of transitions may or may not be meaningful with the current model.)

So now

$$mass_{\text{low state}} = \frac{\sigma}{2\epsilon^2} w_{\text{low state}} \quad (22)$$

$$mass_{\text{middle state}} = \frac{\sigma}{2\epsilon^2} w_{\text{middle state}}$$

$$mass_{\text{high state}} = \frac{\sigma}{2\epsilon^2} w_{\text{high state}}$$

where (for the EMT group as we have here) the "low state" is the electron state, the "middle state" is the muon state, and the "high state" is the tau state. So  $w_{\text{low state}}$  contains integrals with the unexcited (ground-state)  $y_1$  and  $y_2$  fields, and higher-state  $w$ 's contain the excited fields. For this Lagrangian

$$\begin{aligned} w = & \left( \frac{2\pi\epsilon^2}{g_1} \right) \int_0^\infty (y_1)^4 x^2 dx + \left( \frac{2\pi\epsilon^2}{g_2} \right) \int_0^\infty (y_2)^4 x^2 dx \quad (23) \\ & - \left( \frac{f_1}{g_1} \right) \int_0^\infty y^2 (y_1)^2 x^2 dx - \left( \frac{f_2}{g_2} \right) \int_0^\infty y^2 (y_2)^2 x^2 dx \\ & - \left( \frac{4\pi\epsilon^2 f_3}{g_1 g_2} \right) \int_0^\infty (y_1)^2 (y_2)^2 x^2 dx - (\alpha\beta + \gamma) \end{aligned}$$

where

$$\alpha \equiv \int_0^\infty y^2 z x^2 dx \quad (25)$$

$$\beta \equiv \frac{\epsilon \mu}{\sigma}$$

$$\gamma \equiv \int_0^\infty y^2 z^2 x^2 dx$$

as before.

It will also be useful to introduce some abbreviations for the factors appearing in the equation for w:

$$FG1 \equiv -f_1/g_1 \quad (24)$$

$$FG2 \equiv -f_2/g_2$$

$$FE1 \equiv -f_1/4\pi\epsilon^2$$

$$FE2 \equiv -f_2/4\pi\epsilon^2$$

$$FG3 \equiv \frac{+f_1 f_3}{g_1 g_2}$$

$$I_1 \equiv \int_0^\infty y_1^4 x^2 dx / 2$$

$$I_2 \equiv \int_0^\infty y_2^4 x^2 dx / 2$$

$$I_3 \equiv \int_0^\infty y^2 y_1^2 x^2 dx$$

$$I_4 \equiv \int_0^\infty y^2 y_2^2 x^2 dx$$

$$I_5 \equiv \int_0^\infty y_1^2 y_2^2 x^2 dx$$

so that

$$w = \left(\frac{FG1}{FE1}\right) I_1 + \left(\frac{FG2}{FE2}\right) I_2 + (FG1) I_3 + (FG2) I_4 + \left(\frac{FG3}{FE1}\right) I_5 - (\alpha\beta + \gamma). \quad (25)$$

So if we choose our input parameters FG1, FG2, FG3, FE1, and FE2, and then find  $y(x)$ ,  $z(x)$ ,  $y_1(x)$ , and  $y_2(x)$  satisfying their respective boundary conditions and giving particle-like solutions, how do we work backwards to obtain the other parameters?

1)  $\alpha = \int_0^\infty y^2 z x^2 dx$  and  $\gamma = \int_0^\infty y^2 z^2 x^2 dx$  so  $\alpha$  and  $\gamma$  can both be integrated easily. (At least they can be approximated since we can only integrate a finite distance down the  $x$ -axis before the previously-mentioned instability of the  $e^x/x$  term causes the integrals to diverge. But if care is taken, the exponential decay of  $y$  and  $x$  ensures that the integrals can be approximated to enough accuracy with little trouble.)

2) We also know via (11) that

$$z \rightarrow \frac{\alpha}{x} + \beta \quad \text{as } x \rightarrow \infty$$

so once we determine  $\alpha$  we can also determine  $\beta$ .

3)  $w_{\text{low state}}$ ,  $w_{\text{middle state}}$ , and  $w_{\text{high state}}$  are calculated via (25) (given input values for FG1 etc.).

4) From (22) and (21) we have

$$\begin{aligned} \sigma &= \frac{\text{mass}_{\text{low state}}}{w_{\text{low state}}} 2\epsilon^2 \\ &= \frac{\text{mass}_{\text{low state}}}{w_{\text{low state}}} \frac{2\alpha^2}{q^2} \end{aligned}$$

so  $\sigma$  can be evaluated once  $w_{\text{low state}}$  and  $\alpha$  have been calculated (and using the known charges and masses).

5) Then  $\epsilon$  is obtained from

	JON	MOON	TAU
$w$	1.6300	1.6343	
$y_1(0)$	2.2112	2.2025	
$y_2(0)$	4.3383	4.1036	
$y_3(0)$	4.3431	4.341	
$w$	$4.286 \times 10^{-3}$	0.8889	14.974

$$\epsilon = \frac{\alpha}{q}$$

(We assume that  $\epsilon_{\text{low state}} = \epsilon_{\text{middle state}} = \epsilon_{\text{high state}}$ , and that  $q$  (the charge) is the same for all three states, so we have to adjust the parameters to keep  $\alpha_{\text{low state}} = \alpha_{\text{middle state}} = \alpha_{\text{high state}}$ .)

The next step is to adjust the initial values (i.e.,  $y(0)_{\text{low state}}$ ,  $y(0)_{\text{middle state}}$ , and  $y(0)_{\text{high state}}$ ) to try to make  $\alpha_{\text{low state}} = \alpha_{\text{middle state}} = \alpha_{\text{high state}}$  (so that charge is conserved), and to adjust the F's (i.e., FG1, FG2, FG3, FE1, and FE2) to try to make

$$\frac{w_{\text{middle state}}}{w_{\text{low state}}} = \frac{\text{mass}_{\text{middle state}}}{\text{mass}_{\text{low state}}} \tag{26}$$

$$\frac{w_{\text{high state}}}{w_{\text{low state}}} = \frac{\text{mass}_{\text{high state}}}{\text{mass}_{\text{low state}}}$$

Of course these new F's and  $y(0)$ 's will alter the solution so the whole process has to be repeated until, after many trials, everything converges. The aim is to find a set of parameters (the F's and the initial-values) so that: 1) charge is conserved, and 2) the correct mass-ratios (26) are obtained. Then  $1/\sigma$  provides a check of the theory when compared to the particle size.

Table 2 shows the results as found by Cooperstock and Rosen.

so  $e/m \sim 10^{21}$ . However, Cooperstock and Rosen found that when gravity was studied in this model, every charged particle where gravity played a significant

	ELECTRON	MUON	TAU
$y(0)$	1.6300	1.6300	1.6243
$z(0)$	2.2112	2.2112	2.2025
$y_1(0)$	4.3383	14.1038	14.1049
$y_2(0)$	4.3431	4.4341	14.1049
$w$	$4.286 \times 10^{-3}$	0.8869	14.974
$\alpha$	1.9051	1.9051	1.9051
$\beta$	0.00264	0.00264	0.00489
$\gamma$	2.8168	2.8168	2.8077
FG1	= -2.770872	$\times 10^{-05}$	
FG2	= -2.652788	$\times 10^{-03}$	
FE1	= -5.0	$\times 10^{-04}$	
FE2	= -3.0	$\times 10^{-03}$	
$\sigma$	= 6.013	$\times 10^{15}$	cm <sup>-1</sup>
$1/\sigma$	= 1.663	$\times 10^{-16}$	cm
$\epsilon$	= -1.379	$\times 10^{34}$	cm <sup>-1</sup>
$w_{\text{muon}}/w_{\text{electron}}$	= 206.9	error = -0.06%	
$w_{\text{tau}}/w_{\text{electron}}$	= 3494	error = -0.06%	

**TABLE 2:** *The solution for the EMT group as found by Cooperstock and Rosen.*

(Here  $\epsilon$ , which is related to the charge, has been expressed in dimensions of length via the use of "geometric units." See Appendix C.)

In geometric units

$$e = 6.77 \times 10^{-56} \text{ cm}$$

$$\text{mass of an electron} = 1.38 \times 10^{-34} \text{ cm}$$

so  $e/m \sim 10^{21}$ . However, Cooperstock and Rosen found that when gravity was studied in this model, every charged particle where gravity played a significant

role led to an  $e/m$  ratio  $\sim 1$ , with a radius of  $< \sim 10^{-33}$  cm. So they conclude that gravity does *not* play a significant role in lepton structure, and that the electron radius has a lower bound well above  $10^{-33}$  cm.

\* \* \*

### CHAPTER 3: NEW WORK: FITTING LEPTONS AND QUARKS

Now suppose we wish to extend the model to include quarks (which, like the leptons, are also spin 1/2 particles). They can be grouped into families paralleling the EMT (electron/muon/tau) group:

	EMT group	UCT group	DSB group
Third generation:	TAU	TOP	BOTTOM
Second generation:	MUON	CHARM	STRANGE
First generation:	ELECTRON	UP	DOWN
Charge (in units of $e$ ):	-1	+2/3	-1/3

At this point we do not know what parameters can be held constant and which should be allowed to vary. Another problem is that the masses of quarks are model-dependent and not clearly defined. Various estimates of quark masses are given in Table 3.

	Barv mass	Griffiths Effective mass in mesons	Effective mass in baryons	Veltman	Physics Letters
TOP:	>23,000	>23,000	>23,000	23,500	50,000
CHARM:	1,100	1,500	1,500	1,500	1,300
UP:	4.2	310	363	310	1.5
CHARM/UP:	261.9	4.8	4.1	4.8	241.0

### CHAPTER 3: NEW WORK: FITTING LEPTONS AND QUARKS

Now suppose we wish to extend the model to include quarks (which, like the leptons, are also spin 1/2 particles). They can be grouped into families paralleling the EMT (electron/muon/tau) group:

TABLE 3: Different values of estimates for the masses of quarks and the resulting mass ratios (in  $\text{MeV}/c^2$ ) from Griffiths [8], Veltman [13], and Physics Letters [14].

	EMT group	UCT group	DSB group
Third generation:	TAU	TOP	BOTTOM
Second generation:	MUON	CHARM	STRANGE
First generation:	ELECTRON	UP	DOWN
Charge (in units of e):	-1	+2/3	-1/3

At this point we do not know what parameters can be held constant and which should be allowed to vary. Another problem is that the masses of quarks are model-dependent and not clearly defined. Various estimates of quark masses are given in Table 3.

<sup>11</sup> Physics Letters, [14], page 14.

	Griffiths			Veltman	Physics Letters
	Bare mass	Effective mass in mesons	Effective mass in baryons		
TOP:	>23,000	>23,000	>23,000	22,500	50,000
CHARM:	1,100	1,500	1,500	1,500	1,350
UP:	4.2	310	363	310	5.6
CHARM/UP:	261.9	4.8	4.1	4.8	241.0
TOP/UP:	5,476.2	74.2	63.4	72.6	8,928.6
BOTTOM:	4,200	4,700	4,700	≈5,000	5,000
STRANGE:	150	483	538	505	199
DOWN:	7.5	310	363	310	9.9
STRANGE/DOWN:	20.0	1.6	1.5	1.6	20.1
BOTTOM/DOWN:	560.0	15.2	12.9	16.1	505.0

**TABLE 3:** Different values of estimates for the masses of quarks and the resulting mass ratios (in  $\text{MeV}/c^2$ ) from Griffiths [8], Veltman [13], and Physics Letters [14].

The masses of the leptons are<sup>11</sup>

$$\text{electron mass} = 0.51099906 \text{ MeV}/c^2$$

$$\text{muon mass} = 105.65839 \text{ MeV}/c^2$$

$$\text{tau mass} = 1784.1 \text{ MeV}/c^2$$

So to fit the quarks we need to find the correct set of coupling constants (F's) and initial conditions for each particle so that: 1) we get the correct  $\alpha$ 's, and hence the correct charges, and 2) we get the correct w ratios, and hence the correct mass ratios. Then, as an independent check of the model, we will have a prediction of the particle sizes (through  $1/\sigma$ ).

<sup>11</sup> *Physics Letters*, [14], page 14.

The first attempt made was to use Cooperstock and Rosen's values for the EMT group as in Table 2, and to try to find different  $F$ 's to fit the UCT (up/charm/top) group. This proved difficult to do as it forced the search into areas where  $\beta \simeq 1$ , which is numerically unstable, so another approach was tried.

In looking at previously unpublished data [15] of solutions to the pure Rosen equations for negative-mass particles

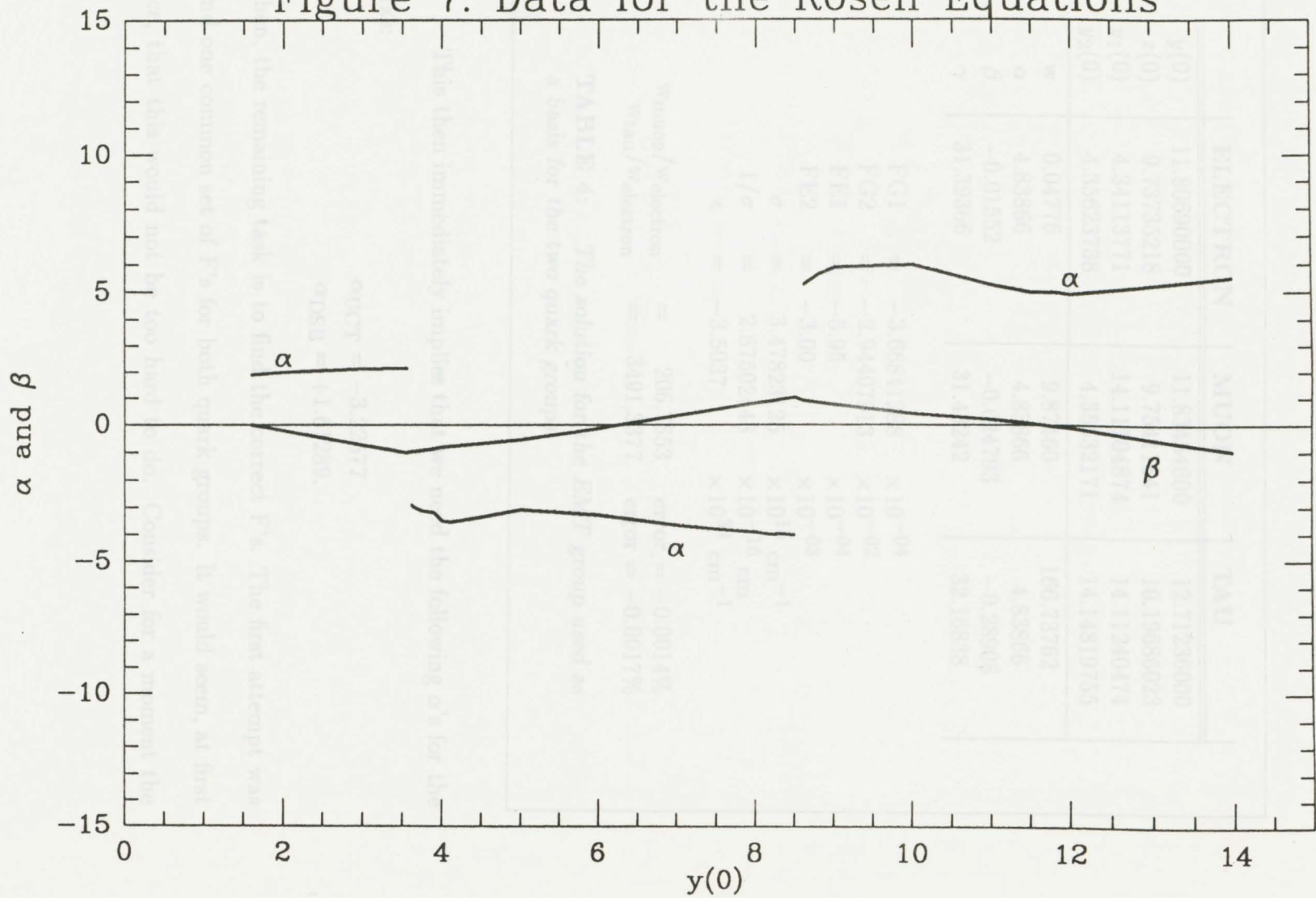
$$y'' + \frac{2}{x}y' - y + yz^2 = 0$$

$$z'' + \frac{2}{x}z' + y^2z = 0$$

it can be seen that there naturally occur three different "regimes" (see Figure 7). In particular notice how  $\alpha$  naturally changes sign in the middle regime. This sign change, because of  $\alpha$ 's connection with the charge via (9), easily accommodates the alternating charges of the three groups. (If  $\alpha$  did not change sign it could still be forced to do so by changing the sign of  $z$ , but the approach used here seems more elegant.)

We assume that solutions to the full equations are not too far removed from solutions to the simpler Rosen equations, so the first step was to fix an EMT solution in the high regime and use that as a basis for the other two quark groups. Once  $\alpha_{\text{EMT}}$  is fixed, then, because of the known charge ratios,  $\alpha_{\text{UCT}}$  and  $\alpha_{\text{DSB}}$  are thereby fixed also. This was done; the solution is shown in Table 4. In particular we have  $\alpha_{\text{EMT}} = +4.83866$ .

Figure 7: Data for the Rosen Equations



	ELECTRON	MUON	TAU
$y(0)$	11.80690000	11.83454600	12.71236000
$z(0)$	9.73735218	9.75411841	10.19686023
$y_1(0)$	4.34113771	14.11104874	14.11240474
$y_2(0)$	4.35623738	4.35632171	14.14819755
w	0.04776	9.87460	166.73762
$\alpha$	4.83866	4.83866	4.83866
$\beta$	-0.01552	-0.024793	-0.28908
$\gamma$	31.39366	31.43242	32.16838
FG1	= -3.66841288	$\times 10^{-04}$	
FG2	= -2.94407913	$\times 10^{-02}$	
FE1	= -5.95	$\times 10^{-04}$	
FE2	= -3.00	$\times 10^{-03}$	
$\sigma$	= 3.47823125	$\times 10^{15} \text{ cm}^{-1}$	
$1/\sigma$	= 2.87502448	$\times 10^{-16} \text{ cm}$	
$\epsilon$	= -3.5037	$\times 10^{34} \text{ cm}^{-1}$	
$w_{\text{muon}}/w_{\text{electron}}$	= 206.7653	error = -0.0014%	
$w_{\text{tau}}/w_{\text{electron}}$	= 3491.3377	error = -0.0017%	

**TABLE 4:** The solution for the EMT group used as a basis for the two quark groups.

This then immediately implies that we need the following  $\alpha$ 's for the quarks:

$$\alpha_{\text{UCT}} = -3.22577$$

$$\alpha_{\text{DSB}} = +1.61289.$$

So then, the remaining task is to find the correct F's. The first attempt was to find one common set of F's for both quark groups. It would seem, at first glance, that this would not be too hard to do. Consider for a moment the

integrals in (23) to be fixed. Then if we consider FG1, FG2, and FE2 to be variables and FE1 and FG3 to be fixed, we can define

$$a \equiv \text{FG1}$$

$$b \equiv \text{FG2/FE2}$$

$$c \equiv \text{FG2.}$$

Then we define

$$A \equiv \frac{I_{1, \text{ low state}}}{\text{FE1}}$$

$$B \equiv I_{2, \text{ low state}}$$

$$C \equiv I_{3, \text{ low state}}$$

$$D \equiv I_{4, \text{ low state}}$$

$$E \equiv \frac{\text{FG3}}{\text{FE1}} \cdot I_{5, \text{ low state}} - (\alpha\beta + \gamma)_{\text{low state}}$$

Then F, G, H, J, K are defined similarly for the middle state, and M, N, P, Q, and R are defined similarly for the high state so that

$$w_{\text{up}} = a \cdot A_{\text{UCT}} + b \cdot B_{\text{UCT}} + a \cdot C_{\text{UCT}} + c \cdot D_{\text{UCT}} + E_{\text{UCT}}$$

$$w_{\text{charm}} = a \cdot F_{\text{UCT}} + b \cdot G_{\text{UCT}} + a \cdot H_{\text{UCT}} + c \cdot J_{\text{UCT}} + K_{\text{UCT}}$$

$$w_{\text{down}} = a \cdot A_{\text{DSB}} + b \cdot B_{\text{DSB}} + a \cdot C_{\text{DSB}} + c \cdot D_{\text{DSB}} + E_{\text{DSB}}$$

$$w_{\text{strange}} = a \cdot F_{\text{DSB}} + b \cdot G_{\text{DSB}} + a \cdot H_{\text{DSB}} + c \cdot J_{\text{DSB}} + K_{\text{DSB}}$$

$$w_{\text{bottom}} = a \cdot M_{\text{DSB}} + b \cdot N_{\text{DSB}} + a \cdot P_{\text{DSB}} + c \cdot Q_{\text{DSB}} + S_{\text{DSB}}$$

If we define

$$R1 \equiv \frac{\text{mass}_{\text{middle state}}}{\text{mass}_{\text{low state}}} = \frac{w_{\text{middle state}}}{w_{\text{low state}}}$$

Of course in actuality the integrals are not fixed, but instead depend on the functions  $y$ ,  $z$ ,  $y_1$ , and  $z_1$ , so these calculated "perfect"  $F$ 's then have to be used to obtain new solutions for  $y$ ,  $z$ ,  $y_1$ , and  $z_1$ , and hence obtain new integrals. The hope was that this process would iteratively converge to an answer with the desired accuracy (and simultaneously an independent estimate of the mass of the top quark). However the problem is that the integrals are not independent, and the order of 1, and with such strong coupling between the equations, solutions are almost impossible to prove that such a set of  $F$ 's does not exist.)

$$R2 \equiv \frac{\text{mass}_{\text{high state}}}{\text{mass}_{\text{low state}}} = \frac{W_{\text{high state}}}{W_{\text{low state}}}$$

then we have

$$R1 \cdot w_{\text{low state}} = W_{\text{middle state}}$$

so that

$$R1 \cdot [a(A + C) + bB + cD + E] = a(F + H) + bG + cJ + K$$

for each family. Therefore

$$a[R1(A + C) - (F + H)] + b[R1 \cdot B - G] + c[R1 \cdot D - J] = [K - R1 \cdot E].$$

Similarly

$$a[R2(A + C) - (M + P)] + b[R2 \cdot B - N] + c[R2 \cdot D - Q] = [S - R2 \cdot E].$$

If we consider the mass of the top quark unknown then we are led to the matrix equation

$$\begin{bmatrix} R1_{\text{UCT}} \cdot (A_{\text{UCT}} + C_{\text{UCT}}) - (F_{\text{UCT}} + H_{\text{UCT}}) & R1_{\text{UCT}} \cdot B_{\text{UCT}} - G_{\text{UCT}} & R1_{\text{UCT}} \cdot D_{\text{UCT}} - J_{\text{UCT}} \\ R1_{\text{DSB}} \cdot (A_{\text{DSB}} + C_{\text{DSB}}) - (F_{\text{DSB}} + H_{\text{DSB}}) & R1_{\text{DSB}} \cdot B_{\text{DSB}} - G_{\text{DSB}} & R1_{\text{DSB}} \cdot D_{\text{DSB}} - J_{\text{DSB}} \\ R2_{\text{DSB}} \cdot (A_{\text{DSB}} + C_{\text{DSB}}) - (M_{\text{DSB}} + P_{\text{DSB}}) & R2_{\text{DSB}} \cdot B_{\text{DSB}} - N_{\text{DSB}} & R2_{\text{DSB}} \cdot D_{\text{DSB}} - Q_{\text{DSB}} \end{bmatrix}$$

and hence generate five methods, and others, producing these various matrix

$$\times \begin{bmatrix} a \\ b \\ c \end{bmatrix} = \begin{bmatrix} K_{\text{UCT}} - R1_{\text{UCT}} \cdot E_{\text{UCT}} \\ K_{\text{DSB}} - R1_{\text{DSB}} \cdot E_{\text{DSB}} \\ S_{\text{DSB}} - R2_{\text{DSB}} \cdot E_{\text{DSB}} \end{bmatrix}.$$

A slightly different approach was also tried. A parameter-fitting program is written which takes the integrals as input, and then iteratively solves for the "perfect"  $F$ 's that, with the given integrals, will give the desired mass ratios: bottom/down, strange/down, and charm/up.

Of course in actuality the integrals are *not* fixed, but instead depend on the functions  $y$ ,  $z$ ,  $y_1$ , and  $y_2$ , which in turn depend on the F's, so these calculated "perfect" F's then have to be used to obtain new solutions for  $y$ ,  $z$ ,  $y_1$ , and  $y_2$ , and hence obtain new integrals. The hope was that this process would iteratively converge to an answer with the desired accuracy (and simultaneously provide an independent estimate of the mass of the top quark). However the problem arises that the calculated "perfect" F's are very large, on the order of 1, and with such strong coupling between the equations, solutions are almost impossible to obtain. The other possible excitation schemes were tried as well, but they also require similar strong couplings. (Unfortunately it seems impossible to prove that such a set of F's does not exist.)

Various other schemes were tried too. Using the two mass ratios ( $w_{\text{middle state}}/w_{\text{low state}}$ , and  $w_{\text{high state}}/w_{\text{low state}}$ ) for the EMT, UCT, and DSB groups yields six equations in five unknowns (the five F's), so the overdetermined system can be solved if one equation is discarded. Or, putting  $FG3=0$  and fitting only the UCT and DSB groups gives four equations in four unknowns. Another approach is to specify (arbitrarily)  $w_{\text{up}}$  and  $w_{\text{down}}$  and then use the known mass ratios to calculate  $w_{\text{charm}}$ ,  $w_{\text{strange}}$ , and  $w_{\text{bottom}}$ , and hence generate five equations in five unknowns. All these various matrix methods, and others, proved fruitless.

A slightly different approach was also tried. A parameter-fitting program<sup>12</sup> was used to try to adjust each parameter in turn to try to converge

<sup>12</sup> Listings and full details of the programs used are in Appendices D and E.

to the desired solution. This has two advantages: it does not require an equal number of equations and unknowns and so can deal effectively with both under-determined and over-determined systems. Also, it does not seek an exact solution only a solution which minimizes the error. Thus one could hope that it might be able to iteratively converge to a solution where the matrix methods fail. But both it and another parameter-fitting program that was written failed to find a common set of  $F$ 's that would fit even just the two quark families.

Such a task does not seem, at first glance, to be so hard; why does it fail? One reason is that when fitting just *one* group the  $F$ 's can be adjusted so that  $w_{\text{low state}}$  is small. Then minor variations in  $w_{\text{middle state}}$  and  $w_{\text{high state}}$  produce large variations in the ratios  $w_{\text{middle state}}/w_{\text{low state}}$  and  $w_{\text{high state}}/w_{\text{low state}}$ . So if  $w_{\text{low state}}$  is small, only minor variations in  $F$ 's can fit the two mass ratios for that group very closely. However the situation is completely different as soon as it is attempted to fit two groups with the same  $F$ 's. Now a typical situation is that the  $F$ 's that make  $w_{\text{low state}}$  small for one group also make  $w_{\text{low state}} \approx 10$  for the other group. And so much larger changes in  $w_{\text{middle state}}$  and  $w_{\text{high state}}$  are needed to fit the two ratios for that group. It is difficulties like these that stop the iterative procedure from working.

But can we achieve something less restrictive? After all, we do not know whether there is any necessary relation between the  $F$ 's of one group and another. So suppose we attempt to find solutions to all three groups where

two of the F's (take FE1 and FE2 for convenience) are the same for all three groups, and FG1 and FG2 vary from group to group. (FG3 is put equal to zero since it is not needed now to achieve the desired mass ratios.) Such solutions can be found. They are obviously not unique since they depend on: 1) the decision of *which* two F's to hold constant, 2) the values for those constant F's, and 3) the value of  $y(0)$ , and hence  $\alpha$ , used for the basis reference group (in this case the EMT group). One such solution is given in Table 5.

This solution gives the small mass ratios of the quarks' effective masses in baryons (as in Table 3). However, with the freedom of two F's to vary for each group, we can adjust the mass ratios to fit either small or large ratios. So another solution is given in Table 6 which fits the large mass ratios of the quarks' bare masses (also as in Table 3).

The graphs of the basic EMT solution is shown in Figures 8 through 10, and the small-ratio UCT and DSB solutions are shown in Figures 11 through 16. (The large-ratio UCT and DSB solutions look very similar.)

Both solutions yield a particle size (through  $1/\sigma$ ) on the order of  $10^{-17}$  or  $10^{-18}$  cm.

So this work shows clearly that this classical field-theory approach is useful for both leptons *and* quarks.

\* \* \*

TABLE 5: Solutions for the UCT and DSB groups which achieve the small mass ratios of the quarks' effective masses in baryons.

	UP	CHARM	TOP
$y(0)$	5.95643900	5.95954400	6.24493500
$z(0)$	5.65725624	5.65907746	5.82410012
$y_1(0)$	4.33858952	14.10955750	14.10987056
$y_2(0)$	4.34342337	4.34342533	14.13530840
w	1.23183	5.090370	78.07905
$\alpha$	-3.22577	-3.22577	-3.22577
$\beta$	-0.15296	-0.15213	-0.07695
$\gamma$	12.76646	12.76822	12.91677
FG1	= -1.44096790	$\times 10^{-04}$	
FG2	= -1.37288390	$\times 10^{-02}$	
FE1	= -5.95	$\times 10^{-04}$	
FE2	= -3.00	$\times 10^{-03}$	
$\sigma$	= 9.57933143	$\times 10^{16} \text{ cm}^{-1}$	
$1/\sigma$	= 1.04391418	$\times 10^{-17} \text{ cm}$	
$\epsilon$	= -3.5037	$\times 10^{34} \text{ cm}^{-1}$	
$w_{\text{charm}}/w_{\text{up}}$	= 4.1324	error = 0.0031%	
$w_{\text{top}}/w_{\text{up}}$	= 63.3845	error = 0.0373%	

	DOWN	STRANGE	BOTTOM
$y(0)$	0.91072000	0.91079000	0.91860000
$z(0)$	1.64385770	1.64386984	1.64485315
$y_1(0)$	4.33796458	14.10395660	14.10395619
$y_2(0)$	4.34029683	4.34029705	14.10545745
w	2.13685	3.16700	27.66712
$\alpha$	1.61289	1.61289	1.61289
$\beta$	0.41928	0.41924	0.41474
$\gamma$	2.00596	2.00596	2.005663
FG1	= -3.84780260	$\times 10^{-05}$	
FG2	= -4.61288970	$\times 10^{-03}$	
FE1	= -5.95	$\times 10^{-04}$	
FE2	= -3.00	$\times 10^{-03}$	
$\sigma$	= 5.52225788	$\times 10^{16} \text{ cm}^{-1}$	
$1/\sigma$	= 1.81085355	$\times 10^{-17} \text{ cm}$	
$\epsilon$	= -3.5037	$\times 10^{34} \text{ cm}^{-1}$	
$w_{\text{strange}}/w_{\text{down}}$	= 1.4821	error = 0.0001%	
$w_{\text{bottom}}/w_{\text{down}}$	= 12.9476	error = 0.0001%	

**TABLE 5:** Solutions for the UCT and DSB groups which achieve the small mass ratios of the quarks' effective masses in baryons.

	UP	CHARM	TOP
$y(0)$	5.94912000	5.95182000	6.21557000
$z(0)$	5.65519001	5.65677482	5.80965891
$y_1(0)$	4.33858466	14.10954358	14.10983931
$y_2(0)$	4.34339885	4.34340053	14.13515024
$w$	0.01286	3.36865	70.44521
$\alpha$	-3.22577	-3.22577	-3.22577
$\beta$	-0.15585	-0.15513	-0.08555
$\gamma$	12.76655	12.76808	12.90706
FG1	=	-1.25321400	$\times 10^{-04}$
FG2	=	-1.26169700	$\times 10^{-02}$
FE1	=	-5.95	$\times 10^{-04}$
FE2	=	-3.00	$\times 10^{-03}$
$\sigma$	=	1.06104654	$\times 10^{17} \text{ cm}^{-1}$
$1/\sigma$	=	9.42465733	$\times 10^{-18} \text{ cm}$
$\epsilon$	=	-3.5037	$\times 10^{34} \text{ cm}^{-1}$
$w_{\text{charm}}/w_{\text{up}}$	=	261.7944	error = -0.0421%
$w_{\text{top}}/w_{\text{up}}$	=	5474.6512	error = -0.0281%

	DOWN	STRANGE	BOTTOM
$y(0)$	0.90281000	0.90284900	0.90736000
$z(0)$	1.63971508	1.63972539	1.64037538
$y_1(0)$	4.33795939	14.10395815	14.10395790
$y_2(0)$	4.34027067	4.34027080	14.10546611
$w$	0.02568	0.51355	14.38032
$\alpha$	1.61289	1.61289	1.61289
$\beta$	0.42185	0.42183	0.41925
$\gamma$	2.00349	2.00350	2.00340
FG1	=	-1.82228970	$\times 10^{-05}$
FG2	=	-2.61084380	$\times 10^{-03}$
FE1	=	-5.95	$\times 10^{-04}$
FE2	=	-3.00	$\times 10^{-03}$
$\sigma$	=	9.49548009	$\times 10^{16} \text{ cm}^{-1}$
$1/\sigma$	=	1.05313264	$\times 10^{-17} \text{ cm}$
$\epsilon$	=	-3.5037	$\times 10^{34} \text{ cm}^{-1}$
$w_{\text{strange}}/w_{\text{down}}$	=	20.0012	error = 0.0058%
$w_{\text{bottom}}/w_{\text{down}}$	=	560.0705	error = 0.0126%

**TABLE 6:** Solutions for the UCT and DSB groups which achieve the large mass ratios of the quarks' bare masses.

Figure 8: The Basic Electron Solution

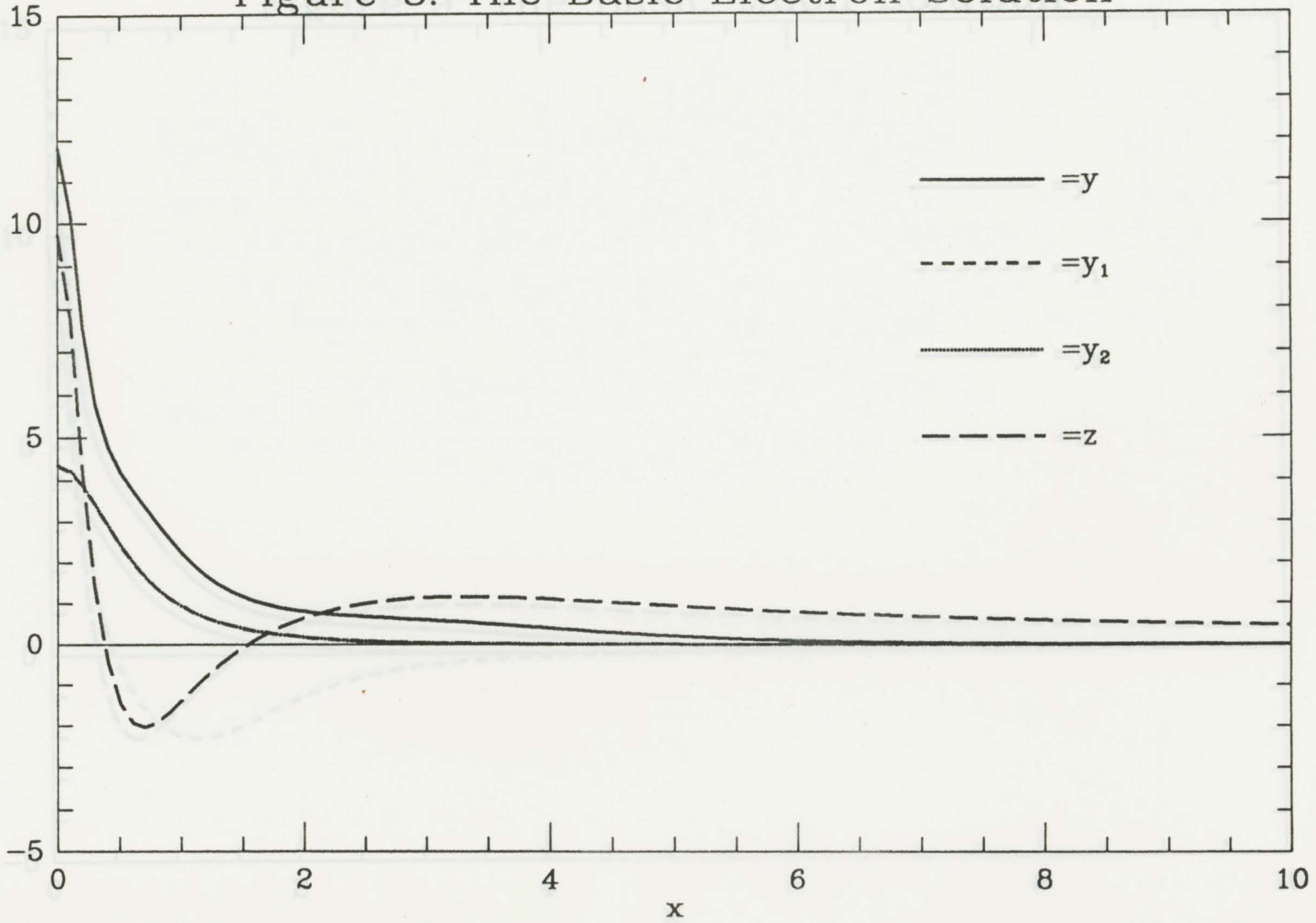


Figure 9: The Basic Muon Solution

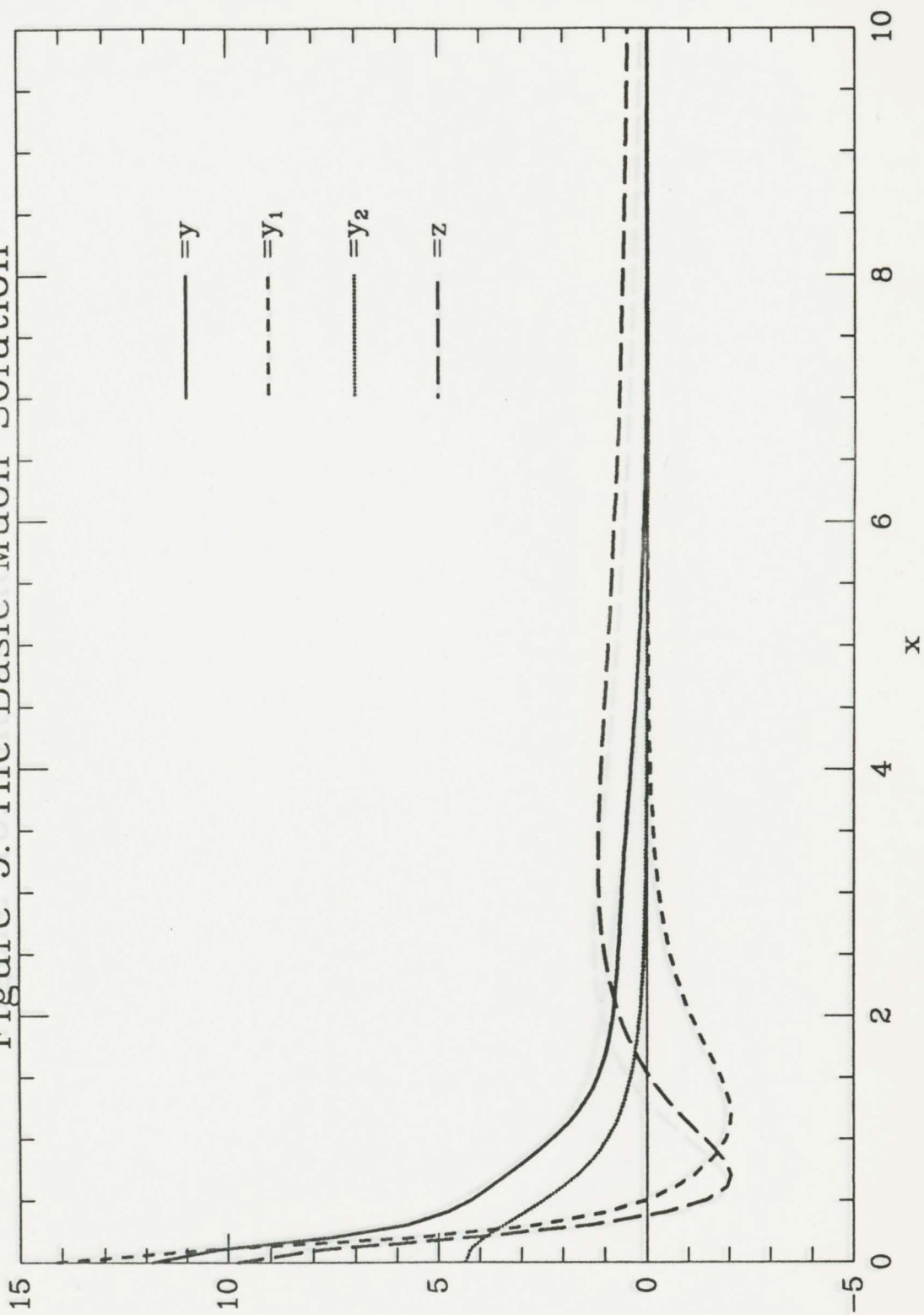


Figure 10: The Basic Tau Solution

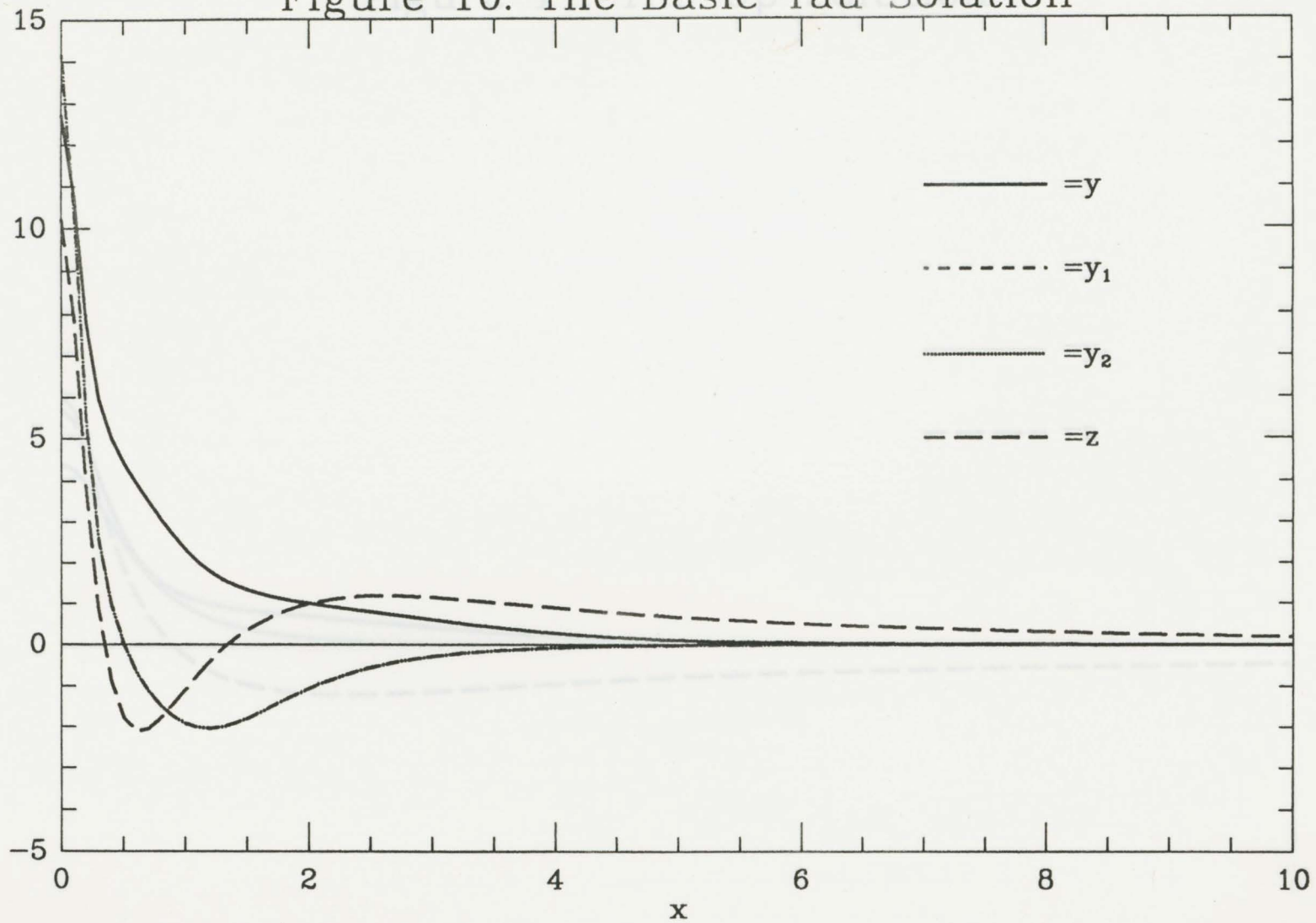


Figure 11: The Up Solution

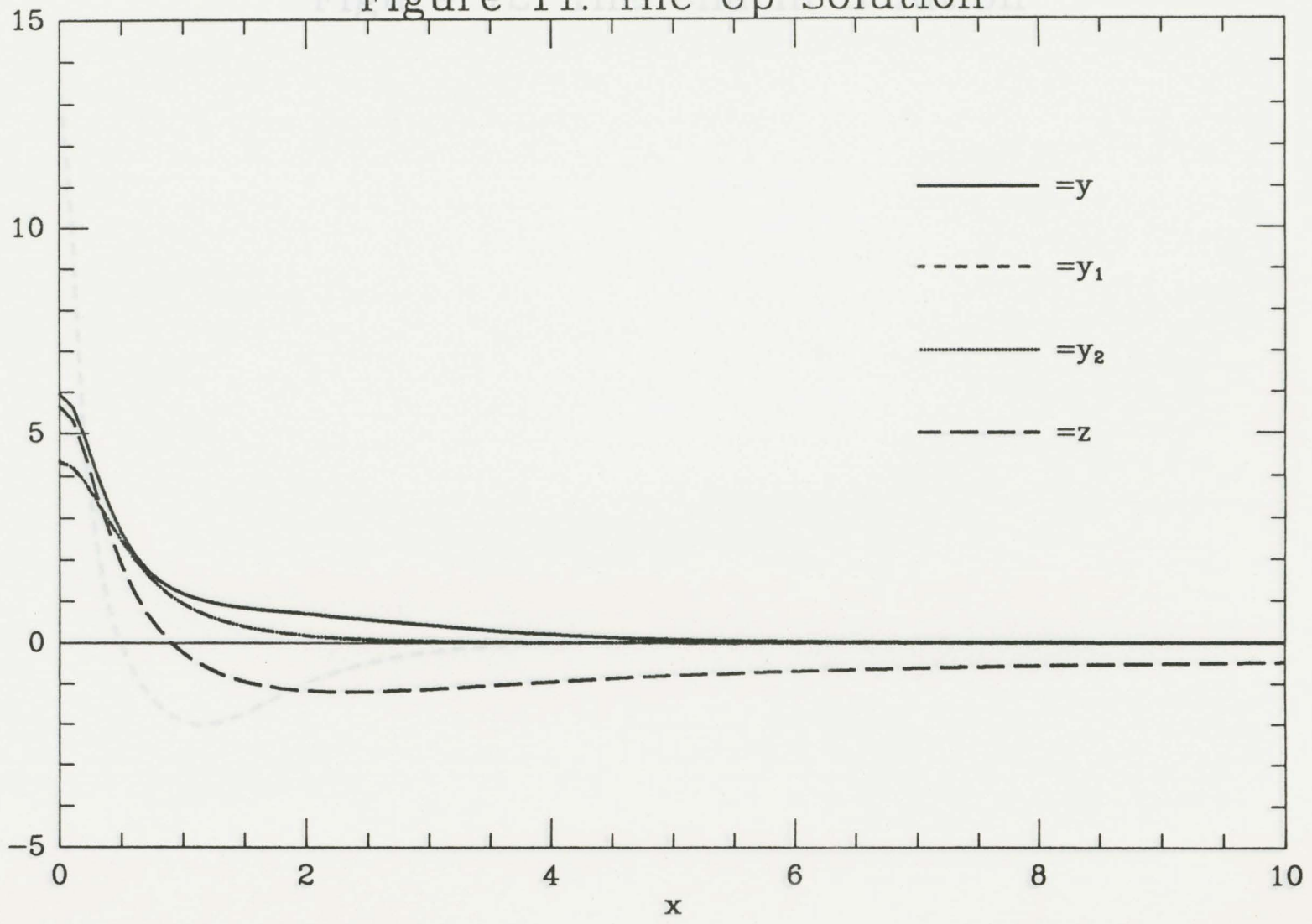


Figure 12: The Charm Solution

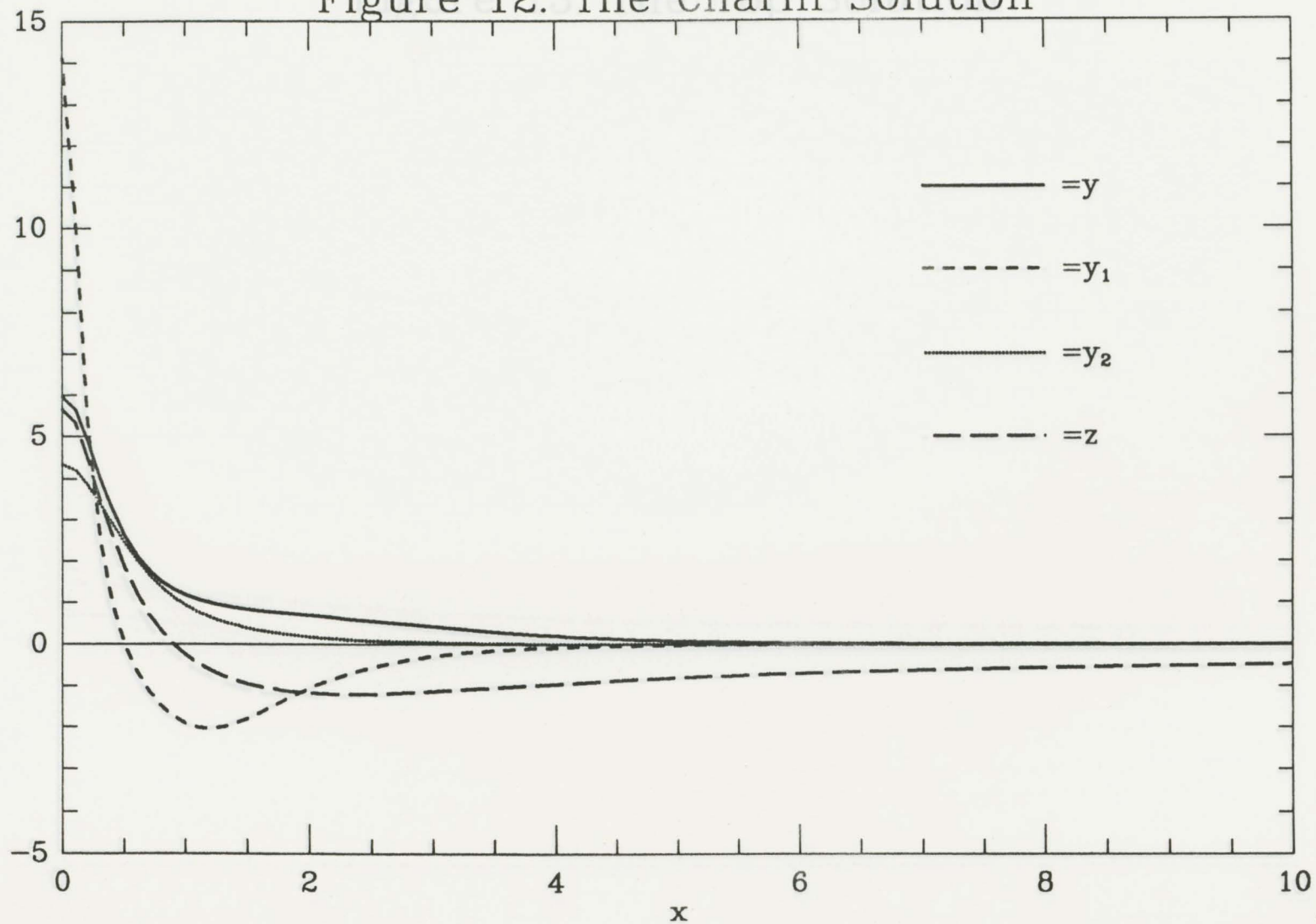


Figure 13: The Top Solution

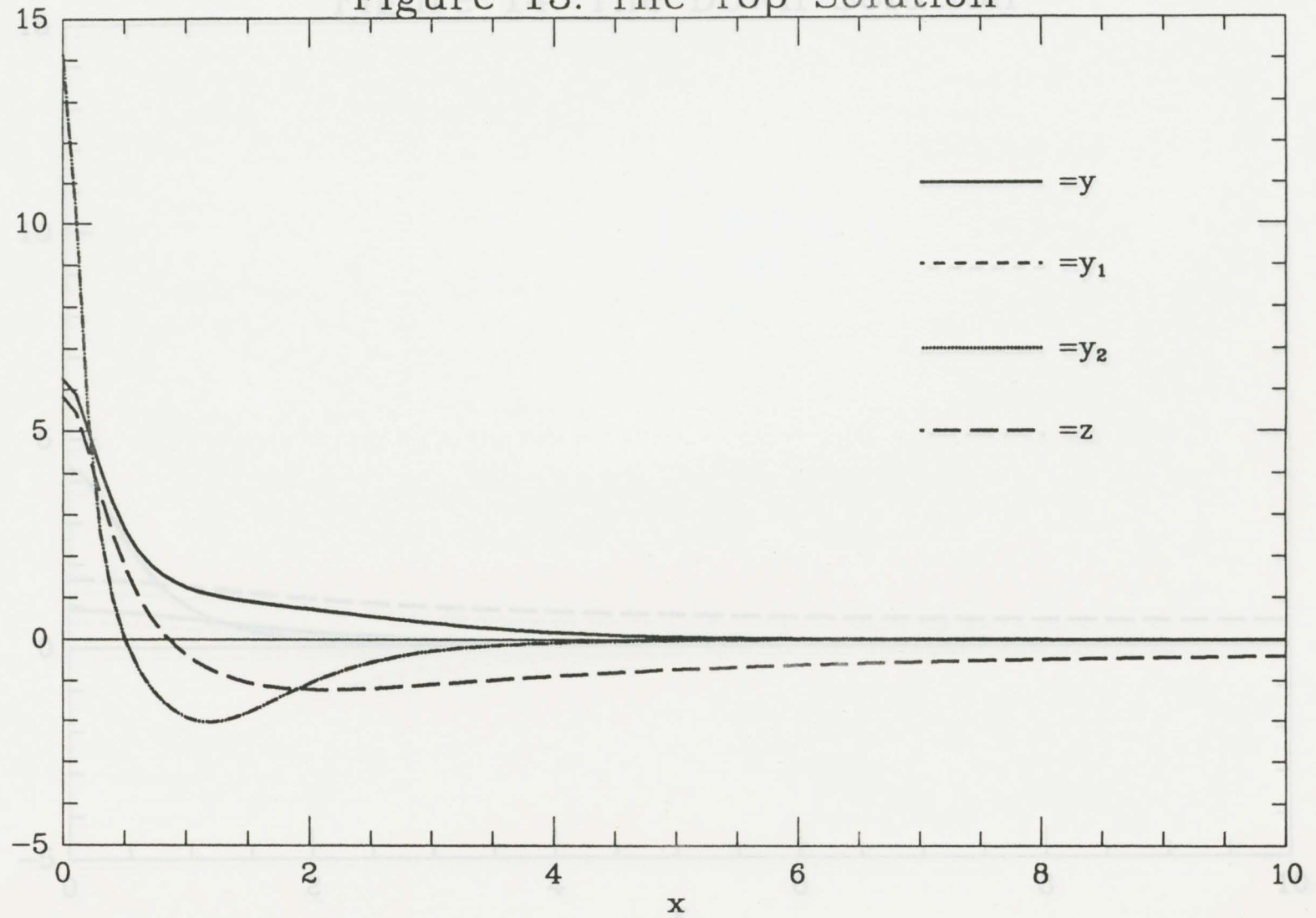


Figure 14: The Down Solution

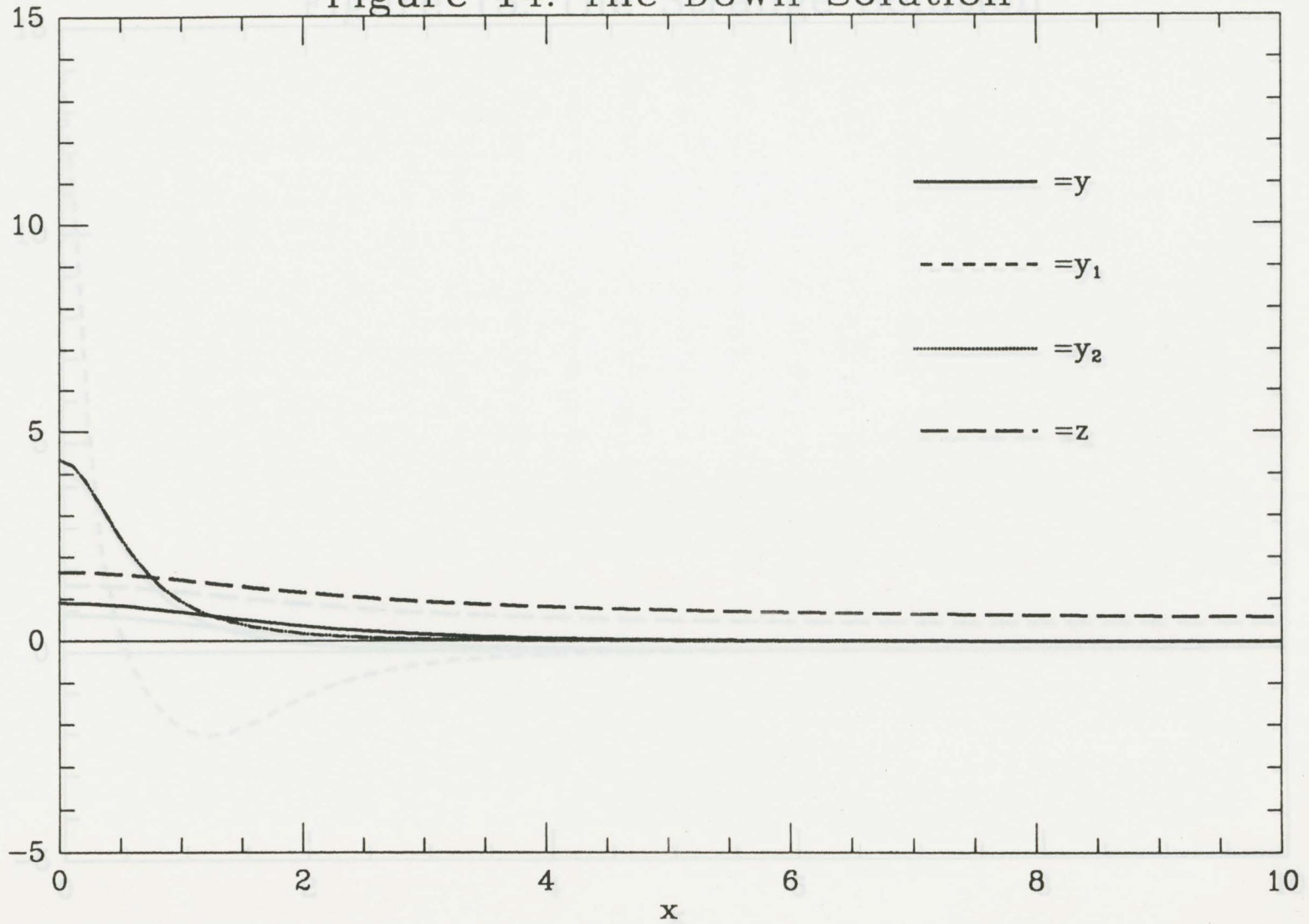


Figure 15: The Strange Solution

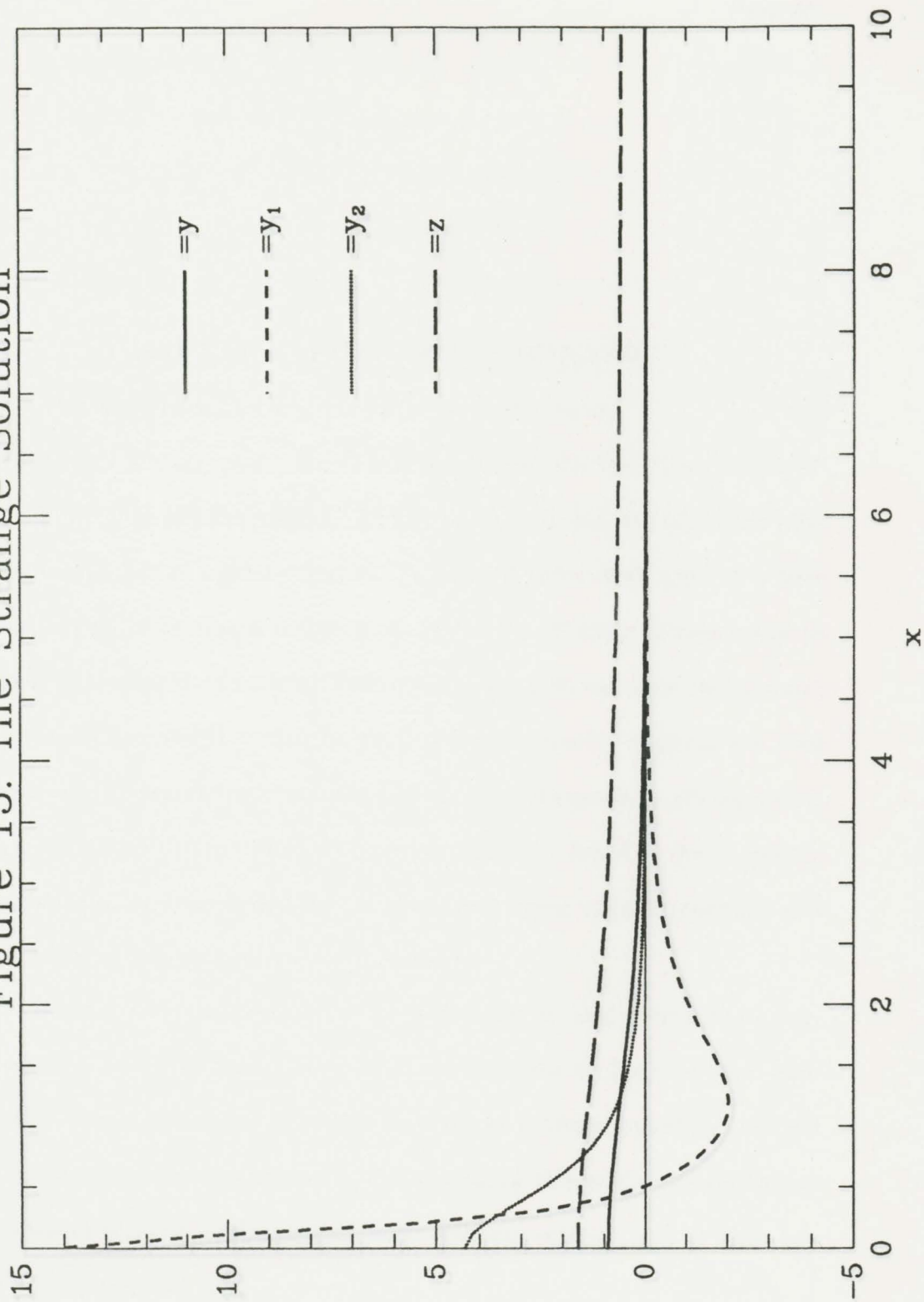
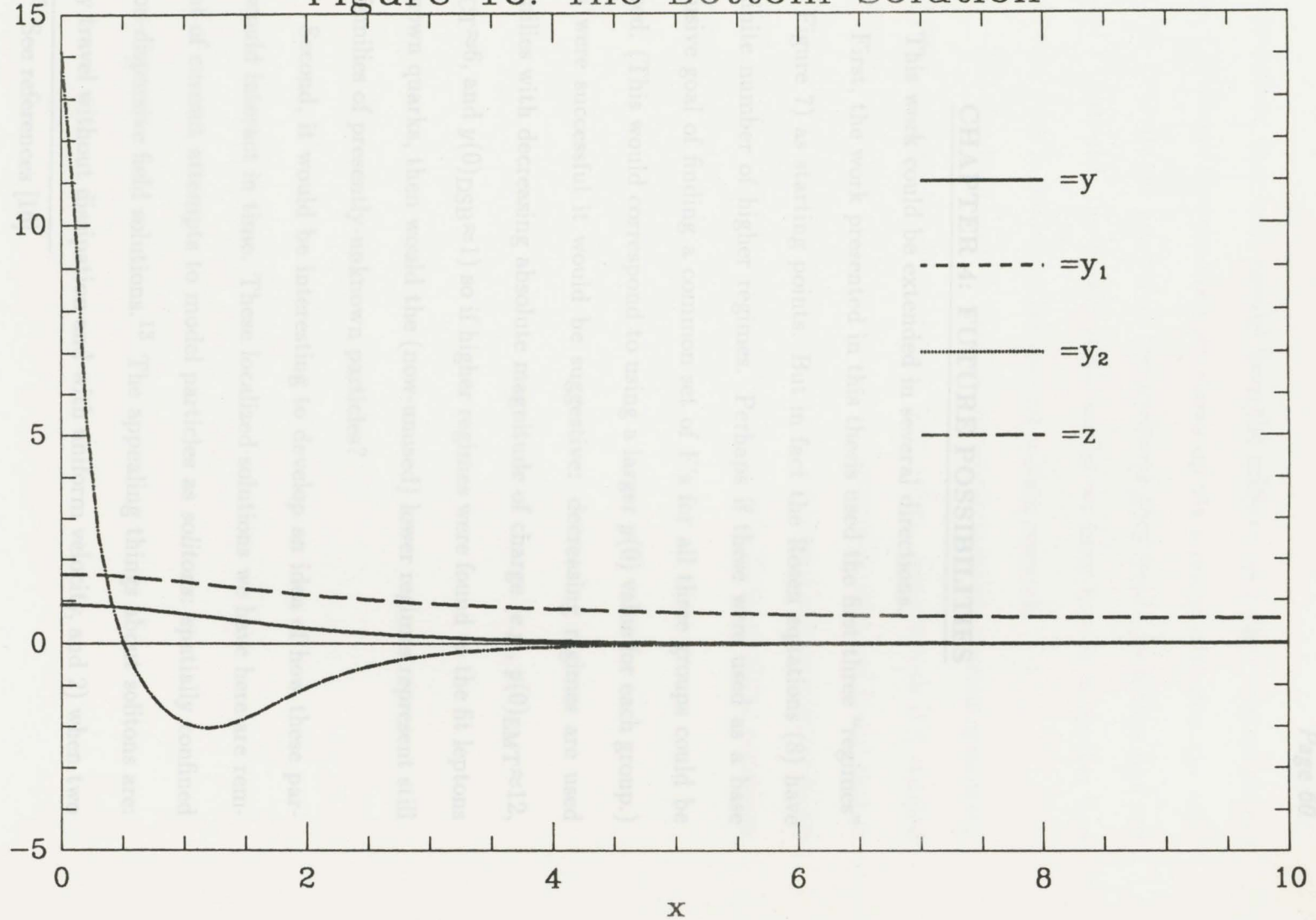


Figure 16: The Bottom Solution



#### CHAPTER 4: FUTURE POSSIBILITIES

This work could be extended in several directions.

First, the work presented in this thesis used the first three “regimes” (as in Figure 7) as starting points. But in fact the Rosen equations (8) have an infinite number of higher regimes. Perhaps if these were used as a base the elusive goal of finding a common set of  $F$ 's for all three groups could be achieved. (This would correspond to using a larger  $y(0)$  value for each group.) If this were successful it would be suggestive: decreasing regimes are used for families with decreasing absolute magnitude of charge (e.g.,  $y(0)_{\text{EMT}} \approx 12$ ,  $y(0)_{\text{UCT}} \approx 6$ , and  $y(0)_{\text{DSB}} \approx 1$ ) so if higher regimes were found to fit leptons and known quarks, then would the (now-unused) lower regimes represent still more families of presently-unknown particles?

Second, it would be interesting to develop an idea of how these particles would interact in time. These localized solutions we have here are reminiscent of current attempts to model particles as *solitons*: spatially confined and non-dispersive field solutions.<sup>13</sup> The appealing things about solitons are: 1) they travel without dissipation and with uniform velocity, and 2) when two

---

<sup>13</sup> See references [16].

solitons collide, they each emerge from the collision with their original shapes intact. (Although the non-linearity shows up via a phase-shift after the collision; the two solitons are not in the positions they would be if they had not collided.) But with the purely *static* model we have here it is unclear how closely this work is related to the current soliton research.

The third approach is to relax the requirement of spherical symmetry. Some steps towards this have already been taken. Cooperstock [17] studied the simpler Lagrangian (17) for the axially-symmetric case and found that it led to a quantized magnetic moment and spin for the particle, in addition to its mass and charge. The complete numerical solutions of the resulting equations are presently unknown.

[1] A. Einstein and N. Rosen, "The Particle Problem in the General Theory of Relativity," *Physical Review*, Volume 48, July 1, 1935.

\* \* \*

[2] G. Mie, *Annalen der Physik*, Volume 37, 1912, p. 511.

G. Mie, *Annalen der Physik*, Volume 39, 1912, p. 1.

G. Mie, *Annalen der Physik*, Volume 40, 1913, p. 10.

[3] M. Born and L. Infeld, "Foundations of the New Field Theory," *Proceedings of the Royal Society*, Volume A144, 1934, pp. 425-451.

M. Born and L. Infeld, "On the Quantization of the New Field Equations I," *Proceedings of the Royal Society*, Volume A147, 1934, pp. 522-546.

M. Born and L. Infeld, "On the Quantization of the New Field Equations II," *Proceedings of the Royal Society*, Volume A150, 1935, pp. 141-166.

[4] N. Rosen, "A Field Theory of Elementary Particles," *Physical Review*, Volume 55, January 1, 1939, pp. 94-101.

[5] David Griffiths, *Introduction To Elementary Particles*, (Harper and Row, New York: 1987).

### BIBLIOGRAPHY

- [1] S. N. Das, "Rishon Constituents of Higher-Generation Leptons and Global Conservation Law," *International Journal of Theoretical Physics*, Volume 25, No. 3, 1986, pp. 247-253.
- Michael A. Shupe, "A Composite Model of Leptons and Quarks," *Physics Letters*, Volume 86B, No. 1, 10 September 1979, pp. 87-92.
- Haim Harari, "A Schematic Model of Quarks and Leptons," *Physics Letters*, Volume 86B, No. 1, 10 September 1979, pp. 83-86.
- Haim Harari, "The Structure of Leptons and Quarks," *Scientific American*, Volume 248, Number 4, April 1983, pp. 56-68.
- [2] D. K. Sen, *Fields and/or Particles*, (Ryerson Press, Toronto: 1968).
- [3] L. Pearce Williams, *The Origins of Field Theory*, (Random House, New York: 1966).
- [4] A. Einstein and N. Rosen, "The Particle Problem in the General Theory of Relativity," *Physical Review*, Volume 48, July 1, 1935.
- [5] G. Mie, *Annalen der Physik*, Volume 37, 1912, p. 511.
- G. Mie, *Annalen der Physik*, Volume 39, 1912, p. 1.
- G. Mie, *Annalen der Physik*, Volume 40, 1913, p. 40.
- [6] M. Born and L. Infeld, "Foundations of the New Field Theory," *Proceedings of the Royal Society*, Volume A144, 1934, pp. 425-451.
- M. Born and L. Infeld, "On the Quantization of the New Field Equations I," *Proceedings of the Royal Society*, Volume A147, 1934, pp. 522-546.
- M. Born and L. Infeld, "On the Quantization of the New Field Equations II," *Proceedings of the Royal Society*, Volume A150, 1935, pp. 141-166.
- [7] N. Rosen, "A Field Theory of Elementary Particles," *Physical Review*, Volume 55, January 1, 1939, pp. 94-101.
- [8] David Griffiths, *Introduction To Elementary Particles*, (Harper and Row, New York: 1987).

- [9] N. Rosen and H. B. Rosenstock, "The Force Between Particles in a Nonlinear Field Theory," *Physical Review*, Volume 85, January 15, 1952, pp. 257-259.
- [10] R. Finkelstein, R. LeLevier, and M. Ruderman, "Nonlinear Spinor Fields," *Physical Review*, Volume 83, No. 2, July 15, 1951, pp. 326-332.
- [11] F. I. Cooperstock and N. Rosen, "A Nonlinear Gauge-Invariant Field Theory of Leptons," *International Journal of Theoretical Physics*, Volume 28, Number 4, 1989.
- [12] Chris Quigg, *Gauge Theories of the Strong, Weak, and Electromagnetic Interactions*, (Benjamin/Cummings, Massachusetts: 1983).
- [13] Martinus J. G. Veltman, "The Higgs Boson," *Scientific American*, Volume 252, Number 11, November 1986, pp. 76-105.
- [14] "Summary Tables of Particle Properties," *Physics Letters B*, Volume 204, 14 April 1988.
- [15] F. I. Cooperstock, unpublished data.
- [16] T. D. Lee, *Particle Physics and Introduction to Field Theory*, (Harwood; Chur, Switzerland: 1981), pp. 117-160.  
R. Rajaraman, *Solitons and Instantons*, (North-Holland, Amsterdam: 1982).  
P. G. Drazin and R. S. Johnson, *Solitons: An Introduction*, (Cambridge University Press, Cambridge: 1988).
- [17] F. I. Cooperstock, "Structure of Elementary Particles in a Non-Linear Field Theory," to appear in *Foundations of Physics Letters*.
- [18] *CERN Program Library* manual, 1987.03.01.
- [19] F. R. Ruckdeschel, *BASIC Scientific Subroutines, Volume II*, (McGraw-Hill, New York: 1981).

APPENDIX A: CONVENTIONS USED

The conventions used here are common ones:

1) Roman subscripts are taken to run 0, 1, 2, 3. Greek subscripts are taken to run 1, 2, 3. Also, the summation convention is used where repeated indices are assumed to have an implied summation.

2)  $f_{,k}$  and  $\partial_k f$  both denote

$$\frac{\partial f}{\partial x^k}$$

and  $f_{;k}$  denotes the covariant derivative of  $f$ .

3) Gaussian electromagnetic units are used.

4) Factors of  $c$  have generally been omitted.

5)  $\mathbf{T} = (x, \vec{y})$  indicates that the components of the four-vector  $\mathbf{T}$  are:

$$T^0 = x, \quad T^1 = y^1$$

$$T^2 = y^2$$

$$T^3 = y^3.$$

6) A metric with a  $[+ \ - \ - \ -]$  signature is used.

\* \* \*

APPENDIX B: THE DIMENSIONS OF VARIOUS SYMBOLS

Quantity:	Dimensions:
$r, \left(\frac{1}{\sigma}\right)$	length
$L$	energy/volume
$\psi, \theta, \varphi, \mu, \left(\frac{\sigma}{\epsilon}\right)$	$\sqrt{\text{energy/length}}$ = charge/distance
$\epsilon$	$\frac{1}{\sqrt{\text{energy} \cdot \text{length}}} = \frac{1}{\text{charge}}$
$f_1, f_2, f_3, g_1, g_2$	$\frac{1}{\text{energy} \cdot \text{length}}$
$x, y, z, y_1, y_2$ FG1, FG2, FG3, FE1, FE2 $w, \alpha, \beta, \gamma$	(dimensionless)

\* \* \*

## APPENDIX C: GEOMETRIC UNITS

It is convenient to use so-called "geometric units" where  $G$  (the gravitational constant) and  $c$  (the speed of light) are used to convert various quantities from conventional dimensions to dimensions of length. So:

1)

$$\text{mass}_{\text{geometric}} = \left( \frac{G}{c^2} \right) \cdot \text{mass}_{\text{conventional}}$$

therefore

$$7.428 \times 10^{-29} \text{ cm} = 1 \text{ g.}$$

2)

$$\text{charge}_{\text{geometric}} = \left( \sqrt{\frac{G}{c^4}} \right) \cdot \text{charge}_{\text{conventional}}$$

(in Gaussian units). Therefore

$$1.381 \times 10^{-34} \text{ cm} = e \quad (\text{the proton charge})$$

3)

$$\text{energy}_{\text{geometric}} = \left( \frac{G}{c^2} \right) \cdot \text{energy}_{\text{conventional}}$$

therefore

$$1.324 \times 10^{-61} \text{ cm} = 1 \text{ eV.}$$

\* \* \*

## APPENDIX D: COMPUTER PROGRAM NOTES

This appendix details some of the computer programs and procedures used for this work. All the programs are written in FORTRAN and were run on a VAX minicomputer. The only non-standard FORTRAN features used are: 1) the `/extend_source` parameter on the VAX FORTRAN compiler is used to make the compiler read input from columns 7 to 132 (instead of stopping at column 72 as in some FORTRAN versions), and 2) the inline-comment feature of the VAX FORTRAN compiler is occasionally used; the rest of the line following an exclamation mark is ignored by the compiler.

I found it convenient to make the programs read, write, and transfer data-files between themselves, thereby reducing the amount of information that must be typed in by hand. Details of the data-file structure are given below, and a diagram of the data-flow is given in Figure 17. Listings of the programs are in Appendix E.

I found it useful to keep a semi-permanent set of input data files (with different files for each family) as well as a set of temporary, generic data files that are the direct input for any particular program. For example, I have three data files `FCONEMT.DAT`, `FCONUCT.DAT`, and `FCONDSB.DAT` which contain the current F's for each group. The program `GETBCD` needs to read this information and needs the correct file depending on which group it is currently doing. Rather than have to tell `GETBCD` which file to read, I have `GETBCD` simply read the generic file `FCONSTANTS.DAT`. It is then a simple matter to copy the desired data file (e.g., `FCONEMT.DAT`) to `FCONSTANTS.DAT` before running

Figure 17: The data flow between various programs.

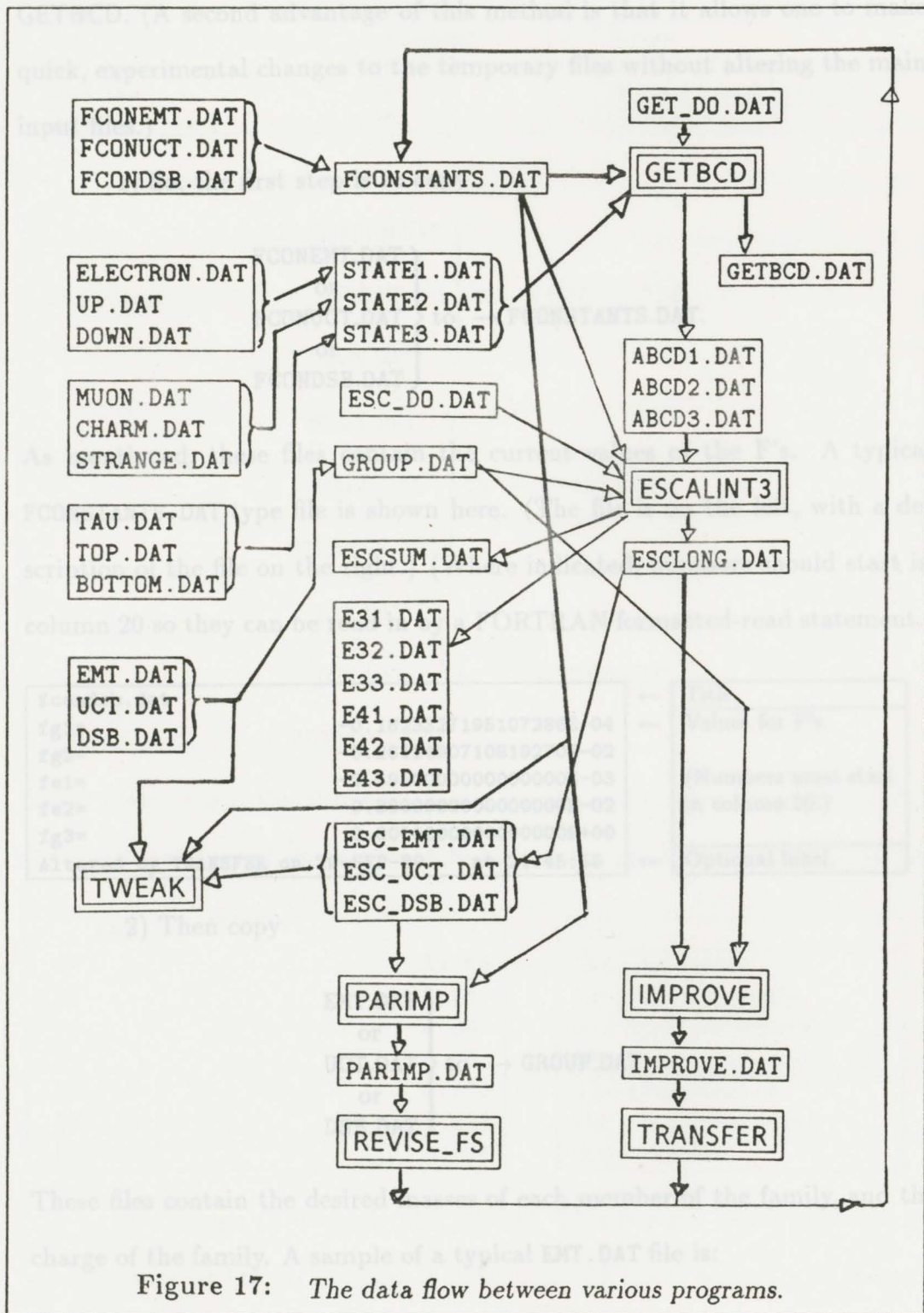


Figure 17: The data flow between various programs.

GETBCD. (A second advantage of this method is that it allows one to make quick, experimental changes to the temporary files without altering the main input files.)

1) So, the first step is to copy

$$\left. \begin{array}{l} \text{FCONEMT.DAT} \\ \text{or} \\ \text{FCONUCT.DAT} \\ \text{or} \\ \text{FCONDSB.DAT} \end{array} \right\} \text{to } \rightarrow \text{FCONSTANTS.DAT.}$$

As mentioned, these files contain the current values of the F's. A typical FCONSTANTS.DAT-type file is shown here. (The file is on the left, with a description of the file on the right.) (Where indicated, numbers should start in column 20 so they can be read in by a FORTRAN formatted-read statement.)

fcondsb.dat		←	Title.
fg1=	-0.1823327195107286E-04	←	Values for F's.
fg2=	-0.2612490710819270E-02		
fe1=	-0.5950000000000000E-03		(Numbers must start in column 20.)
fe2=	-0.3000000000000000E-02		
fg3=	0.0000000000000000E+00		
Altered by TRANSFER on 29-SEP-89	at 21:45:55	←	Optional label.

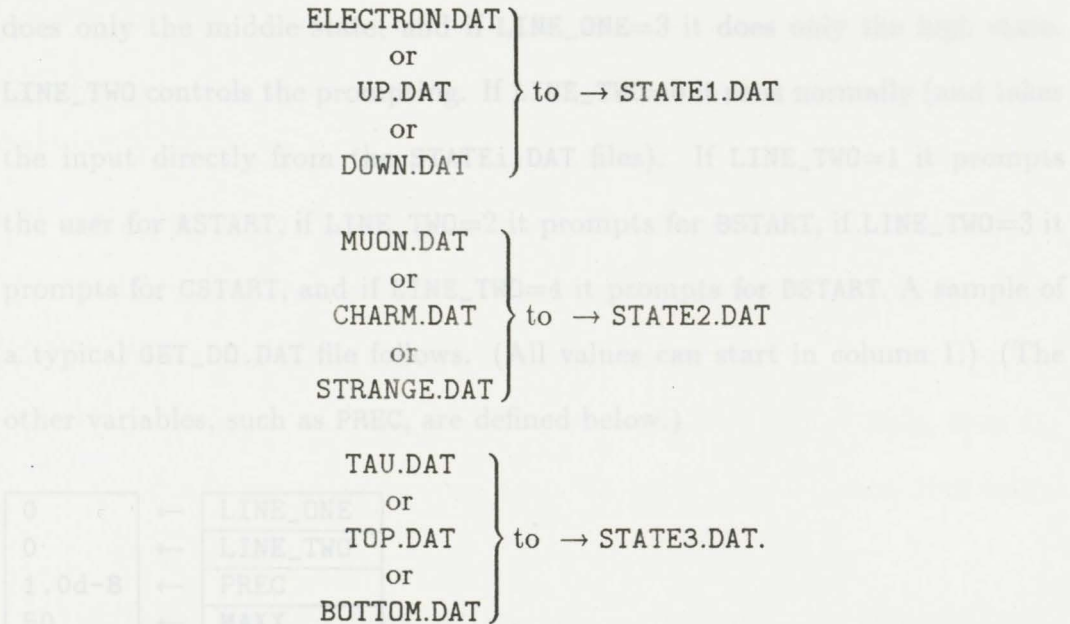
2) Then copy

$$\left. \begin{array}{l} \text{EMT.DAT} \\ \text{or} \\ \text{UCT.DAT} \\ \text{or} \\ \text{DSB.DAT} \end{array} \right\} \text{to } \rightarrow \text{GROUP.DAT.}$$

These files contain the desired masses of each member of the family, and the charge of the family. A sample of a typical EMT.DAT file is:

E/M/T data from Physics Letters Vol 204.	←	Title.
mass electron = 0.51099906	←	Masses (in MeV/c <sup>2</sup> ). (Numbers start in col. 20.)
mass muon = 105.65839		
mass tau = 1784.1		
charge = -1.0	←	Charge (in e).

3) Copy



These files contain the starting guesses for the initial conditions for each of the three members of the family. A sample of a typical ELECTRON.DAT file is:

EMT	←	Title.
A= 11.80	←	Starting values for $y(0)$ , $z(0)$ , $y_1(0)$ , and $y_2(0)$ respectively. (Numbers start in col. 20.)
B= 9.73		
C= 4.34		
D= 4.35		
DELTAB= 0.001	←	The placement of the 1 indicates which digit to test.
DELTAC= 0.001		
DELTAD= 0.001		

4) Change (if desired) GET\_DO.DAT. This is a file which contains instructions for GETBCD. Through it one can run GETBCD normally, or just run

it for one member of the group, and/or have GETBCD prompt for any one of the desired initial values (over-riding what is in the STATEi.DAT file). LINE\_ONE controls the state being done. If LINE\_ONE=0, then it runs normally (and does all three states). If LINE\_ONE=1 it does only the low state, if LINE\_ONE=2 it does only the middle state, and if LINE\_ONE=3 it does only the high state. LINE\_TWO controls the prompting. If LINE\_TWO=0 it runs normally (and takes the input directly from the STATEi.DAT files). If LINE\_TWO=1 it prompts the user for ASTART, if LINE\_TWO=2 it prompts for BSTART, if LINE\_TWO=3 it prompts for CSTART, and if LINE\_TWO=4 it prompts for DSTART. A sample of a typical GET\_DO.DAT file follows. (All values can start in column 1.) (The other variables, such as PREC, are defined below.)

0	←	LINE_ONE
0	←	LINE_TWO
1.0d-8	←	PREC
50	←	MAXX
400	←	1/step size

5) Run GETBCD. (This takes about 45 minutes usually, so it is best done as a batch-job.) GETBCD finds values for  $y(0)$ ,  $z(0)$ ,  $y_1(0)$ , and  $y_2(0)$  that give particle-like solutions using the current F's.

It takes input from: GROUP.DAT, FCONSTANTS.DAT, and GET\_DO.DAT.

As it runs it shows a rough graph on the screen of the equation's behaviour ( $z/y_1/y_2$  are horizontal,  $x$  is vertical). This is useful for spotting instability problems and for generally seeing what is going on.

It works by testing the current guess of the initial value and seeing if that value causes the graph to cross the  $x$ -axis with too many nodes. If that value is acceptable it then increments the last digit and tries again. In this way it finds the next digit of each start value in turn and keeps going until the desired precision is reached. The desired precision is indicated through the variable `PREC`. Set `PREC = 10-n` where  $n$  is the number of desired digits.

`GETBCD` integrates down the  $x$ -axis until

- a)  $x > \text{MAXX}$  (and the program assumes that that initial value is okay), or
- b) the graph crosses with too many nodes (in which case the program knows that initial value is too high), or
- c) the program can tell, via the derivatives of the functions, that the graph is not going to cross again (in which case it knows that initial value is okay).

There are two flags set inside the program: `DO_FORWARD` and `DO_BACKUPS`. If `DO_FORWARD` is set to `TRUE`, `GETBCD` will keep incrementing the last digit until it finds it has gone too far. This is in contrast to the normal behaviour where it will abort if it has to increment the last digit more than ten times. If `DO_BACKUPS` is set to `TRUE`, `GETBCD` will decrement the start-value if it fails on the first iteration. (Normally it aborts if the first iteration fails.) *Both these flags should be set to FALSE for normal operation.* They are intended only as a last resort when strong coupling makes normal operation impossible. (If either flag is enabled there is an iteration-counter built in that will stop the program running after too many iterations, so infinite-loops will

be avoided. But strong couplings can still send the program off searching for non-existent solutions. Normally changing the last digit in one value will not alter the current last digits in the other values, but with strong couplings this is no longer true and each change can mean re-determining all the other values. The problem arises when *those* changes change the value that caused the trouble in the first place. In such instances there are no stable solutions and GETBCD will finally abort through the loop-counter.)

GETBCD uses the subroutine DRKSTP from the CERN Computer Centre Program Library.<sup>14</sup> This is a double-precision Runge-Kutta integration routine for solving the system of  $n$  ( $n \geq 1$ ) simultaneous first-order differential equations of the form:

$$\frac{dy_i}{dx} = f_i(x, y_1, \dots, y_n) \quad (i = 1, 2, \dots, n).$$

It is called as

```
CALL DRKSTP(N,H,X,Y,FCN,W)
```

where the arguments are as follows:

N (Integer.) On entry, N must equal  $n$ , the number of differential equations. (It is left unchanged.)

H (Real, double-precision.) On entry, H must equal the desired step length (i.e.,  $\Delta x$ ). (It is left unchanged.)

X (Real, double-precision.) On entry, X must equal the initial value of  $x$ . On exit it equals  $x + \Delta x$ .

---

<sup>14</sup> CERN Program Library manual, [18], pp. 6.106-6.107.

- Y (Real, double-precision, one-dimensional array.) On entry, Y(i) must contain  $y_i(x)$  ( $i = 1, \dots, n$ ). On exit, Y(i) contains the approximate values of  $y_i(x + \Delta x)$ .
- W (Real, double-precision array.) W should contain  $\geq 3N$  elements. It is used as a scratch work-array.
- FCN This is the name of a user-written, double-precision subroutine (declared EXTERNAL in the calling program) of the form:

$$\text{FCN}(X, Y, F)$$

The subroutine must set

$$f(i) = \frac{dy_i}{dx}(X, Y(1), \dots, Y(N)).$$

So in our case the four second-order differential equations of the field equations (20) are broken into eight first-order differential equations:

$$u_1 \equiv y_1'$$

$$u_2 \equiv y_1''$$

$$u_3 \equiv z_1'$$

$$u_4 \equiv z_1''$$

$$u_5 \equiv y_2'$$

$$u_6 \equiv y_2''$$

$$u_7 \equiv y_3'$$

$$u_8 \equiv y_3''$$

$$u_1 = y'$$

$$u_2 = - \left( \frac{2}{x} y' + z^2 y - y - (FG1)(y_1)^2 y - (FG2)(y_2)^2 y \right)$$

$$u_3 = z'$$

$$u_4 = - \left( \frac{2}{x} z' + y^2 z \right)$$

$$u_5 = y_1'$$

$$u_6 = - \left( \frac{2}{x} y_1 - y_1' + (y_1)^3 + (FE1)y^2 y_1 + \left( \frac{FG3}{FG1} \right) (y_2)^2 y_1 \right)$$

$$u_7 = y_2'$$

$$u_8 = - \left( \frac{2}{x} y_2 - y_2' + (y_2)^3 + (FE2)y^2 y_2 + \left( \frac{FG3 FE2}{FG2 FE1} \right) (y_1)^2 y_2 \right)$$

A good step size for most operations is 1/400 or 1/800. A good MAXX value is 50. (Although 20 speeds things up and is good for rough work.)

GETBCD produces GETBCD.DAT which is a short summary of its attempts to find solutions, and ABCDi.DAT. ( $i = 1, 2, 3$ ). The latter files record the A, B, C, and D values (i.e.,  $y(0)$ ,  $z(0)$ ,  $y_1(0)$ , and  $y_2(0)$ ) that GETBCD found - if it was successful. These files are used for input to other programs.

6) Run ESCALINT3. This program calculates all the various integrals and  $\alpha$ ,  $\omega$ ,  $\sigma$ , etc.

It uses ABCDi.DAT, FCONSTANTS.DAT, and GROUP.DAT as input.

It also takes instructions from ESC\_DO.DAT. This controls the regime being studied and the step size, and contains a flag called RECORD.

A sample of a typical ESC\_DO.DAT file is:

.false.	←	RECORD
0	←	regime
400	←	1/step size

ESCALINT3 automatically finds the point where  $\alpha$  starts diverging and stops integrating there. But it cannot start checking for divergence until  $\alpha$  is the correct side of the  $x$ -axis, so it uses information about the regime being studied. It turns out that if the lowest regime is labelled 'regime 0' then the  $n$ th regime has  $n$  nodes in  $\alpha$ . (This can be seen from the behaviour of the Rosen pair of equations (8).)

RECORD is a flag which will optionally produce an additional two output files E3i.DAT and E4i.DAT. These contain values useful for plotting.

It produces ESCSUM.DAT and ESCLONG.DAT. The former is a short summary of the most-useful quantities. The latter is a more comprehensive data-file designed to be used as input for other programs.

#### 7) Copy

ESCLONG.DAT → to  $\left\{ \begin{array}{l} \text{ESC\_EMT.DAT} \\ \text{or} \\ \text{ESC\_UCT.DAT} \\ \text{or} \\ \text{ESC\_DSB.DAT.} \end{array} \right.$

8) Now at this point we should have a set of initial values that give particle-like solutions for the current group with the current set of F's. The next thing to do is to a) change the F's to give the desired mass-ratios, and/or b) change the  $y(0)$  for each member of the family to bring the  $\alpha$ 's into line. (In practice (a) changes things more than (b) so do it seems best to first get

the F's reasonably close, then worry about the  $\alpha$ 's. In the last stages changes can be made to both at the same time.)

8a) To improve the mass-ratios either run IMPROVE, PARIMP, or TWEAK.

i) Run IMPROVE:

IMPROVE uses FCONSTANTS.DAT, GROUP.DAT, and ESCLONG.DAT as input.

It produces IMPROVE.DAT as output.

This program finds a "perfect" FG1 and FG2 pair that will give the correct mass-ratios. (At least they would be correct if nothing else changed; in practice these new F's will change everything else, hopefully slightly.)

ii) Run TRANSFER. This takes the data from IMPROVE.DAT and inserts the "perfect" F's into FCONSTANTS.DAT.

iii) Copy

$$\text{FCONSTANTS.DAT} \rightarrow \begin{cases} \text{FCONEMT.DAT} \\ \text{or} \\ \text{FCONUCT.DAT} \\ \text{or} \\ \text{FCONUCT.DAT} \end{cases}$$

OR

i) Run PARIMP or TWEAK. They are both parameter-fitting programs that attempt to find common F's to fit two or more groups.

PARIMP is based on PARAFIT<sup>15</sup> and uses ESC\_EMT.DAT, ESC\_UCT.DAT, and ESC\_DSB.DAT for input. It produces PARIMP.DAT which records the F's which give the best (although not necessarily perfect) mass-ratios.

PARAFIT should be called with the following variables set by the calling program:

- L: the number of coefficients to be fitted.
- A(i): ( $i = 1, \dots, L$ ) This array contains the initial guesses for the coefficients. These guesses should be over-estimated in magnitude and of the correct sign.
- N: the number of data pairs.
- X(i), Y(i): ( $i = 1, \dots, N$ ) These pairs hold the data itself.
- E: a convergence criterion. The iteration is stopped when the variance is less than this fraction. (A good starting value is  $E=0.1$ .)
- E1: a convergence factor. ( $0 < E1 \leq 1$ .)
- E1=0 will stall calculation.
- E1=1 will speed up the search, but may cause instability.

PARAFIT returns:

- A(i): ( $i = 1, \dots, L$ ) These are the calculated coefficients that fit the data best.
- D: the standard deviation of the fit.

---

<sup>15</sup> PARAFIT is from *BASIC Scientific Subroutines, Vol. II*, [19], pp. 75-82.

M: the number of iterations performed.

PARAFIT also needs access to a subroutine to evaluate the function

$$y = f(x, A(1), \dots, A(L)).$$

TWEAK uses ESC\_EMT.DAT, ESC\_UCT.DAT, and ESC\_DSB.DAT for input. It produces only output to the screen. TWEAK is an intelligent parameter-tweaker. It assumes there are some variables,  $x(1:n_x)$ , and some functions of these variables,  $y(1:n_y)$  and some desired values,  $d(1:n_y)$ , of these functions. (No assumption is made as to whether the number of  $x$ 's is equal to the number of functions.) It tries to tweak the variables to give the desired values, but instead of just doing each variable in turn (as PARAFIT does for example) it also assumes there is some other criterion for judging what tweakings are better than others. (This might be "don't move the variables far from where they are now", or it might be "keep the variables as small as possible", or "keep the variables as close to certain values as possible.") So the basic strategy is this: for each variable, and each desired output, work out what change would be required of each variable for each output. Then see what change has the least penalty value, do that change, then the next lowest penalty value, and so on. In this way, it is hoped, the variables will converge to some sort of stable solution (and hopefully also one with low undesirability) although it is not guaranteed that this will happen.



APPENDIX E: COMPUTER PROGRAM LISTINGS

## 1) GETBCD

```

c -----
c
c GETBCD
c =====
c Determines B,C, and D. (Initial values of z,y1,y2.)
c Loops three times to do the ground, middle, and highest excited states.
c Reads: STATEi.DAT, FCONSTANTS.DAT, GET_DO.DAT
c Writes: ABCDi.DAT (i=1 for ground ... i=3 for highest),
c GETBCD.DAT.
c
c If DO_BACKUP is enabled, it will back up the attempted solution.
c If DO_FORWARD is enabled, it will go forward more than 10 iterations.
c (Use only for strong coupling problems.)
c
c It reads instructions from GET_DO.DAT:
c 1) Line 1: 0=process normally, i=do only state i. (i=1,2,3)
c 2) Line 2: 0=process normally, i=prompt for A,B,C, or D.
c (i=1 -> A, ..., i=4 -> D)
c 3) Line 3: prec (Controls # of digits. EG. prec=1.0d-4
c means do 4 digits.) (real*8)
c 4) Line 4: maxx (Controls how far down the x-axis to go.)
c (integer)
c
c (All variables are declared, none are implicitly typed.)
c
c Written: - Oct. 1988 by Philip Sharman,
c (based on BVAXB3 by F.I. Cooperstock).
c Last revised: 1988/Nov/30 -scale added for doplot
c Dec/5 -minor changes to scale stuff
c 1989/Jan/17 -STOP changed to EXIT
c Jan/18 -prints out tiles of input STATEi files
c Jan/26 -show f's scaled by 10-4
c Feb/16 -Changed things because of extra
c coupling term...
c -Changed COMMON, FCN; FG3 added.
c -Also added RETURN to DOPLLOT.
c Mar/11 -Added some quit tests.
c -reads GET_DO.DAT to read special
c instructions.
c -writing to GETBCD.DAT moved to DOIT.
c Mar/25 -SAVE added for instructions.
c May/19 -'Nnodes is still okay at' message added.
c -special handling put into FCN to prevent
c division by zero in FG3 term.
c May/28 -MAXX lowered from 30 to 20.
c May/29 -made it back up if it fails on first
c iteration.
c Jun/1 -HMAX deleted.

```

```

c      read(unit=3, fmt=93)      Jun/2      -DO_BACKUP and N_BACKUPS added.
c      read(unit=3, fmt=93)      Jun/4      -SCHEME added.
c      read(unit=3, fmt=93)      Jun/10     -ability added to go more than 10 forward.
c      read(unit=3, fmt=93)      (via NORMAL)
c      close(unit=3)             Jun/14     -NORMAL taken out, DO_FORWARD put in.
c      -formatting cleaned up
c      TAKE SPECIAL INSTRUCTIONS Jun/21     -MAXX and PREC added to GET_DO.DAT.
c      print*, '(Reading FCONSTANTS.DAT)' Aug/29     -if CSTART and DSTART are zero,
c      open(unit=3, file='fconstants.dat') then the program skips C and/or D loops.
c      read(unit=3, fmt=93)      (Useful for doing just uncoupled Rosen
c      read(unit=3, fmt=93)      equations ... remember to also set
c      read(unit=3, fmt=93)      all the F's to zero.)
c      read(unit=3, fmt=93)      -modified to also show the quantities
c      read(unit=3, fmt=93)      (like FG1/FE1 that appear in the field
c      close(unit=3)             equations.
c      format (1)                Aug/31     -step size stuff cleaned up. (JK
c      renamed INVERSE_STEP_SIZE and moved
c      all to one place.)
c      GROUND STATE              Sep/1     -the quit tests changed to only abort
c      call doit(1, prec, inv)    if slopes are > 0.5 (instead of just >0
c      as before).
c      MIDDLE STATE              Sep/3     -reads inverse_step_size from GET_DO.DAT.
c      call doit(2, prec, inv)    Sep/19     -bug fixed in initialization of
c      MIGHTY STATE              backup and forwarding parameters at
c      call doit(3, prec, max, inv) top of KSTEPB loop.
c      And missing line fixed.
c      -----
c      -----

```

```

c      MAIN LOOP
c      real*8      a, b, c, d, zero, one, three,
c      &          fg1, fe1, fg2, fe2, fg3, ten, prec
c      integer    scale, instructions_line1, instructions_line2, maxx,
c      &          inverse_step_size
c      character*100 junk
c      character*12 buffer, buffer2

c      common     a, b, c, d, fg1, fe1, fg2, fe2, fg3
c      common     /const/ zero, one, three
c      common     /plots/ scale
c      common     /special/ instructions_line1, instructions_line2
c      zero = 0.d0
c      one = 1.d0
c      three = 3.d0
c      ten = 10.d0

c      READ FG1 ETC
c      print*, '(Reading FCONSTANTS.DAT)'
c      open(unit=3, file='fconstants.dat', status='old')
c      read(unit=3, fmt=91), junk
c      read(unit=3, fmt=93), fg1

```

```

read(unit=3, fmt=93), fg2
read(unit=3, fmt=93), fe1
read(unit=3, fmt=93), fe2
read(unit=3, fmt=93), fg3
close(unit=3)
c
TAKE SPECIAL INSTRUCTIONS (IF ANY) AND PREC & MAXX FROM 'GET_DO.DAT'
print*, '(Reading GET_DO.DAT)'
open(unit=3, file='get_do.dat', status='old')
read(unit=3, fmt=90), instructions_line1
read(unit=3, fmt=90), instructions_line2
read(unit=3, fmt=*), prec
read(unit=3, fmt=90), maxx
read(unit=3, fmt=90), inverse_step_size
close(unit=3)
90 format (i)

c
GROUND STATE
call doit(1, prec, maxx, inverse_step_size)

c
MIDDLE STATE
call doit(2, prec, maxx, inverse_step_size)

c
HIGHEST STATE
call doit(3, prec, maxx, inverse_step_size)

c
ALL DONE
print*, 'Program is complete.'

c -----formats-----
91 format (a)
83 format (' ', a, t20, e30.16)
93 format (t20, e30.16)
c -----
end

c -----
c
DOIT - does the particular case
subroutine doit(case, prec, maxx, inverse_step_size)
real*8 x, h, y(8), w(99),
& a,b,c,d,
& astart, bstart, cstart, dstart, deltab, deltax,
& deltad, prec, fg1, fg2, fe1, fe2, fg3,
& zero, one, three, ten
logical ok,okay,hunt,tryx,lastpos(3),lastneg(3),backed_up,
& do_backup, do_forward, did_more_than_ten
integer nnodes(3), nstate(3), maxx, case,
& kstepb, istepc, jstepd, inverse_step_size, ki,
& instructions_line1, instructions_line2,
& scale, n_backups, scheme, max_forward

```

```

character*100  junk, title
character*12   buffer,buffer2
common        /special/ instructions_line1, instructions_line2

c             TRYX determines whether to keep searching down the x-axis.
c             OKAY determines whether the axis has been crossed too many times.
c             NNODES(i) records the number of graph crossings for y(1),
c             y(5), and y(7) [y,y1,& y2].
c             NSTATE(i) records the maximum allowable number of graph crossings.

common        a,b,c,d,fg1,fe1,fg2,fe2, fg3
common        /const/ zero,one,three
common        /plots/ scale

external fcn

ten=10.0d0

c             (DO_BACKUP LETS IT BACK UP IF IT FAILS ON FIRST ITERATION)
do_backup= .false.
if (do_backup) print*, 'Ability to do backups is enabled.'
c             (DO_FORWARD LETS IT GO FORWARD MORE THAN 10 ITERATIONS)
do_forward= .false.
if (do_forward) print*, 'Ability to go more than 10 forward is enabled.'
if (do_forward) then
    max_forward= 100
else
    max_forward= 10
endif

print*, 'case=', case
if (case.ne.1 .and. case.ne.2 .and. case.ne.3) call exit(0)

c             TAKE SPECIAL INSTRUCTIONS (IF ANY) FROM 'GET_DO.DAT'
if (instructions_line1.eq.1 .and. case.ne.1) return
if (instructions_line1.eq.2 .and. case.ne.2) return
if (instructions_line1.eq.3 .and. case.ne.3) return
c             USE ZERO TO MEAN PROCESS NORMALLY
if ( instructions_line1.ne.1 .and.
& instructions_line1.ne.2 .and.
& instructions_line1.ne.3) instructions_line1=0
if ( instructions_line2.ne.1 .and.
& instructions_line2.ne.2 .and.
& instructions_line2.ne.3 .and.
& instructions_line2.ne.4) instructions_line2=0

c             READ STARTING VALUES
print*, '(Reading STATEi.DAT)'
if (case.eq.1)
& open(unit=3, file='state1.dat', status='old')
if (case.eq.2)

```

```

&         open(unit=3, file='state2.dat', status='old')
if (case.eq.3)
&         open(unit=3, file='state3.dat', status='old')
read(unit=3, fmt=91), title
print*, 'Title: ', title
read(unit=3, fmt=93), astart
read(unit=3, fmt=93), bstart
read(unit=3, fmt=93), cstart
read(unit=3, fmt=93), dstart
read(unit=3, fmt=93), deltab
read(unit=3, fmt=93), deltac
read(unit=3, fmt=93), deltad
close(unit=3)

c   SPECIAL INSTRUCTIONS MAYBE
if (instructions_line1.ne.0) then
    if (instructions_line2.eq.1) then
        print*, 'astart?'
        read*, astart
        print*, 'astart=', astart
    elseif (instructions_line2.eq.2) then
        print*, 'bstart?'
        read*, bstart
        print*, 'bstart=', bstart
    elseif (instructions_line2.eq.3) then
        print*, 'cstart?'
        read*, cstart
        print*, 'cstart=', cstart
    elseif (instructions_line2.eq.4) then
        print*, 'dstart?'
        read*, dstart
        print*, 'dstart=', dstart
    endif
endif

c   PARAMETERS
ccc   maxx = 20
      print*, 'maxx= ', maxx
cccc  prec = 1.d-8
      print*, 'prec= ', prec
      scheme= 2
c     (SCHEME 0= 000,001, 002)
c     (SCHEME 1= 000,010, 001 ... ROSEN'S)
c     (SCHEME 2= 000,010, 011 ... NORMAL, FRED'S)
print*, 'scheme= ', scheme
if (case.eq.1) then
    nstate(1)=0
    nstate(2)=0
    nstate(3)=0
else if (case.eq.2) then
    if (scheme.eq.0) then
        nstate(1)=0

```

```

nstate(2)=0
nstate(3)=1
print*, 'DOING #0 EXCITATION SCHEME.'
else if (scheme.eq.1) then
nstate(1)=0
nstate(2)=1
nstate(3)=0
print*, 'DOING #1 (ROSEN''S) EXCITATION SCHEME.'
else
nstate(1)=0
nstate(2)=1
nstate(3)=0
print*, 'DOING NORMAL EXCITATION SCHEME.'
endif
else if (case.eq.3) then
if (scheme.eq.0) then
nstate(1)=0
nstate(2)=0
nstate(3)=2
print*, 'DOING #0 EXCITATION SCHEME.'
else if (scheme.eq.1) then
nstate(1)=0
nstate(2)=0
nstate(3)=1
print*, 'DOING #1 (ROSEN''S) EXCITATION SCHEME.'
else
nstate(1)=0
nstate(2)=1
nstate(3)=1
print*, 'DOING NORMAL EXCITATION SCHEME.'
endif
endif
c (NSTATE(I) IS THE NUMBER OF NODES ALLOWED. 0 FOR GROUND STATE,
c 1 FOR NEXT, ETC. I=1 FOR Y, I=2 FOR Y1, I=3 FOR Y2)

call date(buffer)
call time(buffer2)
print*, 'GETBCD running on ', buffer, ' at ', buffer2

a = astart
b = bstart
c = cstart
d = dstart

if (
& (instructions_line1.eq.1).and.(case.eq.1)
& .or. (instructions_line1.eq.2).and.(case.eq.2)
& .or. (instructions_line1.eq.3).and.(case.eq.3)
& .or. (instructions_line1.eq.0).and.(case.eq.1))
& open(unit=4,file='getbcd.dat',status='new')
print*, '(Writing GETBCD.DAT)'

```

```

print*, 'fg1=', fg1, ' = ', (fg1*1.0d4), ' x10^-4'
print*, 'fg2=', fg2, ' = ', (fg2*1.0d4), ' x10^-4'
print*, 'fe1=', fe1, ' = ', (fe1*1.0d4), ' x10^-4'
print*, 'fe2=', fe2, ' = ', (fe2*1.0d4), ' x10^-4'
print*, 'fg3=', fg3, ' = ', (fg3*1.0d4), ' x10^-4'
if (fe1.ne.zero .and. fe2.ne.zero .and. fg1.ne.zero .and. fg2.ne.zero) then
  print*, ' Field equation quantities:'
  print*, 'fg1/fe1 = ', fg1/fe1
  print*, 'fg2/fe2 = ', fg2/fe2
  print*, 'fe1 = ', fe1
  print*, 'fe2 = ', fe2
  print*, 'fg3/fg1 = ', fg3/fg1
  print*, '(fg3*fe2)/(fg2*fe1) = ', (fg3*fe2)/(fg2*fe1)
  print*, ' '
endif
print*, 'astart=', astart
print*, 'bstart=', bstart
print*, 'cstart=', cstart
print*, 'dstart=', dstart

if (case.eq.1)
&   write(unit=4,fmt=*) 'Results of GETBCD run on ',
&   buffer, ' at ', buffer2

write(unit=4,fmt=*) ' '
write(unit=4,fmt=*) 'case=', case
write(unit=4,fmt=*) 'fg1=', fg1, ' = ', (fg1*1.0d4), ' x10^-4'
write(unit=4,fmt=*) 'fg2=', fg2, ' = ', (fg2*1.0d4), ' x10^-4'
write(unit=4,fmt=*) 'fe1=', fe1, ' = ', (fe1*1.0d4), ' x10^-4'
write(unit=4,fmt=*) 'fe2=', fe2, ' = ', (fe2*1.0d4), ' x10^-4'
write(unit=4,fmt=*) 'fg3=', fg3, ' = ', (fg3*1.0d4), ' x10^-4'
write(unit=4,fmt=*) 'astart=', astart
write(unit=4,fmt=*) 'bstart=', bstart
write(unit=4,fmt=*) 'cstart=', cstart
write(unit=4,fmt=*) 'dstart=', dstart

c   MAIN PART
c   =====
  hunt = .true.
  print*, 'Looking for ', nstate(1), nstate(2), nstate(3),
&   ' type solution.'
  write(unit=4,fmt=*) 'Looking for ', nstate(1), nstate(2),
&   nstate(3), ' type solution.'

  h = one/inverse_step_size
  print*, 'Step size = 1/', inverse_step_size, ' = ', h

c   SKIP AHEAD IF OBEYING SPECIAL INSTRUCTIONS
  if (instructions_line2.eq.4) goto 774      ! (GOTO D LOOP)
  if (instructions_line2.eq.3) goto 773      ! (GOTO C LOOP)

```

```

c      KEEP LOOKING FOR B, C, & D UNTIL THE DESIRED PRECISION ("PREC") IS REACHED
do 154 while(hunt)      !(THIS HUGE OUTER LOOP NOT INDENTED)

c -----KSTEPB LOOP-----
backed_up= .false.
did_more_than_ten= .false.
n_backups= 0

270  continue      !(THIS IS A TARGET LINE FOR A LATER GOTO)
if (do_backup .or. do_forward) then
    print*, 'After line 270 (top of KSTEPB loop):'
    print*, ' backed_up = ', backed_up
    print*, ' did_more_than_ten = ', did_more_than_ten
    print*, ' n_backups = ', n_backups
endif

scale= int(abs(a)+1)
okay=.true.
kstepb=0
do 152 while(okay .and. kstepb.le.max_forward)
    lastpos(1)=.true.
    lastpos(2)=.true.
    lastpos(3)=.true.
    lastneg(1)=.false.
    lastneg(2)=.false.
    lastneg(3)=.false.
c      (LASTPOS & LASTNEG RECORD WHAT SIDE OF THE AXIS THE GRAPH WAS LAST ON)
nnodes(1)=0
nnodes(2)=0
nnodes(3)=0
c      (NNODES RECORDS NUMBER OF NODES FOUND)
print*, 'Now inside the KSTEPB loop. KSTEPB=', kstepb
b = bstart + deltab*kstepb
print 044, a
print 144, b, deltab
print 244, c, deltac
print 344, d, deltad
x = zero
y(1) = a
y(2) = zero
y(3) = b
y(4) = zero
y(5) = c
y(6) = zero
y(7) = d
y(8) = zero
ki=1
tryx = .true.
do 151 while(okay .and. tryx .and. x.le.maxx)
    call drkstp(8, h, x, y, fcn, w)
    if (mod(ki, 2*inverse_step_size).eq.1) then

```

```

call drkstp(8, h, x, y, fcn, w)
if (mod(ki, 2*inverse_step_size).eq.1) then
    call doplot(y(1))
endif
c
HAVE WE CROSSED THE AXIS?
&
&
if ((lastpos(1) .and. y(1).lt.0) .or.
    (lastpos(2) .and. y(5).lt.0) .or.
    (lastpos(3) .and. y(7).lt.0)) then
c
    YES WE'VE CROSSED IT FROM + TO -
    call doplot(y(1))
    if (lastpos(1) .and. y(1).lt.0) then
        print*, 'y(1) gone pos to neg. x,y(1)=' ,x,y(1)
        nnodes(1)=nnodes(1)+1
        lastpos(1)=.false.
        lastneg(1)=.true.
    endif
    if (lastpos(2) .and. y(5).lt.0) then
        print*, 'y(5) gone pos to neg. x,y(5)=' ,x,y(5)
        nnodes(2)=nnodes(2)+1
        lastpos(2)=.false.
        lastneg(2)=.true.
    endif
    if (lastpos(3) .and. y(7).lt.0) then
        print*, 'y(7) gone pos to neg. x,y(7)=' ,x,y(7)
        nnodes(3)=nnodes(3)+1
        lastpos(3)=.false.
        lastneg(3)=.true.
    endif
    endif
    if (nnodes(1).gt.nstate(1)) then
&
        print*, 'TOO BIG: nnodes(1),nstate(1)=' ,
            nnodes(1),nstate(1)
        okay=.false.
    endif
    if (nnodes(2).gt.nstate(2)) then
&
        print*, 'TOO BIG: nnodes(2),nstate(2)=' ,
            nnodes(2),nstate(2)
        okay=.false.
    endif
    if (nnodes(3).gt.nstate(3)) then
&
        print*, 'TOO BIG: nnodes(3),nstate(3)=' ,
            nnodes(3),nstate(3)
        okay=.false.
    endif
    endif
    else if ((lastneg(1) .and. y(1).gt.0) .or.
    (lastneg(2) .and. y(5).gt.0) .or.
    (lastneg(3) .and. y(7).gt.0)) then
c
        YES WE'VE CROSSED IT FROM - TO +
        call doplot(y(1))
        if (lastneg(1) .and. y(1).gt.0) then
            print*, 'y(1) gone neg to pos. x,y(1)=' ,x,y(1)
            nnodes(1)=nnodes(1)+1
            lastpos(1)=.true.
        endif
    endif

```

```

181         if (lastneg(2) .and. y(5).gt.0) then
                print*, 'y(5) gone neg to pos. x,y(5)=',x,y(5)
                nnodes(2)=nnodes(2)+1
                lastpos(2)=.true.
                lastneg(2)=.false.
            endif
182         if (lastneg(3) .and. y(7).gt.0) then
                print*, 'y(7) gone neg to pos. x,y(7)=',x,y(7)
                nnodes(3)=nnodes(3)+1
                lastpos(3)=.true.
                lastneg(3)=.false.
            endif
CHECK WHAT WE HAVE
183         if (nnodes(1).gt.nstate(1)) then
                print*, 'TOO BIG: nnodes(1),nstate(1)=',
                nnodes(1),nstate(1)
                okay=.false.
            endif
            if (nnodes(2).gt.nstate(2)) then
                print*, 'TOO BIG: nnodes(2),nstate(2)=',
                nnodes(2),nstate(2)
                okay=.false.
            endif
            if (nnodes(3).gt.nstate(3)) then
                print*, 'TOO BIG: nnodes(3),nstate(3)=',
                nnodes(3),nstate(3)
                okay=.false.
            endif
            endif
c         QUIT TEST -- TEST Y(1)=Y (EVEN THOUGH WE ARE STEPPING Z?)
            if ( nstate(1).eq.0 .and. x.gt.3.0 .and.
            &         y(2).gt.0.5) then
                call doplot(y(1))
                print *, 'Quitting: slope of y(1)=y is >0.5'
                print 1000, '(x= ', x, ' )'
                tryx=.false.
1000             print format (a, f7.2)
            endif
            if ( nstate(1).eq.1 .and.
            &         y(1).lt.(-1*a) ) then
                call doplot(y(1))
                print*, 'Quitting: y(1)=y is less than -A.'
                print 1000, '(x= ', x, ' )'
                tryx=.false.
            endif
            if ( nstate(1).eq.2 .and.
            &         y(1).gt.a) then
                call doplot(y(1))
                print*, 'Quitting: y(1)=y is greater than A.'
                print 1000, '(x= ', x, ' )'
                tryx=.false.
            endif
            endif
            did_more_than_ten = .true.

```

```

151         ki=ki+1 (MAXIMIZE: Graph took more than 10 iterations to cross.)
           continue
           if (okay) then
               kstepb=kstepb+1
               if (tryx) print*, 'Number of nodes is still okay at x= ', x
           endif
152     continue
c -----

c CHECK WHAT WE HAVE
if (kstepb.eq.0) then
c IT BOMBED OUT ON THE FIRST TRY - ABORT OR BACKUP
  if (do_backup) then
      print*, 'WARNING: Graph failed on first iteration.'
      print*, 'BACKUP'
      print*, 'backed_up= .true.'
      n_backups= n_backups + 1
      if (n_backups.gt.100 .or. b.lt.zero) then
          print*, 'FAILURE: backups aborted.'
          call exit(0)
      endif
      LOWER THE LAST DIGIT OF BSTART
      bstart= bstart- deltab
      b= bstart
      print*, 'Re-doing B loop with BSTART= ', bstart
      print*, 'backed_up = ', backed_up
      print*, 'did_more_than_ten = ', did_more_than_ten
      print*, 'n_backups = ', n_backups
      print*, ' '
      goto 270 ! (GO TO KSTEPB LOOP AGAIN)
  else
      ABORT
      print*, 'FAILURE: Graph failed on first iteration. (B loop)'
      print 044, a
      print 144, b, deltab
      print 244, c, deltac
      print 344, d, deltad
      call exit(0)
  endif
else if (kstepb.ge.max_forward+1) then
c WE FAILED TO FIND A GOOD VALUE - ABORT
  print*, 'FAILURE: Graph did not cross axis after ',
  & max_forward, ' iterations. (B loop)'
  print 044, a
  print 144, b, deltab
  print 244, c, deltac
  print 344, d, deltad
  call exit(0)
else
c WE ARE ON THE RIGHT TRACK. NARROW DOWN THE PRECISION.
  if (kstepb.ge.11) then
      did_more_than_ten= .true.
  endif
endif

```

```

        print*, '(WARNING: Graph took more than 10 iterations to cross.)'
    endif
    print*, 'Look for next digit of B.'
    bstart=bstart + (kstepb-1)*deltab
    deltab = deltab / ten
    b = bstart
    kstepb=0
endif

c -----ISTEPC LOOP-----
773 continue      !(TARGET FOR A POSSIBLE GOTO)
if (do_backup .or. do_forward) then
    print*, 'After line 773 (top of ISTEPC loop):'
    print*, ' backed_up = ', backed_up
    print*, ' did_more_than_ten = ', did_more_than_ten
    print*, ' n_backups = ', n_backups
endif

scale= int(abs(c)+1)
print*, ' '
okay = .true.
istepc=0
if (cstart.eq.zero) then
    print*, '(Skipping the C loop as CSTART=0.)'
    goto 774      ! (GOTO JSTEP D LOOP)
endif

do 252 while(okay .and. istepc.le.max_forward)
    lastpos(2)=.true.
    lastneg(2)=.false.
    nnodes(2)=0
    print*, 'Now inside the ISTEPC loop. ISTEPC=', istepc
    c = cstart + deltac*istepc
    print 044, a
    print 144, b, deltab
    print 244, c, deltac
    print 344, d, deltad
    x = zero
    y(1) = a
    y(2) = zero
    y(3) = b
    y(4) = zero
    y(5) = c
    y(6) = zero
    y(7) = d
    y(8) = zero
    ki=1
    tryx = .true.
    do 251 while(okay .and. tryx .and. x.le.maxx)
        call drkstp(8, h, x, y, fcn, w)
        if (mod(ki, 2*inverse_step_size).eq.1) then
            call doplot(y(5))
        endif
    enddo
endif

```

```

c          HAVE WE CROSSED THE AXIS?
c          if (lastpos(2) .and. y(5).lt.0) then
c              YES WE'VE CROSSED IT FROM + TO -
252          continue
c              print*, 'gone pos to neg. x,y(5)=',x,y(5)
c              call doplot(y(5))
c              nnodes(2)=nnodes(2)+1
c              lastpos(2)=.false.
c              lastneg(2)=.true.
c              if (nnodes(2).gt.nstate(2)) then
&                  print*, 'TOO BIG: nnodes(2),nstate(2)=',
c                      nnodes(2),nstate(2)
c                      okay=.false.
c              endif
c              else if (lastneg(2) .and. y(5).gt.0) then
c              YES WE'VE CROSSED IT FROM - TO +
c              print*, 'gone neg to pos. x,y(5)=',x,y(5)
c              call doplot(y(5))
c              nnodes(2)=nnodes(2)+1
c              lastpos(2)=.true.
c              lastneg(2)=.false.
c              if (nnodes(2).gt.nstate(2)) then
&                  print*, 'TOO BIG: nnodes(2),nstate(2)=',
c                      nnodes(2),nstate(2)
c                      okay=.false.
c              endif
c              endif
c          QUIT TEST -- TEST Y(5)=y1
&          if ( nstate(2).eq.0 .and. x.gt.3.0 .and.
c              y(6).gt.0.5) then
c              call doplot(y(5))
c              print*, 'Quitting: slope of y(5)=y1 is >0.5'
c              print 1000, '(x= ', x, ' )'
c              tryx=.false.
c          endif
&          if ( nstate(2).eq.1 .and.
c              y(5).lt.(-1*c) ) then
c              call doplot(y(5))
c              print*, 'Quitting: y(5)=y1 is less than -C.'
c              print 1000, '(x= ', x, ' )'
c              tryx=.false.
c          endif
&          if ( nstate(2).eq.2 .and.
c              y(5).gt.c) then
c              call doplot(y(5))
c              print*, 'Quitting: y(5)=y1 is greater than C.'
c              print 1000, '(x= ', x, ' )'
c              tryx=.false.
c          endif
251          ki=ki+1
          continue
          if (okay) then

```

```

                istepc=istepc+1
                if (tryx) print*, 'Number of nodes is still okay at x= ', x
            endif
252     continue
c -----

c     CHECK WHAT WE HAVE
    if (istepc.eq.0) then
c     IT BOMBED OUT ON THE FIRST TRY - ABORT OR BACKUP
        if (do_backup) then
c     print*, 'WARNING: Graph failed on first iteration.'
            BACKUP
c     backed_up= .true.
            n_backups= n_backups + 1
            if (n_backups.gt.100 .or. c.lt.zero) then
                print*, 'FAILURE: backups aborted'
                call exit(0)
            endif
c     LOWER THE LAST DIGIT OF CSTART
            cstart= cstart- deltax
            c= cstart
            print*, 'Re-doing C loop with CSTART= ', cstart
            print*, ' backed_up = ', backed_up
            print*, ' did_more_than_ten = ', did_more_than_ten
            print*, ' n_backups = ', n_backups
            print*, ' '
            goto 773      !(GOTO C LOOP)
        else
c     ABORT
            print*, 'FAILURE: Graph failed on first iteration. (C loop)'
            print 044, a
            print 144, b, deltax
            print 244, c, deltax
            print 344, d, deltax
            call exit(0)
        endif
c     else if (istepc.ge.max_forward+1) then
        WE FAILED TO FIND A GOOD VALUE - ABORT
        print*, 'FAILURE: Graph did not cross axis after ',
&         max_forward, ' iterations.' (C loop)'
        print 044, a
        print 144, b, deltax
        print 244, c, deltax
        print 344, d, deltax
        call exit(0)
c     else
        WE ARE ON THE RIGHT TRACK.  NARROW DOWN THE PRECISION.
        if (istepc.ge.11) then
            did_more_than_ten= .true.
            print*, '(WARNING: graph took more than 10 iterations to cross.)'
        endif
        print*, 'Look for next digit of C.'
    endif

```



```

c          endif          YES WE'VE CROSSED IT FROM + TO -
352          continue     print*, 'gone pos to neg. x,y(7)=' ,x,y(7)
c          -----       call doplot(y(7))
c                          nnodes(3)=nnodes(3)+1
c          CHECK WHAT WE HAVE      lastpos(3)=.false.
c          IF (jstepd.eq.0) then    lastneg(3)=.true.
c          IT BUMPED OUT (         if (nnodes(3).gt.nstate(3)) then
&          IF (do_backup) then     print*, 'TOO BIG: nnodes(3),nstate(3)=' ,
c          BACKUP                  nnodes(3),nstate(3)
c          backed_up               okay=.false.
c          endif
c          else if (lastneg(3) .and. y(7).gt.0) then
c          IF (n.)               YES WE'VE CROSSED IT FROM - TO +
c          LOWER                 print*, 'gone neg to pos. x,y(7)=' ,x,y(7)
c          data                  call doplot(y(7))
c          data                  nnodes(3)=nnodes(3)+1
c          data                  lastpos(3)=.true.
c          data                  lastneg(3)=.false.
&          data                  if (nnodes(3).gt.nstate(3)) then
c          data                  print*, 'TOO BIG: nnodes(3),nstate(3)=' ,
c          data                  nnodes(3),nstate(3)
c          data                  okay=.false.
c          data                  endif
c          data                  endif
c          QUIT TEST -- TEST Y(7)=y2
c          if ( nstate(3).eq.0 .and. x.gt.3.0 .and.
&          y(8).gt.0.5) then
c          print call doplot(y(7))
c          print print*, 'Quitting: slope of y(7)=y2 is >0.5'
c          print print 1000, '(x= ', x, ' )'
c          print tryx=.false.
c          endif
&          if ( nstate(3).eq.1 .and.
c          y(7).lt.(-1*d) ) then
c          call doplot(y(7))
c          print*, 'Quitting: y(7)=y2 is less than -D.'
c          print 1000, '(x= ', x, ' )'
c          print tryx=.false.
c          endif
&          if ( nstate(3).eq.2 .and.
c          y(7).gt.d) then
c          call doplot(y(7))
c          print*, 'Quitting: y(7)=y2 is greater than D.'
c          print 1000, '(x= ', x, ' )'
c          print tryx=.false.
c          endif
c          if (jstepd.ge.1) endif
c          print*, '(WARNING: graph took more than 10 iterations to cross.)'
c          did not ki=ki+1
351          endif
c          continue
c          if (okay) then
c          DO WE HAVE TO jstepd=jstepd+1 WE FINALLY DONE?
c          if (backed_up) if (tryx) print*, 'Number of nodes is still okay at x= ', x

```

```

                endif
352     continue
c -----
c     CHECK WHAT WE HAVE
    if (jstepd.eq.0) then
c         IT BOMBED OUT ON THE FIRST TRY - ABORT OR BACKUP
        if (do_backup) then
c             print*, 'WARNING: Graph failed on first iteration.'
            BACKUP
c             backed_up= .true.
            n_backups= n_backups + 1
            if (n_backups.gt.100 .or. d.lt.zero) then
                print*, 'FAILURE: backups aborted.'
                call exit(0)
            endif
c             LOWER THE LAST DIGIT OF DSTART
            dstart= dstart- deltax
            d= dstart
            print*, 'Re-doing D loop with DSTART= ', dstart
            print*, ' backed_up = ', backed_up
            print*, ' did_more_than_ten = ', did_more_than_ten
            print*, ' n_backups = ', n_backups
            print*, ' '
            goto 774      !(GOTO D LOOP)
        else
c             ABORT
            print*, 'FAILURE: Graph failed on first iteration. (D loop)'
            print 044, a
            print 144, b, deltax
            print 244, c, deltac
            print 344, d, deltax
            call exit(0)
        endif
c     else if (jstepd.eq.max_forward+1) then
        WE FAILED TO FIND A GOOD VALUE - ABORT
c         print*, 'FAILURE: Graph did not cross axis after',
            &             max_forward, ' iterations. (D loop)'
            print 044, a
            print 144, b, deltax
            print 244, c, deltac
            print 344, d, deltax
            call exit(0)
        endif
c     if (jstepd.ge.11) then
            print*, '(WARNING: graph took more than 10 iterations to cross.)'
            did_more_than_ten= .true.
        endif
c     DO WE HAVE TO LOOP AGAIN, OR ARE WE FINALLY DONE?
        if (backed_up .or. did_more_than_ten) then

```

```

c          LOOP AGAIN
          if (backed_up) print*, 'WARNING: We backed up somewhere.'
          if (did_more_than_ten) print*, 'WARNING: We did more than 10 itera-
tions somewhere.'
          print*, 'Increasing DELTAB and DELTAC and going back to B loop again.'
          deltab= deltab*10
          deltac= deltac*10
          print*, 'Now re-setting flags and counter.'
          backed_up = .false.
          did_more_than_ten = .false.
          n_backups = 0
          print*, ' backed_up = ', backed_up
          print*, ' did_more_than_ten = ', did_more_than_ten
          print*, ' n_backups = ', n_backups
          print*, ' '
          goto 270          !(GO TO STEPB LOOP AGAIN)
else
c          WE ARE ON THE RIGHT TRACK.  NARROW DOWN THE PRECISION.
          print*, '(We are okay.  No backups or >10 iterations.)'
          print*, 'Look for next digit of d.'
          dstart=dstart + (jstepd-1)*deltad
          deltad = deltad / ten
          d = dstart
          jstepd=0
endif

c -----
c -----
850 continue          ! (TARGET FOR AN EARLIER GOTO)
c          HAVE WE ACHEIVED THE DESIRED PRECISION?  IF NOT, LOOP BACK AND
c          GET ANOTHER DIGIT FOR B, C, AND D.
          if (deltab.ge.prec) then
          print*, 'Looping again to increase precision.'
          print 111, bstart, deltab
          print 211, cstart, deltac
          print 311, dstart, deltad
          write (unit=4,fmt=*) 'Looping again to increase precision.'
          write (unit=4,fmt=111) bstart, deltab
          write (unit=4,fmt=211) cstart, deltac
          write (unit=4,fmt=311) dstart, deltad
else
          print*, 'Desired precision (', prec,
&          ') has been reached.'
          print*, 'a,b,c,d='
          print*, a
          print*, b
          print*, c
          print*, d
          write(unit=4,fmt=*) 'Desired precision (', prec,
&          ') has been reached.'
          write(unit=4,fmt=*) 'a,b,c,d='
          write(unit=4,fmt=*) a

```

```

        write(unit=4,fmt=*) b
        write(unit=4,fmt=*) c
        write(unit=4,fmt=*) d
        hunt=.false.
    endif

    154 continue      !(THIS HUGE OUTER LOOP NOT INDENTED)

    print*, '(Writing ABCDi.DAT)'
    if (case.eq.1) then
        open(unit=3, file='abcd1.dat', status='new')
    else if (case.eq.2) then
        open(unit=3, file='abcd2.dat', status='new')
    else if (case.eq.3) then
        open(unit=3, file='abcd3.dat', status='new')
    endif
    call date(buffer)
    call time(buffer2)
    write(unit=3, fmt=*), ' Results found by GETBCD run on ',
&      buffer, ' at ', buffer2
    write(unit=3, fmt=83), 'a=', a
    write(unit=3, fmt=83), 'b=', b
    write(unit=3, fmt=83), 'c=', c
    write(unit=3, fmt=83), 'd=', d
    endfile(unit=3)
    close(unit=3)

c -----
111 format (' new precision: bstart=',f18.15,' deltab=',f18.15)
211 format (' new precision: cstart=',f18.15,' deltac=',f18.15)
311 format (' new precision: dstart=',f18.15,' deltad=',f18.15)
044 format (' A=', F18.15)
144 format (' B=', F18.15, '   DeltaB=', F18.15)
244 format (' C=', F18.15, '   DeltaC=', F18.15)
344 format (' D=', F18.15, '   DeltaD=', F18.15)
91  format (a)
93  format (t20, e30.16)
83  format (' ', a, t20, e30.16)

c -----
    print*, 'Doit ends.'
    write(unit=4,fmt=*) 'Doit ends.'
    if (
&      (instructions_line1.eq.1).and.(case.eq.1)
&      .or. (instructions_line1.eq.2).and.(case.eq.2)
&      .or. (instructions_line1.eq.3).and.(case.eq.3)
&      .or. (instructions_line1.eq.0).and.(case.ge.3))
&      then
        endfile(unit=4)
        close(unit=4)
    endif

    return

```

```

end
yprime(7) = zero
yprime(8) = (one - a**2*fe2 - d*d) *c/three
if (fg3.ne.zero) then

c =====
c FCN
c ===
c subroutine fcn(x,y,yprime)
c (Evaluates derivatives.)
c (Taken from ESCALINT3, 1989/Feb/16. RETURN added at end.)
c (FG3 term not added if FG3=0. This allows the denominator to
c be zero for when the uncoupled case is wanted.)

real*8 a,b,c,d, fg1,fg2,fe1,fe2, fg3, zero,one,three
real*8 p,q,u,v, pb,qb,pbb,qbb
real*8 x,y(8),yprime(8)

common a,b,c,d, fg1,fe1,fg2,fe2, fg3
common /const/ zero,one,three

p = y(1)
q = y(2)
u = y(3)
v = y(4)
pb = y(5)
qb = y(6)
pbb = y(7)
qbb = y(8)
if (x.gt.zero) then
    yprime(1) = q
    yprime(2) = (one-u**2)*p - 2*q/x +(pb**2)*p*fg1 + (pbb**2)*p*fg2
    yprime(3) = v
    yprime(4) = -u*p**2 - 2*v/x
    yprime(5) = qb
    yprime(6) = -2*qb/x + (one - pb*pb)*pb - p*p*pb*fe1
    if (fg3.ne.zero) then
        yprime(6)= yprime(6) - (fg3/fg1)*pbb*pbb*pb
    endif
    yprime(7) = qbb
    yprime(8) = -2*qbb/x + (one -pbb*pbb)*pbb - p*p*pbb*fe2
    if (fg3.ne.zero) then
        yprime(8) = yprime(8) - (fg3*fe2/(fg2*fe1))*pb*pb*pbb
    endif
endif
else
    yprime(1) = zero
    yprime(2) = (one-b*b+c*c*fg1+d*d*fg2)*a/three
    yprime(3) = zero
    yprime(4) = -b*a**2/three
    yprime(5) = zero
    yprime(6) = (one - a*a*fe1 -c*c) *c/three
    if (fg3.ne.zero) then
        yprime(6)= yprime(6) - (fg3/fg1)*d*d *c/three
    endif
endif

```



## 2) ESCALINT3

```

c -----
c -----
c ESCALINT3
c =====
c
c Computes psi, w, beta, sigma, etc.
c Writes output to: ESCSUM.DAT, ESCLONG.DAT.
c Reads: ABCDi.DAT, GROUP.DAT, FCONSTANTS.DAT, ESC_DO.DAT.
c From ESC_DO.DAT it reads: record, regime, inverse_step_size.
c
c If "record" is set to true, it writes out a number of
c variables to "E3i.DAT" and "E4i.DAT" for later plotting.
c
c Stops integrating when alpha starts to diverge (or when
c x .ge. 100 [via ijlast] ).
c (It starts looking for divergence at x+2 after the point
c when alpha has crossed the axis "regime" number of times.)
c
c (All variables declared; none implicit)
c
c Unit          File
c ----          -
c 3             FCONSTANTS.DAT (read)
c "             ABCDi.DAT (read)
c 4             GROUP.DAT (read)
c "             ESC_DO.DAT (read)
c 6             ESCSUM.DAT (write)
c 7             ESCLONG.DAT (write)
c 8             E3i.DAT (write)
c 9             E4i.DAT (write)
c
c Written: 1988/Nov      (by Philip Sharman)
c                      (modified from ESCALINT by F.I. Cooperstock)
c Revised: 1988/Nov/27  (added ESCSUM.DAT)
c                      28  (fixed bug about -b in /CONST/)
c                      (uses abs(alphaf) to test for alpha divergence)
c                      (print out psi instead of w properly in SUM.dat)
c                      (added w1,w2,w3 in /OMEGAS/)
c                      29  (alter order of fg's printout)
c                      1988/Dec/4 (add ESCLONG.DAT writing)
c                      (reversed order of printing pua,pta etc)
c                      (added units to sigma, epsilon)
c                      10  (show more digits in escsum.dat)
c                      11  (added E3.DAT stuff) (deleted ESCALINT3.DAT)
c                      13  (took out the idea of using +b or -b; use a
c                          negative bstart as input instead)
c                      1989/Jan/15 (E5.DAT added)
c                      18  (ESCSUM.DAT altered for better printout)
c                      26  (scale f's by 10^-4 in ESCSUM.DAT)
c                      26  (Added the extra coupling term....
c                      MAIN LOOP: Feb/15

```

```

c      DO GROUDED CASE                Altered FCN to include F3 term.
c      print*, '(Reading ESCSUM.DAT)' Altered COMMON to include FG3.
c      open(unit=3, file='escsum.dat', status='old') Altered FG3 I/O.
c      read(unit=3, fmt=93)           Changed omega, pextra, etc.)
c      read(unit=3, fmt=93) Mar/1     (Added warnings about when things diverge, etc.)
c      read(unit=3, fmt=93) May/19    Changed FCN to only add FG3 term if it is non
c      read(unit=3, fmt=93)           zero. Other places changed for that too.
c      read(unit=3, fmt=93) /21      (Added PUA etc to ESCSUM.DAT)
c      close(unit=3) Jun/4           (Do not compute psif if fe1 or fe2 are 0)
c      call dofs(1)                   (WARN10 and WARN11 added to show when y' and
c                                     z' go positive)
c      DO MIDDLE CASE                (NOTE: this warning not always very useful
c      print*, '(Reading ESCSUM.DAT)' for z. Could be more sophisticated?)
c      open(unit=3, file='escsum.dat', status='old') Jun/12 (More digits shown in ESCSUM.DAT)
c      read(unit=3, fmt=93) Jun/13   (Percentage error shown in ESCSUM.DAT)
c      read(unit=3, fmt=93) Jun/14   (Formatting cleaned up)
c      read(unit=3, fmt=93) Jun/18   (IMPLICIT typing removed)
c      read(unit=3, fmt=93) Aug/4    (E3i.DAT file should have x=0 values in now
c      read(unit=3, fmt=93)           for x,y,z,y1, and y2.)
c      close(unit=3) Aug/29          (Changed y**2*z*x**2 to z-alpha/x (for Beta)
c      call dofs(2)                   in E3i.DAT)
c      read(unit=3, fmt=93) Aug/31   (IJLAST was set to stop at 20; changed it
c      DO HIGHEST EXCITED CASE        to go to 100 if necessary)
c      print*, '(Reading ESCSUM.DAT)' Sep/3 (ESC_DO.DAT added.)
c      -----
c      -----
c      character*80 junk
c      real*8 a, b, c, d, fg1, fe1, fg2, fe2, fg3,
& zero, one, two, three,
& w1, w2, w3
c      common a, b, c, d, fg1, fe1, fg2, fe2, fg3
c      common /const/ zero, one, two, three
c      common /omegas/ w1, w2, w3
c      print*, 'Program ends.'
c      zero = 0.d0
c      one = 1.d0
c      two = 2.d0
c      three = 3.d0
c      READ FG1 ETC
c      print*, '(Reading FCONSTANTS.DAT)'
c      open(unit=3, file='fconstants.dat', status='old')
c      read(unit=3, fmt=80), junk
c      read(unit=3, fmt=93), fg1
c      read(unit=3, fmt=93), fg2
c      read(unit=3, fmt=93), fe1
c      read(unit=3, fmt=93), fe2
c      read(unit=3, fmt=93), fg3
c      close(unit=3)
c      MAIN LOOP:

```

```

c      DO GROUND CASE
      print*, '(Reading ABCD1.DAT)'
      open(unit=3, file='abcd1.dat', status='old')
      read(unit=3, fmt=80), junk
      read(unit=3, fmt=93), a
      read(unit=3, fmt=93), b
      read(unit=3, fmt=93), c
      read(unit=3, fmt=93), d
      close(unit=3)
      call doit(1)

c      DO MIDDLE CASE
      print*, '(Reading ABCD2.DAT)'
      open(unit=3, file='abcd2.dat', status='old')
      read(unit=3, fmt=80), junk
      read(unit=3, fmt=93), a
      read(unit=3, fmt=93), b
      read(unit=3, fmt=93), c
      read(unit=3, fmt=93), d
      close(unit=3)
      call doit(2)

c      DO HIGHEST EXCITED CASE
      print*, '(Reading ABCD3.DAT)'
      open(unit=3, file='abcd3.dat', status='old')
      read(unit=3, fmt=80), junk
      read(unit=3, fmt=93), a
      read(unit=3, fmt=93), b
      read(unit=3, fmt=93), c
      read(unit=3, fmt=93), d
      close(unit=3)
      call doit(3)

c      ALL DONE
      print*, 'Program ends.'

c      -----formats-----
80      format (a)
93      format (t20, e30.16)
c      -----

c      -----
c      -----
c      DOIT
      subroutine doit(case)
      real*8      charge, x, h, delta, alpha, beta, gamma, psi,
&      print*, '01-', fgl      xold, p, u, pb, pbb, deltax,
&      print*, '02-', fg3      deltax, oldalphaf, alphaf, gammaf,
&      print*, '01-', fcl      oldpsif, oldpextraf, aa, olddeltax, oldgammaf,
&      print*, '02-', fe2      omega, charge2, sigma, epsilon, pi, f1, f2, g1, g2, f3,
&      print*, '03-', fgb      pua, pub, pta, ptb, pextra, abg,

```

```

&      if (case.eq.1) then      puaf, pubf, ptaf, ptbf, pextraf, psif,
&      call date(bu)          pua1, pub1, pta1, ptb1, pextra1, abg1,
&      call time(bu)          pua2, pub2, pta2, ptb2, pextra2, abg2,
&      print*, 'END'         pua3, pub3, pta3, ptb3, pextra3, abg3,
&      endif                 pe1, pe2,
&      if (case.eq.1) then    y(8), wk(232), masslow, massmed, masshigh, mass2,
&      (READ MASS OF GROUND   a, b, c, d, fg1, fe1, fg2, fe2, fg3,
&      print*, '(3'          zero, one, two, three,
&      open(unit=4,          a1, b1, c1, d1, w1, alpha1, beta1,
&      read(unit=4,         a2, b2, c2, d2, w2, alpha2, beta2,
&      read(unit=4,         a3, b3, c3, d3, w3, alpha3, beta3,
&      read(unit=4,         check_alpha
integer      read(un)      ij, ijlast, case, regime, inverse_step_size, nodes
logical      read(un)      okay, record, lastpos
logical      read(un)      warn1, warn2, warn3, warn4, warn5, warn6, warn7, warn8, warn9,
&      read(un)            warn10, warn11
character*12  buffer, buffer2, recfile, rec2
character*100 junk
common       a, b, c, d, fg1, fe1, fg2, fe2, fg3
common /const/ zero, one, two, three
common /case1/ a1, b1, c1, d1, w1, alpha1, beta1
common /case2/ a2, b2, c2, d2, w2, alpha2, beta2
common /case3/ a3, b3, c3, d3, w3, alpha3, beta3
save /case1/
save /case2/
save /case3/
external fcn

warn1=.true.
warn2=.true.
warn3=.true.
warn4=.true.
warn5=.true.
warn6=.true.
warn7=.true.
warn8=.true.
warn9=.true.
warn10=.true.
warn11=.true.

print*, ' '
print*, 'case=', case
if (case.ne.1 .and. case.ne.2 .and. case.ne.3) call exit(0)
print *, 'A=', a
print *, 'B=', b
print *, 'C=', c
print *, 'D=', d
print*, 'FG1=', fg1
print*, 'FG2=', fg2
print*, 'FE1=', fe1
print*, 'FE2=', fe2
print*, 'FG3=', fg3

```

```

if (case.eq.1) then
    call date(buffer)
    call time(buffer2)
    print*, 'ESCALINT3 run on ', buffer, ' at ', buffer2
endif
if (case.eq.1) then
C (READ MASS OF GROUND PARTICLE, AND CHARGE OF GROUP)
    print*, '(Reading GROUP.DAT)'
    open(unit=4, file='group.dat', status='old')
    read(unit=4, fmt=80), junk
    read(unit=4, fmt=95), masslow
    read(unit=4, fmt=95), massmed
    read(unit=4, fmt=95), masshigh
    read(unit=4, fmt=95), charge
    close(unit=4)
c (READ ESC_DO.DAT)
    print*, '(Reading ESC_DO.DAT)'
    open(unit=4, file='esc_do.dat', status='old')
    read(unit=4, fmt=*) record
    read(unit=4, fmt=*), regime
    read(unit=4, fmt=*), inverse_step_size
    close(unit=4)
    print*, 'Regime = ', regime
endif

if (record) then
    print*, '(Record is .true.)'
    if (case.eq.1) recfile='e31.dat'
    if (case.eq.2) recfile='e32.dat'
    if (case.eq.3) recfile='e33.dat'
    if (case.eq.1) rec2='e41.dat'
    if (case.eq.2) rec2='e42.dat'
    if (case.eq.3) rec2='e43.dat'
    open(unit=8, file=recfile, status='new')
    print*, 'Opening file: ', recfile
    open(unit=9, file=rec2, status='new')
    print*, 'Opening file: ', rec2
endif

c CALCULATE
x = zero
y(1) = a
y(2) = zero
y(3) = b
y(4) = zero
y(5) = c
y(6) = zero
y(7) = d
y(8) = zero
h = one/inverse_step_size
print*, 'step size =1/', inverse_step_size, ' = ', h
delta = zero

```

```

alpha = zero
gamma = zero
psi = zero
pta = zero
ptb = zero
pua = zero
pub = zero
pextra = zero
ijlast = 20*inverse_step_size
check_alpha = 999.9
c (OKAY STOPS US WANDERING INTO THE REGION WHERE THE FUNCTION BLOWS UP)
okay=.true.
c (LASTPOS RECORDS WHETHER ALPHA IS + OR -)
if (y(3).ge.zero) then
    lastpos = .true.
else
    lastpos = .false.
endif
c (NODES RECORDS HOW MANY TIMES ALPHA HAS CROSSED THE AXIS)
nodes = 0

print*, 'Integrating...'
ij= 1
c -----INTEGRATION LOOP-----
do 30 while (okay .and. ij.le.ijlast)
    xold = x
    call drkstp(8, h, x, y, fcn, wk)
c     PERFORM THE INTEGRATION OF ALPHA, DELTA, GAMMA, PSI,PTA,PTB,ETC
    p = y(1)
    u = y(3)
    pb = y(5)
    pbb = y(7)
    deltax=x-xold
    if ((abs(deltax-h)).gt.1.0d-15) print*, 'deltax,h=',deltax,h
    deltaf = p**2*u*x
    delta = delta + deltaf*deltax
    oldalphaf=alphaf
    alphaf = deltaf*x
    alpha = alpha + alphaf*deltax
c     HAS ALPHA CROSSED THE AXIS?
    if (lastpos .and. alpha.lt.zero) then
        lastpos = .false.
        nodes = nodes + 1
        print 292, 'Alpha has crossed from + to -. Number of nodes =',
            nodes, ', x = ', x
        format (' ', a, i2, a, f5.2)
        if (nodes.eq.regime) check_alpha = x + 2.0
    endif
    if (.not.lastpos .and. alpha.gt.zero) then
        lastpos = .true.

```

```

endif nodes = nodes + 1
print 292, 'Alpha has crossed from - to +. Number of nodes =',
& nodes, ', x = ', x
if (nodes.eq.regime) check_alpha = x + 2.0
endif
if (nodes.gt.regime) then
print*, 'Something is wrong. Nodes=', nodes, 'regime=', regime
stop
endif
c (STOP INTEGRATING WHEN ALPHA STARTS DIVERGING)
if (x.ge.check_alpha .and. abs(alphaf).gt.abs(olddalphaf) ) then
print*, 'Stopping. Alpha starts diverging at x=', x
print*, '(check_alpha = ', check_alpha, ')'
print*, 'Here y,y1,y2,z= ', y(1), y(5), y(7), y(3)
okay=.false.
endif

gammaaf = alphaf*u
gamma = gamma + gammaaf*deltax
if (fe1.ne.zero .and. fe2.ne.zero) then
& psif = fg1*x*x*pb*pb*(p*p+ pb*pb/(two*fe1)) +
fg2*x*x*pbb*pbb*(p*p+pbb*pbb/(two*fe2))
else
c (PSI NOT COMPUTED TO AVOID DIVISION BY ZERO ERROR)
psif= zero
endif
psi = psi + psif*deltax
ptaf = x*x*p*p*pb*pb
pta = pta+ ptaf*deltax
ptbf = x*x*p*p*pbb*pbb
ptb = ptb + ptbf*deltax
puaf = x*x*pb*pb*pb*pb/two
pua = pua + puaf*deltax
pubf = x*x*pbb*pbb*pbb*pbb/two
pub = pub + pubf*deltax
pextraf = pb*pb* pbb*pbb* x*x
pextra= pextra + pextraf*deltax
olddeltaf=deltaf
oldgammaf=gammaaf
oldpsif=psif
oldpextraf=pextraf

c WARNINGS
if (x.ge.check_alpha .and. abs(alphaf).gt.abs(olddalphaf) .and. warn1) then
warn1=.false.
print*, 'WARNING: alpha starts to diverge at x=',x
endif
c (DON'T EVER USE DELTA, SO DON'T REALLY CARE ABOUT IT.)
c if (x.gt.5 .and. abs(deltaf).gt.abs(olddeltaaf) .and. warn2) then
c warn2=.false.
c print*, 'WARNING: delta starts to diverge at x=',x

```

```

c          endif
c          CHECK: if (x.gt.5 .and. abs(gammaf).gt.abs(oldgammaf) .and. warn3) then
                warn3=.false.
                print*,'WARNING: gamma starts to diverge at x=',x
            endif
            &endif if (x.gt.5 .and. abs(psif).gt.abs(oldpsif) .and. warn4) then
                warn4=.false.
                print*,'WARNING: psi starts to diverge at x=',x
            endif
            &endif if (x.gt.5 .and. abs(pextraf).gt.abs(oldpextraf) .and. warn9) then
                warn9=.false.
                print*,'WARNING: psi starts to diverge at x=',x
            endif
c          PRINTO CHECK DERIVATIVES OF Y AND Z
            if (x.ge.check_alpha) then
                &print*,'ALPHA' if (lastpos .and. y(2).gt.zero .and. warn10) then
                    print*,'y''>0 at x=',x
                    warn10= .false.
                &endif
                &print*,'PSI' if (.not.lastpos .and. y(2).lt.zero .and. warn10) then
                    print*,'y''<0 at x=',x
                    warn10= .false.
                &endif
                &print*,'Z' if (lastpos .and. y(4).gt.zero .and. warn11) then
                    print*,'z''>0 at x=',x
                    warn11= .false.
                &endif
                &print*,'ALPHA' if (.not.lastpos .and. y(4).lt.zero .and. warn11) then
                    print*,'z''<0 at x=',x
                    warn11= .false.
                &endif
            endif
c          (SELECTIVE PRINTOUT CAN GO HERE IF DESIRED)
c          COMPUTE WRITE TO "E3i.DAT" IF WANTED.
c          (DO x=0 STUFF)
            if (record .and. ij.le.1) then
                &write(unit=8, fmt=654), zero, a, b, c, d,
                &zero, zero, zero, b
            endif
c          (WRITE OUT IN INCREMENTS OF 0.1 ALONG THE X AXIS)
            if (record .and. int(mod(ij, int(inverse_step_size*0.1))).eq.1 ) then
                &write(unit=8, fmt=654), x,y(1),y(3),y(5),y(7),
                &alpha,gamma,psi, y(3)- alpha/x
654          &format (f8.3, 1x, f8.3, 1x, f8.3, 1x, f8.3, 1x,
                &f8.3, 1x, f8.3, 1x, f8.3, 1x, f12.3, 1x, f12.3)
            endif
            ij=ij+1
30      continue
c -----(END OF INTEGRATION LOOP)-----

```

```

c      CHECK: DID Z CROSS ENOUGH TIMES?
      if (nodes.lt.regime) then
          print*, 'ERROR! Nodes = ', nodes, ' Regime = ', regime
          stop
      endif

c      CLOSE E3i.DAT FILE IF NECESSARY
      if (record) then
          endfile(unit=8)
          close(unit=8)
      endif

c      PRINTOUT
      print*, 'Stopped at X=', x
      print*, 'ALPHA=', alpha
      print*, 'GAMMA=', gamma
      print*, 'PSI=', psi

c      COMPUTE BETA
      beta=y(3)-(alpha/x)
      print*, 'BETA=', beta
      print*, 'z,x=', y(3), x

      print*, 'alpha,beta+gamma=', (alpha*beta+gamma)

c      CHECK ALPHA
      aa=(y(3) - beta)*x
c      print*, 'ALPHA should equal:', aa, ' Difference=', (alpha-aa)
      if (abs(aa-alpha) .gt. 1.0d-12) then
          print*, 'ERROR!!!! aa does not equal alpha!!!!'
          call exit(0)
      endif

c      COMPUTE OMEGA
      omega=psi - (alpha*beta + gamma)
      if (fg3.ne.zero) then
          omega = omega + (fg3/fe1)*pextra
      endif
      print*, 'OMEGA(w)=', omega
      if (case.eq.1) w1=omega
      if (case.eq.2) w2=omega
      if (case.eq.3) w3=omega

c      if (case.eq.1) then
c          COMPUTE SIGMA ETC
          print*, 'mass of lowest state=', masslow, ' MeV'
          print*, 'charge of group=', charge, ' (in units of e)'
          (convert mass from MeV to cm, scaled by 10^60)
          (1ev=1.324x10^-61cm [MTW], 1Mev=10^6ev)
          mass2 = masslow * (1.324d-1) * 1d6
c          (convert charge to cm, scaled by 10^30)

```

```

c      (e=1.381x10-34cm [MTW])
      charge2 = charge * (1.381d-4)
      print*, 'mass of ground state= ',mass2, ' x10-60 cm'
      print*, 'charge of ground state=',charge2,' x10-30 cm'

c      (equation 4.19, page 22)
c      (the scaling factors cancel out)
      sigma=(2.0d0 * alpha**2/omega)*(mass2/charge2**2)
      print*, 'SIGMA=', sigma, ' cm-1'
      print*, '1/SIGMA=', (1/sigma), ' cm'

c      COMPUTE EPSILON
c      (Scaled)
      epsilon=alpha/charge2
      print*, 'EPSILON=', epsilon, 'x1030 cm-1'

c      COMPUTE F1 ETC
      pi=3.14159265358979323846 d0
c      (SKIP THIS SECTION IF THERE IS NO COUPLING)
      if (fg1.ne.zero .and. fg2.ne.zero .and. fg3.ne.zero
&      .and. fe1.ne.zero .and. fe2.ne.zero) then
          f1 = (-fe1)*(4.0d0*pi*epsilon**2)
          f2 = (-fe2)*(4.0d0*pi*epsilon**2)
          g1 = f1/(-fg1)
          g2 = f2/(-fg2)
          f3 = -g2*(fg3/fg1)
          print99,'F1=', f1, 'x1060'
          print99,'F2=', f2, 'x1060'
          print99,'G1=', g1, 'x1060'
          print99,'G2=', g2, 'x1060'
          print99,'F3=', f3, 'x1060'
          if ( abs(-f1/g1 - (fg1) ).gt.1.0d-15) call exit(0)
          if ( abs(-f2/g2 - (fg2) ).gt.1.0d-15) call exit(0)
          if ( abs(-f3/g2 - (fg3/fg1) ).gt.1.0d-15) call exit(0)
          if ( abs(-f3/g1 - (fg3*fe2/(fg2*fe1) ) ).gt.1.0d-15) call exit(0)
      endif
endif

c      PRINT INTEGRALS TO SCREEN
      print*, 'pua=', pua
      print*, 'pub=', pub
      print*, 'pta=', pta
      print*, 'ptb=', ptb
      print*, 'pextra=', pextra

c      SAVE THE VARIABLES
      if (case.eq.1) then
          a1=a
          b1=b
          c1=c
          d1=d
          w1=omega

```

```

        alpha1=alpha
        beta1=beta
        pua1=pua
        pub1=pub
        pta1=pta
        ptb1=ptb
        pextra1=pextra
        abg1=alpha*beta + gamma
    endif

    if (case.eq.2) then
        a2=a
        b2=b
        c2=c
        d2=d
        w2=omega
        alpha2=alpha
        beta2=beta
        pua2=pua
        pub2=pub
        pta2=pta
        ptb2=ptb
        pextra2=pextra
        abg2=alpha*beta + gamma
    endif

    if (case.eq.3) then
        a3=a
        b3=b
        c3=c
        d3=d
        w3=omega
        alpha3=alpha
        beta3=beta
        pua3=pua
        pub3=pub
        pta3=pta
        ptb3=ptb
        pextra3=pextra
        abg3=alpha*beta + gamma
    endif

c    WRITE TO LONG FILE
    if (case.eq.1) then
        print*, '(Writing to ESCLONG.DAT)'
        open(unit=7, file='esclong.dat', status='new')
        write(unit=7, fmt=*)
        &         'Long output of ESCALINT3 run on ', buffer,
        &         ' at ', buffer2
        write(unit=7,fmt=92),'fg1=', fg1

```

```

          write(unit=7,fmt=92),'fg2=', fg2
          write(unit=7,fmt=92),'fe1=', fe1
          write(unit=7,fmt=92),'fe2=', fe2
          write(unit=7,fmt=92),'fg3=', fg3
        endif
c      (NEXT FIVE LINES ARE WRITTEN FOR ALL 3 LEVELS)
        write(unit=7,fmt=92),'pua=', pua
        write(unit=7,fmt=92),'pub=', pub
        write(unit=7,fmt=92),'pta=', pta
        write(unit=7,fmt=92),'ptb=', ptb
        write(unit=7,fmt=92),'pextra=', pextra
        write(unit=7,fmt=92),'alpha*beta+gamma=',
&      (alpha*beta + gamma)
        write(unit=7, fmt=92),'psi=', psi
        write(unit=7,fmt=92),'w=', omega
        if (case.ge.3) then
            endfile(unit=7)
            close(unit=7)
        endif

        if (case.eq.3) then
            print*, '(Writing to ESCSUM.DAT)'
            open(unit=6,file='escsum.dat',status='new')
            write(unit=6,fmt=*),'Summary of ESCALINT3 ',
&      buffer, ' ', buffer2
            write(unit=6, fmt=200), 'A=', a1, a2, a3
            write(unit=6, fmt=200), 'B=', b1, b2, b3
            write(unit=6, fmt=200), 'C=', c1, c2, c3
            write(unit=6, fmt=200), 'D=', d1, d2, d3
            write(unit=6, fmt=*),
&      '-----'
            write(unit=6, fmt=207), 'w=', w1, w2, w3
            write(unit=6, fmt=208), 'Alpha=', alpha1, alpha2, alpha3
            write(unit=6, fmt=207), 'Beta=', beta1, beta2, beta3
            write(unit=6, fmt=201), (fg1*1.0d4), (fe1*1.0d4)
            write(unit=6, fmt=202), (fg2*1.0d4), (fe2*1.0d4),
&      (fg3*1.0d4)
            write(unit=6, fmt=203), (w2/w1), (w3/w1)
            pe1= 100* ( ((w2/w1)-(massmed/masslow)) / (massmed/masslow) )
            pe2= 100* ( ((w3/w1)-(masshigh/masslow)) / (masshigh/masslow) )
            write(unit=6, fmt=209), pe1, pe2
            write(unit=6, fmt=*), ' '
            write(unit=6, fmt=*), 'pua, pub, pta, ptb, pextra, abg...'
            write(unit=6, fmt=901), '1', pua1, pub1, pta1,
&      ptb1, pextra1, abg1
            write(unit=6, fmt=901), '2', pua2, pub2, pta2,
&      ptb2, pextra2, abg2
            write(unit=6, fmt=901), '3', pua3, pub3, pta3,
&      ptb3, pextra3, abg3
            endfile(unit=6)
            close(unit=6)

```

```

endif
901 format ( ' ', a1, ': ', 6(f7.3, 1x) )
200 format ( ' ', a6, ' : ', 3(f17.10, ' : ') )
201 format ( ' ', 'FG1=', f12.3, 3x, 'FE1=', f12.3 , ' [x10^-4 for all F''s]' )
202 format ( ' ', 'FG2=', f12.3, 3x, 'FE2=', f12.3 ,
&          3x, 'FG3=', f12.3)
203 format ( ' ', 'w2/w1=', f12.4, 3x, 'w3/w1=', f12.4)
207 format ( ' ', a6, ' : ', 3(f10.3, 7x, ' : ') )
208 format ( ' ', a6, ' : ', 3(f12.5, 5x, ' : ') )
209 format ( ' ', 'error=' f12.4, '

c      WRITE TO E4i.DAT
      if (record) then
          do 345, x=0.2, 10, 0.2
              write(unit=9,fmt=*), x, (alpha/x + beta), beta
345      continue
c      CLOSE E4i.DAT FILE IF NECESSARY
      endfile(unit=9)
      close(unit=9)
endif

c      -----FORMATS-----
101 format ( ' ', 'A=', f9.5, ' B=', f9.5,
&          ' C=', f9.5, ' D=', f9.5)
109 format ( ' ', 'w=', f12.3, ' alpha=', f12.4,
&          ' psi=', e12.4)
102 format ( ' ', 'psi=',
&          '(, f10.2, '*', f5.2, ')',
&          ' + ',
&          '(, f10.2, '*', f5.2, ')',
&          ' + ',
&          '(, f10.2, '*', f5.2, ')',
&          ' + ',
&          '(, f10.2, '*', f5.2, ')')
103 format( ' ', ' = ',
&          f10.2, '+', f10.2, '+', f10.2, '+', f10.2, '.')
104 format( ' ', 'w= (, f10.2, ')-(, f10.2, '))
105 format( ' ', ' = ', f10.2, '.')

90 format (a, f12.4)
99 format ( ' ', a, e15.6, a)
92 format ( ' ', a, t20, e30.16)
80 format (a)
95 format (t20, e30.16)

c -----
      return
      end

c -----
c -----

```

```

c      FCN      (EVALUATES DERIVATIVES)
      subroutine fcn(x, y, yprime)
      real*8      y(8), yprime(8), p, q, u, v, pb, pbb, qb, qbb,
&              a, b, c, d, fg1, fe1, fg2, fe2, fg3,
&              zero, one, two, three, x
      common      a, b, c, d, fg1, fe1, fg2, fe2, fg3
      common /const/ zero, one, two, three

      p = y(1)
      q = y(2)
      u = y(3)
      v = y(4)
      pb = y(5)
      qb = y(6)
      pbb = y(7)
      qbb = y(8)
      if (x.gt.zero) then
         yprime(1) = q
         yprime(2) = (one-u**2)*p - 2*q/x +(pb**2)*p*fg1 + (pbb**2)*p*fg2
         yprime(3) = v
         yprime(4) = -u*p**2 - 2*v/x
         yprime(5) = qb
         yprime(6) = -2*qb/x + (one - pb*pb)*pb - p*p*pb*fe1
         if (fg3.ne.zero) then
            yprime(6)= yprime(6) - (fg3/fg1)*pbb*pbb*pb
         endif
         yprime(7) = qbb
         yprime(8) = -2*qbb/x + (one -pbb*pbb)*pbb - p*p*pbb*fe2
         if (fg3.ne.zero) then
            yprime(8) = yprime(8) - (fg3*fe2/(fg2*fe1))*pb*pb*pbb
         endif
      endif
      else
         yprime(1) = zero
         yprime(2) = (one-b*b+c*c*fg1+d*d*fg2)*a/three
         yprime(3) = zero
         yprime(4) = -b*a**2/three
         yprime(5) = zero
         yprime(6) = (one - a*a*fe1 -c*c )*c/three
         if (fg3.ne.zero) then
            yprime(6)= yprime(6) - (fg3/fg1)*d*d *c/three
         endif
         yprime(7) = zero
         yprime(8) = (one - a*a*fe2 -d*d )*d/three
         if (fg3.ne.zero) then
            yprime(8) = yprime(8) - (fg3*fe2/(fg2*fe1))*c*c *d/three
         endif
      endif
      return
      end

```

```

c -----
c -----

```

## 3) IMPROVE

```

c -----
c -----
c      IMPROVE
c      =====
c
c      Calculates better FG1, FG2 given PUA, etc.
c      Takes input from ESCLONG.DAT, GROUP.DAT.
c
c      Writes better FG's to file IMPROVE.DAT.
c
c      Written:
c      1988/Nov          -by Philip Sharman.
c      Revised:
c      1988/Dec/9       - bug fixed (w1 not constant)
c                      - FG3 and PEXTRA terms added.
c -----
c -----

```

```

      implicit real*8 (a-z)
      character*79   title
      character*12  buffer, buffer2

      call date(buffer)
      call time(buffer2)
      print*, 'IMPROVE run on ', buffer, ' at ', buffer2

```

```

c      READ DATA
      print*, 'Reading ESCLONG.DAT:'
      open(unit=3, file='esclong.dat', status='old')
      read(unit=3, fmt=90), title
      print*, title
      read(unit=3,fmt=93), fg1
      read(unit=3,fmt=93), fg2
      read(unit=3,fmt=93), fe1
      read(unit=3,fmt=93), fe2
      read(unit=3,fmt=93), fg3

      read(unit=3,fmt=93), pua1
      read(unit=3,fmt=93), pub1
      read(unit=3,fmt=93), pta1
      read(unit=3,fmt=93), ptb1
      read(unit=3,fmt=93), pextra1
      read(unit=3,fmt=93), abg1
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), w1

      read(unit=3,fmt=93), pua2
      read(unit=3,fmt=93), pub2
      read(unit=3,fmt=93), pta2
      read(unit=3,fmt=93), ptb2

```

```

read(unit=3,fmt=93), pextra2
read(unit=3,fmt=93), abg2
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), w2

read(unit=3,fmt=93), pua3
read(unit=3,fmt=93), pub3
read(unit=3,fmt=93), pta3
read(unit=3,fmt=93), ptb3
read(unit=3,fmt=93), pextra3
read(unit=3,fmt=93), abg3
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), w3

close(unit=3)

print*, '      pua      pub      pta      ptb      pextra      abg'
print 901, '1', pua1, pub1, pta1,
&          ptb1, pextra1, abg1
print 901, '2', pua2, pub2, pta2,
&          ptb2, pextra2, abg2
print 901, '3', pua3, pub3, pta3,
&          ptb3, pextra3, abg3

print*
print*, 'old fg1=', fg1
print*, 'old fg2=', fg2
print*, 'fe1=', fe1
print*, 'fe2=', fe2
print*, 'fg3=', fg3
901 format (' ', a1, ': ', 6(f7.3, 1x) )

print*, '(Reading GOUP.DAT)'
open(unit=3, file='group.dat', status='old')
read(unit=3, fmt=90), title
read(unit=3, fmt=93), mass1
read(unit=3, fmt=93), mass2
read(unit=3, fmt=93), mass3
close(unit=3)

c -----
c (NEW VARIABLE NAMES)
b1=(mass2/mass1)
b2=(mass3/mass1)
print*, 'desired wm/we=', b1
print*, 'desired wt/we=', b2

a=      pua1
b=      pub1
c=      pta1
d=      ptb1
e=      abg1 - (fg3/fe1)*pextra1

```

```

if (abs(lhs-rhs).gt.criteria) then
f=      pua2
g=      pub2
h=      pta2
j=      ptb2
k=      abg2 - (fg3/fe1)*pextra2

      null

l=      pua3
m=      pub3
n=      pta3
p=      ptb3
q=      abg3 - (fg3/fe1)*pextra3

ccc      print 901, '1', a,b,c,d,e, w1
ccc      print 901, '2', f,g,h,j,k, w2
ccc      print 901, '3', l,m,n,p,q, w3

c1= (b1*a - f)/fe1 + b1*c - h
c2= (b1*b - g)/fe2 + b1*d - j
c3= b1*e - k

c4= (b2*a - l)/fe1 + b2*c - n
c5= (b2*b - m)/fe2 + b2*d - p
c6= b2*e - q

c -----
c      CALCULATE
c      Solve: xC1 + yC2 = C3
c              xC4 + yC5 = C6

c      x= (c3*c5 - c6*c2)/(c1*c5 - c4*c2)
c      y= (c3*c4 - c6*c1)/(c2*c4 - c5*c1)
c      print*, 'Perfect fg1=', x
c      print*, 'Perfect fg2=', y

c      WRITE TO OUTPUT FILE
c      print*, '(Writing to IMPROVE.DAT)'
c      open(unit=3,file='improve.dat',status='new')
c      write(unit=3,fmt=*) , 'IMPROVE run on ',
&          buffer, ' at ', buffer2
c      write(unit=3,fmt=80), 'Perfect fg1=', x
c      write(unit=3,fmt=80), 'Perfect fg2=', y
80     format (' ', a, t20, e30.16)

c      CHECK
c      (lhs should equal rhs in all these checks)
c      criteria= 1.0 d-10
c      (error criteria)

lhs= x*c1 + y*c2
rhs= c3

```

```

if (abs(lhs-rhs).gt.criteria) then
    print*,'**** WARNING! *** Check #1 fails'
    print*,'Difference = ', (lhs-rhs)
    write(unit=3,fmt=*),'**** WARNING! *** Check #1 fails'
    print*,'lhs rhs=', lhs, rhs
    write(unit=3,fmt=*),'Difference = ', (lhs-rhs)
endif

lhs= x*c4 + y*c5
rhs= c6
if (abs(lhs-rhs).gt.criteria) then
    print*,'**** WARNING! *** Check #2 fails'
    print*,'Difference = ', (lhs-rhs)
    print*,'lhs rhs=', lhs, rhs
    write(unit=3,fmt=*),'**** WARNING! *** Check #2 fails'
    write(unit=3,fmt=*),'Difference = ', (lhs-rhs)
endif

c      (use the old fg's)
lhs=w1
rhs=(fg1*pua1)/fe1 + (fg2*pub1)/fe2
&      + fg1*pta1 + fg2*ptb1 - (abg1 - (fg3/fe1)*pextra1)
if (abs(lhs-rhs).gt.criteria) then
    print*,'**** WARNING! *** Check #3 fails'
    print*,'Difference = ', (lhs-rhs)
    print*,'lhs rhs=', lhs, rhs
    write(unit=3,fmt=*),'**** WARNING! *** Check #3 fails'
    write(unit=3,fmt=*),'Difference = ', (lhs-rhs)
endif

c      (use the new fg's)
fg1=x
fg2=y
c      (figure out what w1 is now with the new fg's)
w1new=(fg1*pua1)/fe1 + (fg2*pub1)/fe2
&      + fg1*pta1 + fg2*ptb1 - (abg1 - (fg3/fe1)*pextra1)

lhs=(b1*w1new)
rhs=(fg1*pua2)/fe1 + (fg2*pub2)/fe2
&      + fg1*pta2 + fg2*ptb2 - (abg2 - (fg3/fe1)*pextra2)
if (abs(lhs-rhs).gt.criteria) then
    print*,'**** WARNING! *** Check #4 fails'
    print*,'Difference = ', (lhs-rhs)
    print*,'lhs rhs=', lhs, rhs
    write(unit=3,fmt=*),'**** WARNING! *** Check #4 fails'
    write(unit=3,fmt=*),'Difference = ', (lhs-rhs)
endif

lhs=(b2*w1new)
rhs=(fg1*pua3)/fe1 + (fg2*pub3)/fe2

```

```

& 4) TRANS + fg1*pta3 + fg2*ptb3 - (abg3 - (fg3/fe1)*pextra3)
if (abs(lhs-rhs).gt.criteria) then
    print*,'**** WARNING! *** Check #5 fails'
    print*,'Difference = ', (lhs-rhs)
    print*,'lhs rhs=', lhs, rhs
    write(unit=3,fmt=*),'**** WARNING! *** Check #5 fails'
    write(unit=3,fmt=*),'Difference = ', (lhs-rhs)
endif
print*,'Checks are done.'

write(unit=3,fmt=*), 'Checks are done.'
endfile(unit=3)
close(unit=3)

-----
c
90 format (a)
93 format ( t20, e30.16)
print*,'Program ends.'
end

call date(buffer)
call time(buffer2)

READ THE OLD VALUES
open(unit=3, file='constants.dat', status='old')
read(unit=3, fmt=90), title
read(unit=3, fmt=91), junk
read(unit=3, fmt=91), junk
read(unit=3, fmt=91), fe1
read(unit=3, fmt=91), fe2
read(unit=3, fmt=91), fg3
close(unit=3)

90 format (a)
91 format (t20, e30.16)

READ THE NEW VALUES
open(unit=3, file='improve.dat', status='old')
read(unit=3, fmt=90), character_junk
read(unit=3, fmt=91), fg1
read(unit=3, fmt=91), fg2
close(unit=3)

WRITE THE FILE BACK OUT
open(unit=3, file='constants.dat', status='new')
write(unit=3, fmt=90), title
write(unit=3, fmt=91), 'fg1', fg1
write(unit=3, fmt=91), 'fg2', fg2
write(unit=3, fmt=91), 'fe1', fe1
write(unit=3, fmt=91), 'fe2', fe2

```

## 4) TRANSFER

```

-----
c
c
c   TRANSFER
c   =====
c
c   Takes "perfect" FG's from IMPROVE.DAT and inserts them
c   into FCONSTANTS.DAT.
c
c   Written:
c   1988/Dec/10   -by Philip Sharman.
c   Revised:
c   1989/Jan/19   -keep original title; add comment at bottom
c                   of output file.
c                   Jun/12   -FG3 term added.
c
-----
c
c   implicit      real*8 (a-z)
c   character*40   title
c   character*100  character_junk
c   character*12   buffer, buffer2
c
c   call date(buffer)
c   call time(buffer2)
c
c   READ THE OLD VALUES
c   open(unit=3, file='fconstants.dat', status='old')
c   read(unit=3, fmt=90), title
c   read(unit=3, fmt=91), junk
c   read(unit=3, fmt=91), junk
c   read(unit=3, fmt=91), fe1
c   read(unit=3, fmt=91), fe2
c   read(unit=3, fmt=91), fg3
c   close(unit=3)
90  format (a)
91  format (t20, e30.16)
c
c   READ THE NEW VALUES
c   open(unit=3, file='improve.dat', status='old')
c   read(unit=3, fmt=90), character_junk
c   read(unit=3, fmt=91), fg1
c   read(unit=3, fmt=91), fg2
c   close(unit=3)
c
c   WRITE THE FILE BACK OUT
c   open(unit=3, file='fconstants.dat', status='new')
c   write(unit=3, fmt=*, title
c   write(unit=3, fmt=81), 'fg1=', fg1
c   write(unit=3, fmt=81), 'fg2=', fg2
c   write(unit=3, fmt=81), 'fe1=', fe1
c   write(unit=3, fmt=81), 'fe2=', fe2

```

```

write(unit=3, fmt=81), 'fg3=', fg3
write(unit=3, fmt=*) , 'Altered by TRANSFER on ',
&      buffer, ' at ', buffer2
-----
endfile(unit=3)
close(unit=3)
81  format (' ', a, t20, e30.16)
-----
print*, '"Perfect" F's have been copied from IMPROVE.DAT to FCONSTANTS.DAT.'
-----
! See only the 2 quark groups.
end
! is included.

(See Basic Scientific Subroutines, Vol 3, page 80)

Reads:  ESC_INT.DAT, ESC_UCT.DAT, ESC_DSB.DAT,
        FCONSTANTS.DAT
Writes: PARINP.DAT

Written:
1989/Jan/16      (by Philip Sharman)

Revised:
1989/Jan/16      (changed from BASIC to FORTRAN)
17              (added input/output)
19              (keep original title; add comment to bottom)
20              (character_junk fixed to absorb unwanted titles)
28              (REAL changed to REAL*8 or DIMENSION)
28              (printout scaled. And aborts if iterations too big)
Feb/16          (FSZ term added, and DSB added.)
21              (Top quark equation deactivated.)
21              (LT changed to LE when making w's positive)
23              (Read mass ratios from files.)
Mar/5           (Tried not making negative w's positive)
14              (Changed to make sure initial guesses are over-estimated)
24              (Bug changed: had wrong chara/up desired mass ratios)
28              (Tried making the initial guesses all -1)
28              (Tried making the initial guesses all -300d-4)
May/22         (EQ changed to LE in RESIDUALS)
               (check for proper xx added in FVWC)
-----
-----

implicit real*8(a-z)
common L, n, x, y, s, e1, d, m, s01
common /parms/ a
common /fs/ fg1_initial_guess, fg2_initial_guess,
&           fs1_initial_guess, fs2_initial_guess,
&           fg3_initial_guess
common /title/ fs_title
integer n, L, I, m
dimension      x(8), y(8), s(8), e1(8)
character*80   fs_title
character*12   buffer, buffer2

```

## 5) PARIMP

```

c -----
c -----
c
c PARIMP.FOR SECTION ---
c =====
c This is "Parafit" modified to improve F's.
c Does only the 2 quark groups.
c FG3 is included.
c
c (See Basic Scientific Subroutines, Vol 2, page 80)
c
c Reads:  ESC_EMT.DAT, ESC_UCT.DAT, ESC_DSB.DAT,
c         FCONSTANTS.DAT.
c Writes: PARIMP.DAT.
c
c Written:
c 1989/Jan/16      (by Philip Sharman)
c
c Revised:
c 1989/Jan/16      (changed from BASIC to FORTRAN)
c 17              (added input/output)
c 19              (keep original title; add comment to bottom)
c 26              (character_junk fixed to absorb unwanted titles)
c 28              (REAL changed to REAL*8 or DIMENSION)
c 28              (printout scaled. And aborts if #iterations too big)
c Feb/18          (FG3 term added, and DSB added.)
c 21              (Top quark equation deactivated.)
c 21              (LT changed to LE when making w's positive)
c 23              (Read mass ratios from files.)
c Mar/8           (Tried not making negative w's positive)
c 14              (Changed to make sure initial guess are over-estimates)
c 24              (Bug changed: had wrong charm/up desired mass ratio)
c 28              (Tried making the initial guesses all -1)
c 28              (Tried making the initial guesses all -200d-4)
c May/22          (EQ changed to LE in RESIDUALS)
c                (check for proper xx added in FUNC)
c -----
c -----
c
c format (a)
c
c implicit real*8(a-z)
c common L, n, x, y, e, e1, d, m, ee1
c common /params/ a
c common /fs/ fg1_initial_guess, fg2_initial_guess,
c &          fe1_initial_guess, fe2_initial_guess,
c &          fg3_initial_guess
c common /title/ fs_title
c integer n, L, I, m
c dimension      x(6), y(6), a(5), e1(6)
c character*80   fs_title
c character*12   buffer, buffer2

```

```

call date(buffer)
call time(buffer2)
print*, 'PARIMP run on ', buffer, ' at ', buffer2

c      !---DRIVER SECTION---
n=6
c      !(number of equations)
L=5
c      !(number of parameters)

c      READ IN pua's ETC
call readstuff

c      ! (DATA)
x(1)=1
x(2)=2
x(3)=3
x(4)=4
x(5)=5
x(6)=6
c      READ MASS DATA
open(unit=3, file='emt.dat', status='old')
read(unit=3, fmt=90), title
read(unit=3, fmt=93), mass_1_emt
read(unit=3, fmt=93), mass_2_emt
read(unit=3, fmt=93), mass_3_emt
close(unit=3)
open(unit=3, file='uct.dat', status='old')
read(unit=3, fmt=90), title
read(unit=3, fmt=93), mass_1_uct
read(unit=3, fmt=93), mass_2_uct
read(unit=3, fmt=93), mass_3_uct
close(unit=3)
open(unit=3, file='dsb.dat', status='old')
read(unit=3, fmt=90), title
read(unit=3, fmt=93), mass_1_dsb
read(unit=3, fmt=93), mass_2_dsb
read(unit=3, fmt=93), mass_3_dsb
close(unit=3)
90      format (a)
93      format ( t20, e30.16)

print*, 'uct masses:', mass_1_uct, mass_2_uct, mass_3_uct
print*, 'dsb masses:', mass_1_dsb, mass_2_dsb, mass_3_dsb

cccc   y(1)=mass_2_emt/mass_1_emt
cccc   y(2)=mass_3_emt/mass_1_emt
y(1)=1.0
y(2)=2.0
y(3)=mass_2_uct/mass_1_uct
y(4)=mass_3_uct/mass_1_uct

```

```

y(5)=mass_2_dsb/mass_1_dsb
y(6)=mass_3_dsb/mass_1_dsb

print*, 'Input data:'
print*, '#      x      y'
do 111, I=1, n
    print*, I, x(I), y(I)
111  continue

e= 1e-6
ee1= 0.9
c  SHOW THE OLD VALUES
a(1)= fg1_initial_guess
a(2)= fg2_initial_guess
a(3)= fe1_initial_guess
a(4)= fe2_initial_guess
a(5)= fg3_initial_guess
print*, 'The old values are...'
print*, 'FG1= ', a(1), ' = ', a(1)*1.0d4, ' x10^-4'
print*, 'FG2= ', a(2), ' = ', a(2)*1.0d4, ' x10^-4'
print*, 'FE1= ', a(3), ' = ', a(3)*1.0d4, ' x10^-4'
print*, 'FE2= ', a(4), ' = ', a(4)*1.0d4, ' x10^-4'
print*, 'FG3= ', a(5), ' = ', a(5)*1.0d4, ' x10^-4'
print*, 'With the old values...'
print*, 'x      Calculated',
&      '      Desired mass ratio'
do 224, I=1, n
    xx=x(I)
    call func(xx,yy)
    print*, x(I), yy, y(I)
224  continue
print*, ' '

c  STARTING GUESSES
c  a(1)= 10.0 * fg1_initial_guess
c  a(2)= 10.0 * fg2_initial_guess
c  a(3)= 10.0 * fe1_initial_guess
c  a(4)= 10.0 * fe2_initial_guess
c  a(5)= 10.0 * fg3_initial_guess

do 135, i=1, 5
    a(i)= -200.0 d-4
135  continue

c  (CHANGE ANY INITIAL GUESS OF ZERO TO SOMETHING ELSE SO PARAFIT WILL WORK)
do 3, i=1, 5
    if (a(i).eq.0.0d0) a(i)= -1.0d-4
3  continue
print*, 'The initial guesses are...'
print*, (a(i), i=1,5)

call parafit

```

```

print*, 'Solving for only the 2 quark groups....'

print*, ' '
print*, 'The coefficients found by PARAFIT are:'
print*, 'FG1= ', a(1), ' = ', a(1)*1.0d4, ' x10^-4'
print*, 'FG2= ', a(2), ' = ', a(2)*1.0d4, ' x10^-4'
print*, 'FE1= ', a(3), ' = ', a(3)*1.0d4, ' x10^-4'
print*, 'FE2= ', a(4), ' = ', a(4)*1.0d4, ' x10^-4'
print*, 'FG3= ', a(5), ' = ', a(5)*1.0d4, ' x10^-4'
print*, ' '
print*, 'The standard deviation of the fit is:', d
print*, 'The number of iterations was: ', m
print*, ' '
print*, 'With these parameters...'
print*, 'x                               Calculated',
&      'Desired mass ratio'
do 124, I=1, n
  xx=x(I)
  call func(xx,yy)
  print*, x(I), yy, y(I)
124 continue

c      WRITE TO OUTPUT FILE
open(unit=3,file='parimp.dat',status='new')
write(unit=3,fmt=*), fs_title
write(unit=3,fmt=80), 'fg1=', a(1)
write(unit=3,fmt=80), 'fg2=', a(2)
write(unit=3,fmt=80), 'fe1=', a(3)
write(unit=3,fmt=80), 'fe2=', a(4)
write(unit=3,fmt=80), 'fg3=', a(5)
write(unit=3,fmt=*), 'Fs found by PARIMP on ',
&      buffer, ' at ', buffer2
80  format (' ', a, t20, e30.16)
endfile(unit=3)
close(unit=3)

end

-----
c
c
c      !---THE PARAFIT SUBROUTINE---
subroutine parafit
implicit real*8(a-z)
common L, n, x, y, e, e1, d, m, ee1
common /params/ a
integer n, L, I, mod_iterations, m
dimension      x(6), y(6), a(5), e1(6)

do 333, I=1, L
  e1(I)=ee1
128  continue
333  continue

```

```

m=0
L1=1000000
! SWEEP
4257 do 129, I=1, L
      a0=a(I)
      a(I)=a0
      (Why is that line there?)
      ! MIDDLE
4261 call residuals(L2)
      m0=L2
      a(I)=a0*(1-e1(I))
      call residuals(L2)
      m1=L2

      if (m1.gt.m0) e1(I)= -e1(I)/2
      if (m1.lt.m0) e1(I)=1.2*e1(I)
      if (m1.gt.m0) a(I)=a0
      if (m1.gt.m0) goto 4261
129  continue

      m=m+1
      mod_iterations= mod(m, 500)
      if (mod_iterations.eq.0) print*,'# of iterations= ', m
      if (m.gt.10000) call exit(0)

      if (L2.eq.0) return
      if (abs((L1-L2)/L2).lt.e) return

      L1=L2
      goto 4257
c    !(goto Sweep)
      end

c -----
c    ! RESIDUALS
c    subroutine residuals(L2)
c    (Returns L2, and modifies d [in common block].)
c    implicit real*8(a-z)
c    common L, n, x, y, e, e1, d, m, ee1
c    common /params/ a
c    integer n, L, j, m
c    dimension      x(6), y(6), a(5), e1(6)

      L2=0
      do 128,j=1, n
          xx=x(j)
          call func(xx, yy)
          L2=L2+(y(j)-yy)*(y(j)-yy)
128  continue

```

```

      if (n.gt.L) d=sqrt(L2/(n-L))
      if (n.le.L) d=sqrt(L2)
c      !(fudged. Not in original pgm.)
      return
      end
c -----
c      READ IN THE PUA'S ETC
      subroutine readstuff
      implicit real*8(a-z)
      common /con_1_empt/ pua_1_empt, pub_1_empt,
&          pta_1_empt, ptb_1_empt, abg_1_empt, pextra_1_empt
      common /con_2_empt/ pua_2_empt, pub_2_empt,
&          pta_2_empt, ptb_2_empt, abg_2_empt, pextra_2_empt
      common /con_3_empt/ pua_3_empt, pub_3_empt,
&          pta_3_empt, ptb_3_empt, abg_3_empt, pextra_3_empt

      common /con_1_uct/ pua_1_uct, pub_1_uct,
&          pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct
      common /con_2_uct/ pua_2_uct, pub_2_uct,
&          pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct
      common /con_3_uct/ pua_3_uct, pub_3_uct,
&          pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct

      common /con_1_dsb/ pua_1_dsb, pub_1_dsb,
&          pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb
      common /con_2_dsb/ pua_2_dsb, pub_2_dsb,
&          pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb
      common /con_3_dsb/ pua_3_dsb, pub_3_dsb,
&          pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb

      common /fs/ fg1_initial_guess, fg2_initial_guess,
&          fe1_initial_guess, fe2_initial_guess,
&          fg3_initial_guess

      common /title/ fs_title
      character*80 fs_title
      character*100 character_junk

      save /con_1_empt/, /con_2_empt/, /con_3_empt/
      save /con_1_uct/, /con_2_uct/, /con_3_uct/
      save /con_1_dsb/, /con_2_dsb/, /con_3_dsb/
      save /title/

c      READ F'S
      open(unit=3, file='fconstants.dat', status='old')
      read(unit=3, fmt=90), fs_title
      read(unit=3, fmt=93), fg1_initial_guess
      read(unit=3, fmt=93), fg2_initial_guess
      read(unit=3, fmt=93), fe1_initial_guess
      read(unit=3, fmt=93), fe2_initial_guess
      read(unit=3, fmt=93), fg3_initial_guess

```

```

close(unit=3)

c READ EMT DATA
open(unit=3, file='esc_empt.dat', access='sequential',
&      form='formatted', status='old')
read(unit=3, fmt=90), character_junk
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk

read(unit=3, fmt=93), pua_1_empt
read(unit=3, fmt=93), pub_1_empt
read(unit=3, fmt=93), pta_1_empt
read(unit=3, fmt=93), ptb_1_empt
read(unit=3, fmt=93), pextra_1_empt
read(unit=3, fmt=93), abg_1_empt
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk

read(unit=3, fmt=93), pua_2_empt
read(unit=3, fmt=93), pub_2_empt
read(unit=3, fmt=93), pta_2_empt
read(unit=3, fmt=93), ptb_2_empt
read(unit=3, fmt=93), pextra_2_empt
read(unit=3, fmt=93), abg_2_empt
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk

read(unit=3, fmt=93), pua_3_empt
read(unit=3, fmt=93), pub_3_empt
read(unit=3, fmt=93), pta_3_empt
read(unit=3, fmt=93), ptb_3_empt
read(unit=3, fmt=93), pextra_3_empt
read(unit=3, fmt=93), abg_3_empt
read(unit=3, fmt=93), junk
read(unit=3, fmt=93), junk

close(unit=3)

print*, 'EMT GROUP:'
print*, '  pua, pub, pta, ptb, pextra, abg ...'
print 901, '1', pua_1_empt, pub_1_empt, pta_1_empt,
&          ptb_1_empt, pextra_1_empt, abg_1_empt
print 901, '2', pua_2_empt, pub_2_empt, pta_2_empt,
&          ptb_2_empt, pextra_2_empt, abg_2_empt
print 901, '3', pua_3_empt, pub_3_empt, pta_3_empt,
&          ptb_3_empt, pextra_3_empt, abg_3_empt

901 format (' ', a1, ': ',

```

```

&          6(f7.3, 1x) )
c      READ UCT DATA
      open(unit=3, file='esc_uct.dat', access='sequential',
&          form='formatted', status='old')
      read(unit=3, fmt=90), character_junk
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk

      read(unit=3,fmt=93), pua_1_uct
      read(unit=3,fmt=93), pub_1_uct
      read(unit=3,fmt=93), pta_1_uct
      read(unit=3,fmt=93), ptb_1_uct
      read(unit=3,fmt=93), pextra_1_uct
      read(unit=3,fmt=93), abg_1_uct
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk

      read(unit=3,fmt=93), pua_2_uct
      read(unit=3,fmt=93), pub_2_uct
      read(unit=3,fmt=93), pta_2_uct
      read(unit=3,fmt=93), ptb_2_uct
      read(unit=3,fmt=93), pextra_2_uct
      read(unit=3,fmt=93), abg_2_uct
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk

      read(unit=3,fmt=93), pua_3_uct
      read(unit=3,fmt=93), pub_3_uct
      read(unit=3,fmt=93), pta_3_uct
      read(unit=3,fmt=93), ptb_3_uct
      read(unit=3,fmt=93), pextra_3_uct
      read(unit=3,fmt=93), abg_3_uct
      read(unit=3,fmt=93), junk
      read(unit=3,fmt=93), junk

      close(unit=3)

      print*, 'UCT GROUP:'
      print*, '  pua, pub, pta, ptb, pextra, abg ...'
      print 901, '1', pua_1_uct, pub_1_uct, pta_1_uct,
&          ptb_1_uct, pextra_1_uct, abg_1_uct
      print 901, '2', pua_2_uct, pub_2_uct, pta_2_uct,
&          ptb_2_uct, pextra_2_uct, abg_2_uct
      print 901, '3', pua_3_uct, pub_3_uct, pta_3_uct,
&          ptb_3_uct, pextra_3_uct, abg_3_uct

c      READ DSB DATA
      open(unit=3, file='esc_dsb.dat', access='sequential',

```

```

&      form='formatted', status='old')
c      read(unit=3,fmt=90), character_junk
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk

c      read(unit=3,fmt=93), pua_1_dsb
c      read(unit=3,fmt=93), pub_1_dsb
c      read(unit=3,fmt=93), pta_1_dsb
c      read(unit=3,fmt=93), ptb_1_dsb
c      read(unit=3,fmt=93), pextra_1_dsb
c      read(unit=3,fmt=93), abg_1_dsb
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk

c      read(unit=3,fmt=93), pua_2_dsb
c      read(unit=3,fmt=93), pub_2_dsb
c      read(unit=3,fmt=93), pta_2_dsb
c      read(unit=3,fmt=93), ptb_2_dsb
c      read(unit=3,fmt=93), pextra_2_dsb
c      read(unit=3,fmt=93), abg_2_dsb
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk

c      read(unit=3,fmt=93), pua_3_dsb
c      read(unit=3,fmt=93), pub_3_dsb
c      read(unit=3,fmt=93), pta_3_dsb
c      read(unit=3,fmt=93), ptb_3_dsb
c      read(unit=3,fmt=93), pextra_3_dsb
c      read(unit=3,fmt=93), abg_3_dsb
c      read(unit=3,fmt=93), junk
c      read(unit=3,fmt=93), junk

c      close(unit=3)

c      print*, 'DSB GROUP:'
c      print*, '   pua, pub, pta, ptb, pextra, abg   ...'
c      print 901, '1', pua_1_dsb, pub_1_dsb, pta_1_dsb,
&                ptb_1_dsb, pextra_1_dsb, abg_1_dsb
c      print 901, '2', pua_2_dsb, pub_2_dsb, pta_2_dsb,
&                ptb_2_dsb, pextra_2_dsb, abg_2_dsb
c      print 901, '3', pua_3_dsb, pub_3_dsb, pta_3_dsb,
&                ptb_3_dsb, pextra_3_dsb, abg_3_dsb

c      -----
c      format (a)
c      format ( t20, e30.16)
c      -----

c      return

```

```

end
c -----
c -----
c ! ---THE FUNCTION IN QUESTION---
subroutine func(xx,yy)
implicit real*8(a-z)
common /params/ a

common /con_1_empt/ pua_1_empt, pub_1_empt,
& pta_1_empt, ptb_1_empt, abg_1_empt, pextra_1_empt
common /con_2_empt/ pua_2_empt, pub_2_empt,
& pta_2_empt, ptb_2_empt, abg_2_empt, pextra_2_empt
common /con_3_empt/ pua_3_empt, pub_3_empt,
& pta_3_empt, ptb_3_empt, abg_3_empt, pextra_3_empt

common /con_1_uct/ pua_1_uct, pub_1_uct,
& pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct
common /con_2_uct/ pua_2_uct, pub_2_uct,
& pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct
common /con_3_uct/ pua_3_uct, pub_3_uct,
& pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct

common /con_1_dsb/ pua_1_dsb, pub_1_dsb,
& pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb
common /con_2_dsb/ pua_2_dsb, pub_2_dsb,
& pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb
common /con_3_dsb/ pua_3_dsb, pub_3_dsb,
& pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb

real*8 a(5), xx, yy

save /con_1_empt/, /con_2_empt/, /con_3_empt/
save /con_1_uct/, /con_2_uct/, /con_3_uct/
save /con_1_dsb/, /con_2_dsb/, /con_3_dsb/

fg1=a(1)
fg2=a(2)
fe1=a(3)
fe2=a(4)
fg3=a(5)

c CALCULATE
w1_empt= (fg1/fe1)*pua_1_empt + (fg2/fe2)*pub_1_empt
& + fg1*pta_1_empt + fg2*ptb_1_empt - abg_1_empt
& + (fg3/fe1)*pextra_1_empt
w2_empt= (fg1/fe1)*pua_2_empt + (fg2/fe2)*pub_2_empt
& + fg1*pta_2_empt + fg2*ptb_2_empt - abg_2_empt
& + (fg3/fe1)*pextra_2_empt
w3_empt= (fg1/fe1)*pua_3_empt + (fg2/fe2)*pub_3_empt
& + fg1*pta_3_empt + fg2*ptb_3_empt - abg_3_empt

```

```

&      + (fg3/fe1)*pextra_3_empt
cccc
w1_uct= (fg1/fe1)*pua_1_uct + (fg2/fe2)*pub_1_uct
&      + fg1*pta_1_uct + fg2*ptb_1_uct - abg_1_uct
&      + (fg3/fe1)*pextra_1_uct
w2_uct= (fg1/fe1)*pua_2_uct + (fg2/fe2)*pub_2_uct
&      + fg1*pta_2_uct + fg2*ptb_2_uct - abg_2_uct
&      + (fg3/fe1)*pextra_2_uct
w3_uct= (fg1/fe1)*pua_3_uct + (fg2/fe2)*pub_3_uct
&      + fg1*pta_3_uct + fg2*ptb_3_uct - abg_3_uct
&      + (fg3/fe1)*pextra_3_uct

w1_dsb= (fg1/fe1)*pua_1_dsb + (fg2/fe2)*pub_1_dsb
&      + fg1*pta_1_dsb + fg2*ptb_1_dsb - abg_1_dsb
&      + (fg3/fe1)*pextra_1_dsb
w2_dsb= (fg1/fe1)*pua_2_dsb + (fg2/fe2)*pub_2_dsb
&      + fg1*pta_2_dsb + fg2*ptb_2_dsb - abg_2_dsb
&      + (fg3/fe1)*pextra_2_dsb
w3_dsb= (fg1/fe1)*pua_3_dsb + (fg2/fe2)*pub_3_dsb
&      + fg1*pta_3_dsb + fg2*ptb_3_dsb - abg_3_dsb
&      + (fg3/fe1)*pextra_3_dsb

c      FUDGE ... IF W'S ARE ZERO OR NEGATIVE, MAKE THEM VERY SMALL BUT POSITIVE
cccc   if (w1_empt.le.0) w1_empt=0.0001*exp(w1_empt)
cccc   if (w2_empt.le.0) w2_empt=0.0001*exp(w2_empt)
cccc   if (w3_empt.le.0) w3_empt=0.0001*exp(w3_empt)
cccc   if (w1_uct.le.0) w1_uct=0.0001*exp(w1_uct)
cccc   if (w2_uct.le.0) w2_uct=0.0001*exp(w2_uct)
cccc   if (w3_uct.le.0) w3_uct=0.0001*exp(w3_uct)
cccc   if (w1_dsb.le.0) w1_dsb=0.0001*exp(w1_dsb)
cccc   if (w2_dsb.le.0) w2_dsb=0.0001*exp(w2_dsb)
cccc   if (w3_dsb.le.0) w3_dsb=0.0001*exp(w3_dsb)
c      TRY A DIFFERENT METHOD
      if (w1_empt.le.0) then
          w1_empt=0.00001
          w2_empt=9.0d4 *exp(abs(w1_empt))
          w3_empt=9.0d4 *exp(abs(w1_empt))
      endif
      if (w1_uct.le.0) then
          w1_uct=0.00001
          w2_uct=9.0d4 *exp(abs(w1_uct))
          w3_uct=9.0d4 *exp(abs(w1_uct))
      endif
      if (w1_dsb.le.0) then
          w1_dsb=0.00001
          w2_dsb=9.0d4 *exp(abs(w1_dsb))
          w3_dsb=9.0d4 *exp(abs(w1_dsb))
      endif

c      RETURN VALUES

```

```

cccc  if (xx.eq.1) yy=(w2_empt/w1_empt)
cccc  if (xx.eq.2) yy=(w3_empt/w1_empt)
c     if (xx.eq.1) yy=1.0
c     if (xx.eq.2) yy=2.0
c     if (xx.eq.3) yy=(w2_uct/w1_uct)
c     if (xx.eq.4) yy=(w3_uct/w1_uct)
c     if (xx.eq.5) yy=(w2_dsb/w1_dsb)
c     if (xx.eq.6) yy=(w3_dsb/w1_dsb)
c     if (xx.ne.1 .and. xx.ne.2 .and. xx.ne.3 .and. xx.ne.4
&      .and. xx.ne.5 .and. xx.ne.6) call exit(0)
c     write(*,*)
c     return
c     -----
c     -----
c     end
c     -----
c     implicit real*8 (a-z)
c     character*80  title, comment
c     character*12  buffer, buffer2

c     call date(buffer)
c     call time(buffer2)
c     print*, 'XSWISS_FS run on ', buffer, ' at ', buffer2

c     READ DATA
c     open(unit=3, file='parimp.dat', status='old')
c     read(unit=3, iat=90), title
c     read(unit=3, iat=93), fg1
c     read(unit=3, iat=93), fg2
c     read(unit=3, iat=93), fe1
c     read(unit=3, iat=93), fe2
c     read(unit=3, iat=93), fg3
c     read(unit=3, iat=90), comment
c     close(unit=3)

c     WRITE DATA
c     open(unit=3, file='fconstants.dat', status='new')
c     write(unit=3, iat=*)  title
c     write(unit=3, iat=93), 'FG1=', fg1
c     write(unit=3, iat=93), 'FG2=', fg2
c     write(unit=3, iat=93), 'FE1=', fe1
c     write(unit=3, iat=93), 'FE2=', fe2
c     write(unit=3, iat=93), 'FG3=', fg3
c     write(unit=3, iat=*)  comment
c     endfile(unit=3)
c     close(unit=3)

c     print*, 'PARIMP.DAT has been transferred to FCONSTANTS.DAT.'

c     -----
c     format (a)
c     format (' ', x, t20, e30.16)

```

## 6) REVISE\_FS

```

c -----
c -----
c   REVISE_FS
c   =====
c
c   Takes data from PARIMP.DAT
c   and writes a revised FCONSTANTS.DAT
c
c   Written:
c   1989/Jan/19   -by Philip Sharman.
c   Revised:
c   1989/Feb/20   -FG3 term added.
c -----
c -----
c
c   implicit real*8 (a-z)
c   character*80   title, comment
c   character*12   buffer, buffer2
c
c   call date(buffer)
c   call time(buffer2)
c   print*, 'REVISE_FS run on ', buffer, ' at ', buffer2
c
c   READ DATA
c   open(unit=3, file='parimp.dat', status='old')
c   read(unit=3, fmt=90), title
c   read(unit=3, fmt=93), fg1
c   read(unit=3, fmt=93), fg2
c   read(unit=3, fmt=93), fe1
c   read(unit=3, fmt=93), fe2
c   read(unit=3, fmt=93), fg3
c   read(unit=3, fmt=90), comment
c   close(unit=3)
c
c   WRITE DATA
c   open(unit=3, file='fconstants.dat', status='new')
c   write(unit=3, fmt=*), title
c   write(unit=3, fmt=92), 'FG1=', fg1
c   write(unit=3, fmt=92), 'FG2=', fg2
c   write(unit=3, fmt=92), 'FE1=', fe1
c   write(unit=3, fmt=92), 'FE2=', fe2
c   write(unit=3, fmt=92), 'FG3=', fg3
c   write(unit=3, fmt=*), comment
c   endfile(unit=3)
c   close(unit=3)
c
c   print*, 'PARIMP.DAT has been transferred to FCONSTANTS.DAT.'
c
c -----
c   format (a)
c   format (' ', a, t20, e30.16)

```

```

93      format (t20, e30.16)
c      -----
c      end
c
c      PAR_ALPHA
c      =====
c      This is "Parafit" modified to try to fit a parabola
c      to alpha versus I data.
c      y=I, x=alpha, y=a*x^2 + b*x +c
c
c      (See Basic Scientific Subroutines, Vol 2, page 80)
c
c      (Alter I_DESIRED to be the desired alpha.)
c
c      Written:
c      1983/May/30      -by Philip Sherman.
c                    -Modified from PARIMP.FOB.
c                    (Variable names etc cleaned up a lot.)
c
c      Revised:
c      1989/Jun/18     -Modified to handle variable number of data pairs.
c                    Jun/18     -quick plot added.
c                    Jun/17     -min/max added.
c                    Jun/18     -INT() bug fixed.
c                    Jun/18     -PLOT subroutine added.
c                    Jun/19     -initial guesses changed to put the parabola
c                    through the first two data points.
c
c      -----
c
c      GLOBAL VARIABLES
c      real*8          x(99), y(99), a(99), standard_deviation
c      integer         n_parameters, n_data_pairs, n_iterations
c
c      LOCAL VARIABLES
c      integer         i
c      real*8          xa, ya
c      character*100   array(99)
c      integer         ncols, nrow, row, col
c      real*8          b_col, b_row, a_col, a_row, xmin, xmax, ymin, ymax,
c                    x_desired, y_desired
c
c      GREETINGS
c      print*, 'Welcome to Par_Alpha'
c      print*, '-----'
c
c      DRIVER SECTION
c      print*, 'How many data pairs are there?'
c      read*, n_data_pairs
c      print*, 'n_data_pairs= ', n_data_pairs
c      do i= 1, n_data_pairs

```

## 7) PAR\_ALPHA

```

c -----
c -----
c   PAR_ALPHA
c   =====
c   This is "Parafit" modified to try to fit a parabola
c   to alpha versus A data.
c   y=A, x=alpha, y=a*x^2 + b*x +c
c
c   (See Basic Scientific Subroutines, Vol 2, page 80)
c
c   (Alter X_DESIRED to be the desired alpha.)
c
c   Written:
c   1989/May/30   -by Philip Sharman.
c                -Modified from PARIMP.FOR.
c                (Variable names etc cleaned up a lot.)
c
c   Revised:
c   1989/Jun/15   -Modified to handle variable number of data pairs.
c                Jun/16   -quick plot added.
c                Jun/17   -min/max added.
c                Jun/18   -INT() bug fixed.
c                Jun/18   -PLOT subroutine added.
c                Jun/19   -initial guesses changed to put the parabola
c                through the first two data points.
c -----
c -----
c
c   GLOBAL VARIABLES
c   real*8      x(99), y(99), a(99), standard_deviation
c   integer     n_parameters, n_data_pairs, n_iterations
c
c   LOCAL VARIABLES
c   integer     i (THESE PUT THE PARABOLA THROUGH THE FIRST TWO
c   real*8      xx, yy
c   character*100 array(99)
c   integer     ncols, nrows, row, col
c   real*8      b_col, b_row, m_col, m_row, xmin, xmax, ymin, ymax,
c   &           x_desired, y_desired
c
c   GREETINGS
c   print*, 'Welcome to Par_Alpha'
c   print*, '=====
c
c   DRIVER SECTION
c   print*, 'How many data pairs are there?'
c   read*, n_data_pairs
c   print*, 'n_data_pairs= ', n_data_pairs
c   do i= 1, n_data_pairs

```

```

      print 50, 'A(', i, ')?'
      read*, y(i)
      print*, y(i)
      print 50, 'alpha(', i, ')?'
      read*, x(i)
      print*, x(i)
    end do
50  format (' ', a, i2, a)

    n_parameters= 3

c    DATA
c    (STATE 2)
c    n_data_pairs= 3
c    x(1)= -3.27457
c    x(2)= -3.20235
c    x(3)= -3.21508
c    y(1)= 6.07
c    y(2)= 5.89
c    y(3)= 5.9
c    (STATE 3)
c    n_data_pairs= 2
c    x(1)= 4.83885
c    x(2)= 4.83511
c    y(1)= 12.71303
c    y(2)= 12.70

    print*, 'Input data:'
    print*, '      #      x      y'
    do i= 1, n_data_pairs
      print*, i, x(i), y(i)
    end do
    print*, ' '

c    STARTING GUESSES (THESE PUT THE PARABOLA THROUGH THE FIRST TWO
C    DATA POINTS)
    a(1)= (y(1)-y(2))/(x(1)**2 - x(2)**2)
    a(2)= 0.001      !(PARAFIT WON'T ACCEPT 0)
    a(3)= (y(1)*x(2)**2 - y(2)*x(1)**2)/(x(2)**2 - x(1)**2)

    print*, 'The initial guesses are...'
    print*, (a(i), i= 1,3)
    print*, 'With these parameters...'
    print*, 'x      Calculated',
&    print*, '      Desired '
    do i= 1, n_data_pairs
      xx= x(i)
      call func(xx, yy, a)
      print*, x(i), yy, y(i)
    end do

    call parafit(x, y, a, n_data_pairs, n_parameters,

```

```

& standard_deviation, n_iterations)

print*, ' '
print*, 'The coefficients found by PARAFIT are:'
print*, 'a= ', a(1)
print*, 'b= ', a(2)
print*, 'c= ', a(3)
print*, 'y= a*x^2 + b*x + c'
print*, ' '
print*, 'The standard deviation of the fit is:', standard_deviation
print*, 'The number of iterations was: ', n_iterations
print*, ' '
print*, 'With these parameters...'
print*, 'x          Calculated',
&          ,          Desired '

do i= 1, n_data_pairs
  xx= x(i)
  call func(xx, yy, a)
  print*, x(i), yy, y(i)
end do

print*, 'The vertex is:'
print*, 'x= ', -a(2)/(2*a(1))
print*, 'y= ', (4*a(1)*a(3) - a(2)*a(2))/(4*a(1))

c  WHAT WOULD GIVE DESIRED XX?
x_desired= +1.612885
ccc x_desired= -3.22577055
print*, 'To get alpha= ', x_desired
call func(x_desired, y_desired, a)
print*, '...choose A= ', y_desired

c  DO A QUICK PLOT
c  INITIALIZE
xmax= x_desired
xmin= x_desired
ymax= y_desired
ymin= y_desired
do i= 1, n_data_pairs-1
  xmax= max(x(i), x(i+1), xmax )
  ymax= max(y(i), y(i+1), ymax )
  xmin= min(x(i), x(i+1), xmin )
  ymin= min(y(i), y(i+1), ymin )
end do
print*, 'X goes from ', xmin, ' to ', xmax
print*, 'Y goes from ', ymin, ' to ', ymax
ncols= 78
nrows= 19
m_col= (ncols-1)/(xmax-xmin)
b_col= 1 - m_col*xmin
m_row= (nrows-1)/(ymax-ymin)
b_row= 1 - m_row*ymin

```

```

do row=1, n_rows
  array(row)= '.....'
&      // '.....'
end do

c PLOT PARABOLA
do col=1, n_cols
  xx= (col - b_col)/m_col
  call func(xx, yy, a)
  row= int(m_row*yy + b_row)
  if (col.ge.1 .and. col.le.n_cols .and. row.ge.1 .and. row.le.n_rows) then
    call plot('+', row, col, array)
  endif
end do

c PLOT X_DESIRED LINE
do row= 1, n_rows
  col= int(m_col*x_desired + b_col)
  if (col.ge.1 .and. col.le.n_cols .and. row.ge.1 .and. row.le.n_rows) then
    call plot('|', row, col, array)
  endif
end do

c PLOT O AT X,Y_DESIRED
col= int(m_col*x_desired + b_col)
row= int(m_row*y_desired + b_row)
call plot('O', row, col, array)

c PLOT DATA POINTS
do i=1, n_data_pairs
  col= int(m_col*x(i) + b_col)
  row= int(m_row*y(i) + b_row)
  if (col.ge.1 .and. col.le.n_cols .and. row.ge.1 .and. row.le.n_rows) then
    call plot(char(i+48), row, col, array)
  else
    print*, 'Data point out of range: ', x(i), y(i)
    print*, 'row= ', row, ' col= ', col
  endif
end do

c SHOW IT
do row= n_rows, 1, -1
  print*, array(row)(1:n_cols)
end do

end

c -----
c -----
c THE PARAFIT SUBROUTINE
subroutine parafit(x, y, a, n_data_pairs, n_parameters,
&      standard_deviation, n_iterations)

```

```

c      INPUT VARIABLES
real*8      x(*), y(*)
integer     n_data_pairs, n_parameters

c      IN/OUT VARIABLES
real*8      a(*)

c      OUTPUT VARIABLES
real*8      standard_deviation
integer     n_iterations

c      LOCAL VARIABLES
integer     i, mod_iterations
real*8      e1(99), max_variance, convergence_factor, a0, m0, m1,
&          l1, xx, yy, residual

c      SOME OF RUCKDESCHEL'S VARIABLE NAMES HAVE BEEN CHANGED:
C      E      -> MAX_VARIANCE
C      E1     -> CONVERGENCE_FACTOR
C      D      -> STANDARD DEVIATION
C      M      -> N_ITERATIONS
C      N      -> N_DATA_PAIRS
C      L      -> N_PARAMETERS
C      L2     -> RESIDUAL

c      MAX VARIANCE DETERMINES WHEN TO STOP
max_variance= 1.0d-6
c      CONVERGENCE_FACTOR SHOULD BE BETWEEN 0 AND 1.  (1 IS
C      FASTEST BUT IT MAY NOT BE STABLE)
convergence_factor= 0.1

c      LOCAL VARIABLES
do i= 1, n_parameters
    e1(i)= convergence_factor
end do

c      n_iterations= 0
l1= 1000000

c      SWEEP
4257 do i= 1, n_parameters
    a0= a(i)
    a(i)= a0
    (WHY IS THAT LINE THERE?)
c      MIDDLE
4261 call residuals(x, y, a, n_data_pairs, n_parameters,
&          residual, standard_deviation)
&          m0= residual
c      (AND)
c      a(i)= a0*(1-e1(i))
c      call residuals(x, y, a, n_data_pairs, n_parameters,
&          residual, standard_deviation)
&          m1= residual

c      if (m1.gt.m0) e1(i)= -e1(i)/2

```

```

        if (m1.lt.m0) e1(i)= 1.2*e1(i)
        if (m1.gt.m0) a(i)= a0
        if (m1.gt.m0) goto 4261          !(GOTO MIDDLE)
    end do

    n_iterations= n_iterations + 1
    mod_iterations= mod(n_iterations, 500)
    if (mod_iterations.eq.0) print*,'# of iterations= ', n_iterations
    if (n_iterations.gt.90000) call exit(0)

    if (residual.eq.0) return
    if (abs((l1-residual)/residual).lt.max_variance) return

    l1= residual
    goto 4257          !(GOTO SWEEP)

end
-----
c
c
c   RESIDUALS
c   subroutine residuals(x, y, a, n_data_pairs, n_parameters,
&   residual, standard_deviation)
c
c   INPUT VARIABLES
c   integer n_data_pairs, n_parameters
c   real*8   x(*), y(*), a(*)

c   OUTPUT VARIABLES
c   real*8   residual, standard_deviation

c   LOCAL VARIABLES
c   integer   j
c   real*8    xx, yy

c   COMPUTE
c   residual= 0
c   do j= 1, n_data_pairs
c       xx= x(j)
c       call func(xx, yy, a)
c       residual= residual+(y(j)-yy)*(y(j)-yy)
c   end do

c   if (n_data_pairs.gt.n_parameters) standard_deviation=
&   sqrt(residual/(n_data_pairs - n_parameters))
c   if (n_data_pairs.le.n_parameters) standard_deviation= sqrt(residual)
c   (ABOVE LINE NOT IN ORIGINAL PGM.)

c   return
c   end
-----

```



## 8) TWEAK

```

c -----
c -----
c TWEAK
c =====
c WELCOME TO TWEAK
c TWEAK is an itelligent parameter-tweaker. It assumes there
c are some variables ,x(1:n_x), and some functions of these variables,
c y(1:n_y)
c and some desired values, desired(1:n_y), of these functions.
c (No assumption is made as to whether the number of x's is
c equal to the number of functions.)
c It tries to tweak the variables to give the desired values, but
c instead of just doing each variable in turn
c (as PARAFIT does for example) it also assumes there is some other
c criterion for judging what tweakings are better than others.
c (This might be 'don't move the variables far from where they are
c now', or it might be 'keep the variables as small as possible', or
c 'keep the variables as close to certain values as possible.')

```

```

c      GLOBAL PARAMETERS
      debug= .true.
      n_x= 5
      n_y= 4

c      WELCOME TO TWEAK
      call date(buffer1)
      call time(buffer2)
      print*, '*****'
      print*, 'TWEAK running on ', buffer1, ' at ', buffer2
      print*, '*****'

c      READ F'S AND SET Y() ARRAY
      call read_stuff

c      SHOW STUFF
      print*, 'The values with the old F's are...'
      call evaluate_y(0) !(SHOW W'S)
      call find_actual_ys
      print*, 'y(1)=', y(1)
      print*, 'y(2)=', y(2)
      print*, 'y(3)=', y(3)
      print*, 'y(4)=', y(4)
      call show_fs

c      NOW SEARCH
      keep_searching= .true.
      n_loops= 1
      minimum_error= 9.9d30

c      MAIN LOOP
      do while (keep_searching .and. n_loops.le.maxloops)
         print*, ' '
         call get_matrix
         call show_table
         call evaluate_y(0)
         call change_lowest_penalty
c      ARE WE DONE?
         sum_error= 0.0d0
         do col=1, n_y
            sum_error= sum_error + abs(desired(col)-y(col))
         end do
         print*, 'Sum of errors is now= ', sum_error
         if (sum_error .lt. minimum_error) then
            minimum_error= sum_error
            best_trial= n_loops
         endif
         if (best_trial.gt.0) print*, '(Current best= ',
            &             minimum_error, 'at trial', best_trial, ')'
         print '(a, i3)', ' Trial: ', n_loops
         criterion= 1.0
         if (sum_error.lt.criterion) then
            print*, 'Errors in y's are acceptable. '

```

```

logical          print*, 'We are successful.'
integer          keep_searching= .false.
endif
n_loops= n_loops + 1
end do
666 print*, 'Program is done.'
end
-----
c
c
subroutine get_matrix
c WORK OUT MATRIX
c MATRIX(r,c) SHOWS HOW X(R) VARIES IF Y(C) IS MADE PERFECT
c MATRIX(r,0) HOLDS ORIGINAL UNCHANGED X'S

c GLOBAL VARIABLES
common /global/ debug, n_x, n_y, x, y, desired
logical          debug
integer          n_x, n_y
real*8          x(5), y(4), desired(4)

c OTHER COMMON VARIABLES
common /mat/     matrix
real*8          matrix(5,0:4)

c LOCAL VARIABLES
integer          row, col
real*8          x_old

do row= 1, n_x
  x_old= x(row)
  matrix(row, 0)= x_old
  do col= 1, n_y
    if (debug) print*, 'Finding matrix ', row, col
    call solve_x(row, col)
    matrix(row, col)= x(row)
  (NOW RESTORE ORIGINAL VALUE)
  x(row)= x_old
end do

return
end
-----
c
c
subroutine change_lowest_penalty
c FINDS ENTRY IN MATRIX WITH LOWEST PENALTY-FUNCTION
c CHANGES THAT X AND Y

c GLOBAL VARIABLES
common /global/ debug, n_x, n_y, x, y, desired

```

```

logical      debug
integer      n_x, n_y
real*8      x(5), y(4), desired(4)

c      OTHER COMMON VARIABLES
common /mat/ matrix
real*8      matrix(5,0:4)

c      LOCAL VARIABLES
integer      row, col, min_row, min_col, max_col, col1, col2
real*8      minimum, pen, zero, x_new, x_old, penalty(5,4),
&           maximum, size, b, d, e, f
parameter    (zero= 0.0d0)
save        col1, col2

c      BASE PENALTY ON ABSOLUTE MAGNITUDE OF THE VARIABLES.
do row= 1, n_x
  do col= 1, n_y
    x_new= matrix(row, col)
    size= abs(x_new)
    b= 1000.0
    d= 1.0
    (WOULD REALLY LIKE D ABOUT 1.0d-4)
    e= 100.0
    f= 10.0
  c      ***ABOVE VARIABLES NOT USED AT PRESENT***
  c      penalty(row,col)= size
  c      end do
end do

c      PRINT PENALTY MATRIX
ccc     print*, 'Penalty matrix:'
ccc     print*, '          (1)          ', col_last_time
ccc     & '          (2)          (3)          (4) '
ccc     do row= 1, n_x
ccc     c      print 150, row, (penalty(row, col), col=1, n_y)
ccc     end do
150     format (' ', 'x(', i1, ')', t21, 4(':', 'e14.3), ':')

c      FIND LOWEST PENALTY ROW IN COLUMN COL
c      minimum = 9.9d30
c      CHANGE THE COLUMN WHICH HAS THE least non-zero ERROR, BUT
c      AVOID LAST TWO COLUMNS DONE (COL1 AND COL2)
c      (INITIALIZE IF THIS IS THE FIRST TIME THROUGH)
if (col1.lt.1 .and. col1.gt.n_y) col1= 0
if (col2.lt.1 .and. col2.gt.n_y) col2= 0
205     print*, 'Will not do columns: ', col1, col2
minimum= +9.9d30
do col= 1, n_y
  pen= abs(desired(col) - y(col))
  if (pen.ne.zero .and.

```



```

c      print*, 'Minimum= ', minimum, ' so go back and try again.'
      col2= 0
      goto 205      !(BACK UP)
      else if (col1.ne.0 .and. col2.eq.0) then
c      WE MAY HAVE ALREADY BACKED UP ONCE
      print*, 'Minimum= ', minimum, ' so go back and try again.'
      col1= 0      !(NOW BOTH COL1 AND COL2 = 0)
      goto 205      !(BACK UP)
c      elseif (col1.eq.0 .and. col2.eq.0 .and. minimum.gt.1.0) then
      WE ARE IN TROUBLE NOW
c      print*, 'We have no choice.'
      else
      print*, '????', col1, col2
      call exit(0)
      endif
endif
print*, 'Minimum penalty= ', minimum
print '(a, i1, a, i1, a)',
&      ' Changing x(', min_row, ') to improve y(', col, ')'
call solve_x(min_row, col)
col2= col1
col1= col

return
end
-----
c
c      subroutine show_table
c      DISPLAY STUFF (AND SET Y() TO ACTUAL VALUES AT END)
c
c      GLOBAL VARIABLES
common /global/ debug, n_x, n_y, x, y, desired
logical      debug
integer      n_x, n_y
real*8      x(5), y(4), desired(4)
c
c      OTHER COMMON VARIABLES
common /mat/      matrix
real*8      matrix(5,0:4)
c
c      LOCAL VARIABLES
integer      row, col
c
c      PRINT MATRIX
print*, ' '
print*, '          no change      (1)      ',
&      '      (2)      (3)      (4) '
do row= 1, n_x
      print 103, row, (matrix(row, col), col=0, n_y)
end do
print 105, 'Desired y''s->', (desired(col), col= 1, n_y)

```

```

c      NOW SET Y() TO ACTUAL VALUES
      call find_actual_ys
      print 105, 'Actual y''s-> ', (y(col), col= 1, n_y)
      print 106, 'Difference-> ', ((desired(col)-y(col)), col= 1, n_y)
103    format (' ', 'x(', i1, ')', 1(':', f14.3), 4(':', e14.3), ':')
105    format (' ', a13, t21, 4(':', e14.3), ':')
106    format (' ', a13, t21, 4(':', f14.3), ':')

      call show_fs

      return
      end
-----
c
c      subroutine read_stuff
c      READ IN THE PUA'S ETC. ALSO WORK OUT DESIRED() ARRAY.
c      AND SET INITIAL x() ARRAY.

c      GLOBAL VARIABLES
      common /global/ debug, n_x, n_y, x, y, desired
      logical      debug
      integer      n_x, n_y
      real*8       x(5), y(4), desired(4)

c      COMMON VARIABLES
      common /con_1_uct/ pua_1_uct, pub_1_uct,
&          pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct
      common /con_2_uct/ pua_2_uct, pub_2_uct,
&          pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct
      common /con_3_uct/ pua_3_uct, pub_3_uct,
&          pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct
      common /con_1_dsb/ pua_1_dsb, pub_1_dsb,
&          pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb
      common /con_2_dsb/ pua_2_dsb, pub_2_dsb,
&          pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb
      common /con_3_dsb/ pua_3_dsb, pub_3_dsb,
&          pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb
      save   /con_1_uct/, /con_2_uct/, /con_3_uct/
      save   /con_1_dsb/, /con_2_dsb/, /con_3_dsb/
      real*8 pua_1_uct, pub_1_uct,
&          pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct,
&          pua_2_uct, pub_2_uct,
&          pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct,
&          pua_3_uct, pub_3_uct,
&          pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct,
&          pua_1_dsb, pub_1_dsb,
&          pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb,
&          pua_2_dsb, pub_2_dsb,
&          pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb,
&          pua_3_dsb, pub_3_dsb,
&          pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb

```

```

c      LOCAL VARIABLES
      character*80    fs_title
      character*100   character_junk, title
      real*8          junk, mass_1_uct, mass_2_uct, mass_3_uct,
&                   mass_1_dsb, mass_2_dsb, mass_3_dsb,
&                   fg1, fg2, fe1, fe2, fg3
      integer         i

c      READ F'S
      print*, '(Reading FCONSTANTS.DAT)'
      open(unit=3, file='fconstants.dat', status='old')
      read(unit=3, fmt=90), fs_title
      read(unit=3, fmt=93), fg1
      read(unit=3, fmt=93), fg2
      read(unit=3, fmt=93), fe1
      read(unit=3, fmt=93), fe2
      read(unit=3, fmt=93), fg3
      close(unit=3)

      print*, 'fg1=', fg1
      print*, 'fg2=', fg2
      print*, 'fe1=', fe1
      print*, 'fe2=', fe2
      print*, 'fg3=', fg3

901    format (' ', a1, ': ',
&          6(f7.3, 1x) )

c      READ UCT DATA
      print*, ' '
      open(unit=3, file='esc_uct.dat', status='old')
      read(unit=3, fmt=90), title
      print*, 'Reading ESC_UCT.DAT: ', title
      read(unit=3, fmt=93), junk
      read(unit=3, fmt=93), junk
      read(unit=3, fmt=93), junk
      read(unit=3, fmt=93), junk
      read(unit=3, fmt=93), junk

      read(unit=3, fmt=93), pua_1_uct
      read(unit=3, fmt=93), pub_1_uct
      read(unit=3, fmt=93), pta_1_uct
      read(unit=3, fmt=93), ptb_1_uct
      read(unit=3, fmt=93), pextra_1_uct
      read(unit=3, fmt=93), abg_1_uct
      read(unit=3, fmt=93), junk
      read(unit=3, fmt=93), junk

      read(unit=3, fmt=93), pua_2_uct
      read(unit=3, fmt=93), pub_2_uct

```

```

read(unit=3,fmt=93), pta_2_uct
read(unit=3,fmt=93), ptb_2_uct
read(unit=3,fmt=93), pextra_2_uct
read(unit=3,fmt=93), abg_2_uct
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

read(unit=3,fmt=93), pua_3_uct
read(unit=3,fmt=93), pub_3_uct
read(unit=3,fmt=93), pta_3_uct
read(unit=3,fmt=93), ptb_3_uct
read(unit=3,fmt=93), pextra_3_uct
read(unit=3,fmt=93), abg_3_uct
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

close(unit=3)

print*, 'UCT GROUP:'
print*, '  pua, pub, pta, ptb, pextra, abg  ...'
print 901, '1', pua_1_uct, pub_1_uct, pta_1_uct,
&          ptb_1_uct, pextra_1_uct, abg_1_uct
print 901, '2', pua_2_uct, pub_2_uct, pta_2_uct,
&          ptb_2_uct, pextra_2_uct, abg_2_uct
print 901, '3', pua_3_uct, pub_3_uct, pta_3_uct,
&          ptb_3_uct, pextra_3_uct, abg_3_uct

c READ DSB DATA
print*, ' '
open(unit=3, file='esc_dsb.dat', status='old')
read(unit=3, fmt=90), title
print*, 'Reading ESC_DSB.DAT: ', title
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

read(unit=3,fmt=93), pua_1_dsb
read(unit=3,fmt=93), pub_1_dsb
read(unit=3,fmt=93), pta_1_dsb
read(unit=3,fmt=93), ptb_1_dsb
read(unit=3,fmt=93), pextra_1_dsb
read(unit=3,fmt=93), abg_1_dsb
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

read(unit=3,fmt=93), pua_2_dsb
read(unit=3,fmt=93), pub_2_dsb
read(unit=3,fmt=93), pta_2_dsb
read(unit=3,fmt=93), ptb_2_dsb
read(unit=3,fmt=93), pextra_2_dsb

```

```

read(unit=3,fmt=93), abg_2_dsb
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

read(unit=3,fmt=93), pua_3_dsb
read(unit=3,fmt=93), pub_3_dsb
read(unit=3,fmt=93), pta_3_dsb
read(unit=3,fmt=93), ptb_3_dsb
read(unit=3,fmt=93), pextra_3_dsb
read(unit=3,fmt=93), abg_3_dsb
read(unit=3,fmt=93), junk
read(unit=3,fmt=93), junk

close(unit=3)

print*, 'DSB GROUP:'
print*, '   pua, pub, pta, ptb, pextra, abg   ...'
print 901, '1', pua_1_dsb, pub_1_dsb, pta_1_dsb,
&          ptb_1_dsb, pextra_1_dsb, abg_1_dsb
print 901, '2', pua_2_dsb, pub_2_dsb, pta_2_dsb,
&          ptb_2_dsb, pextra_2_dsb, abg_2_dsb
print 901, '3', pua_3_dsb, pub_3_dsb, pta_3_dsb,
&          ptb_3_dsb, pextra_3_dsb, abg_3_dsb

c -----
90 format (a)
93 format ( t20, e30.16)
c -----

c SET X() ARRAY
x(1)= fg1
x(2)= fg2
x(3)= fe1
x(4)= fe2
x(5)= fg3

c READ MASS DATA
print*, '(Reading UCT.DAT)'
open(unit=3, file='uct.dat', status='old')
  read(unit=3, fmt=90), title
  read(unit=3, fmt=93), mass_1_uct
  read(unit=3, fmt=93), mass_2_uct
  read(unit=3, fmt=93), mass_3_uct
close(unit=3)
print*, '(Reading DSB.DAT)'
open(unit=3, file='dsb.dat', status='old')
  read(unit=3, fmt=90), title
  read(unit=3, fmt=93), mass_1_dsb
  read(unit=3, fmt=93), mass_2_dsb
  read(unit=3, fmt=93), mass_3_dsb
close(unit=3)

```

```

c      WORK OUT DESIRED VALUES
      desired(1)= 0.0d0
      desired(2)= mass_2_uct/mass_1_uct
      desired(3)= mass_2_dsb/mass_1_dsb
      desired(4)= mass_3_dsb/mass_1_dsb

      print*, 'Desired values:'
      print*, 'desired(1)=', desired(1), '(IS UP POSITIVE?)'
      print*, 'desired(2)=', desired(2), '(CHARM/UP)'
      print*, 'desired(3)=', desired(3), '(STRANGE/DOWN)'
      print*, 'desired(4)=', desired(4), '(BOTTOM/DOWN)'

      print*, 'Variables:'
      do i=1, n_x
      print '(a, i1, a, e25.15)', ' x(', i, ')', x(i)
      end do

      return
      end
-----
c
c      subroutine solve_x(m, n)
c      ADJUSTS X(m) TO GIVE DESIRED Y(n)
c      (ALTERS Y(n) BUT RESTORES IT TO ORIGINAL CONDITION AT THE END)
c      IF IT FAILS IT SETS X(m)=1D6

c      GLOBAL VARIABLES
      common /global/ debug, n_x, n_y, x, y, desired
      logical      debug
      integer      n_x, n_y
      real*8       x(5), y(4), desired(4)

c      INPUT VARIABLES
      integer      m, n

c      LOCAL VARIABLES
      integer      n_moves, maxmoves
      parameter    (maxmoves=100)
      real*8       x_original, delta_x, y_original, error,
&                error_last_time, two, criterion
      parameter    (two=2.0d0, criterion=1.0d-4)
      logical      keep_moving

c      SET UP
      n_moves= 1
      keep_moving= .true.
      x_original= x(m)
      y_original= y(n)
      delta_x= x_original
      if (delta_x.eq.0.0d0) delta_x= 1.0d-4

```

```

c          x(n)= 1.0 d6
c          y(n)= y_original
c          x(m)= x_original + delta_x
c          call evaluate_y(n)
c          error_last_time= desired(n) - y(n)

c          do while(keep_moving .and. n_moves.le.maxmoves)
c              NOW MOVE AGAIN IN SAME DIRECTION
c              x(m)= x(m) + delta_x
c              call evaluate_y(n)
c              error= desired(n) - y(n)
c              DID IT GET BETTER OR WORSE?
c              if (abs(error).gt.abs(error_last_time)) then
c                  WORSE
c                  HALVE DELTA AND CHANGE DIRECTION
c                  delta_x= -(delta_x/two)
c                  if (debug) then
c                      print*, 'worse'
c                      print*, 'x(m)= ', x(m), ' y(n)= ', y(n),
c                      & 'error=', error
c                      print*, 'delta_x=', delta_x
c                  endif
c              else
c                  BETTER
c                  KEEP GOING
c                  delta_x= 1.2 * delta_x
c                  if (debug) then
c                      print*, 'better'
c                      print*, 'x(m)=', x(m), 'y(n)=', y(n),
c                      & 'error=', error
c                      print*, 'delta_x=', delta_x
c                  endif
c              endif
c              error_last_time= error
c              n_moves= n_moves + 1
c              TEST FOR ACCEPTABLE CONVERGENCE
c              if (abs(error).lt.criterion) keep_moving= .false.
c              ABORT IF WE ARE WANDERING OFF INTO INFINITY
c              if (abs(x(m)).gt.1.0d6) then
c                  keep_moving= .false.
c                  print*, 'SOLVE_X not converging. x(m)= ', x(m)
c                  print*, 'Number of iterations= ', n_moves
c                  y(n)= y_original
c                  return
c              endif
c              ABORT IF WE HAVE GONE ON TOO LONG
c              if (n_moves.gt.maxmoves) then
c                  keep_moving= .false.
c                  print*, 'SOLVE_X failed.'
c                  print*, 'Number of iterations= ', n_moves

```

```

c          x(m)= 1.0 d6
c          y(n)= y_original
c          (WE REALLY JUST WANT TO PUT A LARGE VALUE IN THE
c          PENALTY MATRIX ... DO IT DIRECTLY?)
c          print*, '(Setting x(m)=1.0d6)'
c          return
c      endif
c  end do

c  FINISH UP - WE DID OKAY
c  if (debug) print*, 'SOLVE_X is done. X(', m, ') is now: ', x(m)
c  if (debug) print*, 'Number of iterations= ', n_moves

c  (PUT ORIGINAL Y VALUE BACK INTO PLACE)
c  y(n)= y_original

c  return
c  end
c-----
c-----
c  subroutine evaluate_y(n)
c  FINDS Y(n) (AND CHANGES IT)
c  THIS IS WHERE THE ACTUAL EQUATIONS GO
c  (IF n=0 IT JUST SHOWS THE w's)

c  GLOBAL VARIABLES
c  common /global/ debug, n_x, n_y, x, y, desired
c  logical      debug
c  integer      n_x, n_y
c  real*8      x(5), y(4), desired(4)

c  OTHER COMMON VARIABLES (OBTAINED BY READ_STUFF)
c  common /con_1_uct/ pua_1_uct, pub_1_uct,
c  &                pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct
c  common /con_2_uct/ pua_2_uct, pub_2_uct,
c  &                pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct
c  common /con_3_uct/ pua_3_uct, pub_3_uct,
c  &                pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct
c  common /con_1_dsb/ pua_1_dsb, pub_1_dsb,
c  &                pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb
c  common /con_2_dsb/ pua_2_dsb, pub_2_dsb,
c  &                pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb
c  common /con_3_dsb/ pua_3_dsb, pub_3_dsb,
c  &                pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb
c  save  /con_1_uct/, /con_2_uct/, /con_3_uct/
c  save  /con_1_dsb/, /con_2_dsb/, /con_3_dsb/
c  real*8 pua_1_uct, pub_1_uct,
c  &      pta_1_uct, ptb_1_uct, abg_1_uct, pextra_1_uct,
c  &      pua_2_uct, pub_2_uct,
c  &      pta_2_uct, ptb_2_uct, abg_2_uct, pextra_2_uct,
c  &      pua_3_uct, pub_3_uct,
c  &      pta_3_uct, ptb_3_uct, abg_3_uct, pextra_3_uct,

```

```

&      if (n.eq.0) pua_1_dsb, pub_1_dsb,
&                pta_1_dsb, ptb_1_dsb, abg_1_dsb, pextra_1_dsb,
&                pua_2_dsb, pub_2_dsb,
&                pta_2_dsb, ptb_2_dsb, abg_2_dsb, pextra_2_dsb,
&                pua_3_dsb, pub_3_dsb,
&                pta_3_dsb, ptb_3_dsb, abg_3_dsb, pextra_3_dsb
c      INPUT VARIABLES
integer      n
c      LOCAL VARIABLES
real*8      zero, w1_uct, w2_uct, w3_uct,
&          w1_dsb, w2_dsb, w3_dsb,
&          fg1, fg2, fe1, fe2, fg3
parameter   (zero=0.0d0)
c      THE NITTY GRITTY
fg1= x(1)
fg2= x(2)
fe1= x(3)
fe2= x(4)
fg3= x(5)
if (fe1.ne.zero .and. fe2.ne.zero) then
&      w1_uct=      (fg1/fe1)*pua_1_uct + (fg2/fe2)*pub_1_uct
&                + fg1*pta_1_uct + fg2*ptb_1_uct - abg_1_uct
&                + (fg3/fe1)*pextra_1_uct
&      w2_uct=      (fg1/fe1)*pua_2_uct + (fg2/fe2)*pub_2_uct
&                + fg1*pta_2_uct + fg2*ptb_2_uct - abg_2_uct
&                + (fg3/fe1)*pextra_2_uct
&      w3_uct=      (fg1/fe1)*pua_3_uct + (fg2/fe2)*pub_3_uct
&                + fg1*pta_3_uct + fg2*ptb_3_uct - abg_3_uct
&                + (fg3/fe1)*pextra_3_uct
&      w1_dsb=      (fg1/fe1)*pua_1_dsb + (fg2/fe2)*pub_1_dsb
&                + fg1*pta_1_dsb + fg2*ptb_1_dsb - abg_1_dsb
&                + (fg3/fe1)*pextra_1_dsb
&      w2_dsb=      (fg1/fe1)*pua_2_dsb + (fg2/fe2)*pub_2_dsb
&                + fg1*pta_2_dsb + fg2*ptb_2_dsb - abg_2_dsb
&                + (fg3/fe1)*pextra_2_dsb
&      w3_dsb=      (fg1/fe1)*pua_3_dsb + (fg2/fe2)*pub_3_dsb
&                + fg1*pta_3_dsb + fg2*ptb_3_dsb - abg_3_dsb
&                + (fg3/fe1)*pextra_3_dsb
else
&      print*, 'PROBLEM in EVALUATE. (Division by zero) FE1= ',
&            fe1, ' FE2= ', fe2
c      THEN CARRY ON? STOP? RETURN? FORCE THEM NON-ZERO?
do call exit(0)
endif
c      JUST SHOW W'S IF N=0

```

```

-----
if (n.eq.0)then
-----
    print*, 'w1_uct= ', w1_uct
    print*, 'w2_uct= ', w2_uct
    print*, 'w3_uct= ', w3_uct
    print*, 'w1_dsb= ', w1_dsb
    print*, 'w2_dsb= ', w2_dsb
    print*, 'w3_dsb= ', w3_dsb
c OTHERWISE UPDATE Y(n)
else if (n.eq.1) then
c   IS UP MASS POSITIVE?
    if (w1_uct.lt.zero) then
        y(1)= abs(w1_uct)
    else if (w1_uct.eq.zero) then
        y(1)= 9.9d6
    else
        y(1)= zero
    endif
else if (n.eq.2) then
    y(2)= w2_uct/w1_uct      !(CHARM/UP)
else if (n.eq.3) then
    y(3)= w2_dsb/w1_dsb    !(STRANGE/DOWN)
else if (n.eq.4) then
    y(4)= w3_dsb/w1_dsb    !(BOTTOM/DOWN)
else
    print*, 'ERROR: n= ', n, 'n_x= ', n_x
    call exit(0)
endif

return
end

-----
-----
c
c   subroutine find_actual_ys
c   UPDATE ALL THE Y() ARRAY

c
c   GLOBAL VARIABLES
c   common /global/ debug, n_x, n_y, x, y, desired
c   logical      debug
c   integer      n_x, n_y
c   real*8      x(5), y(4), desired(4)

c
c   LOCAL VARIABLE
c   integer col

c
c   FIND THEM
do col= 1, n_y
    call evaluate_y(col)
end do

return
end

```

```

-----
c
c
      subroutine show_fs
c      SHOW THE FG'S AND FE'S ETC.

c      GLOBAL VARIABLES
      common /global/ debug, n_x, n_y, x, y, desired
      logical      debug
      integer      n_x, n_y
      real*8       x(5), y(4), desired(4)

c      LOCAL VARIABLES
      real*8       fg1, fg2, fe1, fe2, fg3, eq(6)
      integer      i

c      WORK THEM OUT
      fg1= x(1)
      fg2= x(2)
      fe1= x(3)
      fe2= x(4)
      fg3= x(5)

c      ALSO WORK OUT THE QUANTITIES APPEARING IN THE FIELD EQUATIONS
      eq(1)= fg1
      eq(2)= fg2
      eq(3)= fe1
      eq(4)= fg3/fg1
      eq(5)= fe2
      eq(6)= fg3*fe2/(fg2*fe1)

c      PRINT THEM
      print*, '      fg1      fg2      fe1      fe2      fg3'
      print 755, fg1, fg2, fe1, fe2, fg3
      print 855, fg1, fg2, fe1, fe2, fg3

ccc     print*,
ccc     & '      fg1      fg2',
ccc     & '      fe1      (fg3/fg1)      fe2 (fg3*fe2)/(fg2*fe1)'
ccc     print 760, (eq(i), i=1, 6)
ccc     print 860, (eq(i), i=1, 6)

755     format (' ', 5(f12.7) )
855     format (' ', 5(e12.4) )
760     format (' ', 6(f12.7) )
860     format (' ', 6(e12.4) )

      return
      end
-----
c
-----

```

## VITA

Surname: Sharman. Given names: Philip Howard.

Place of birth: London, England.

Date of birth: 1956/October/18.

### Educational institutions attended:

Ryerson Polytechnical Institute, Toronto, 1975 - 1978,

University of Calgary, Calgary, 1982 - 1986,

University of Victoria, Victoria, 1987 - 1989.

### Degrees awarded:

Bachelor of Applied Arts, Ryerson Polytechnical Institute, 1978,

Bachelor of Science (Honours), University of Calgary, 1986.

### Honors:

All-Canada Award (Ryerson),

Andrew Kirkor Memorial Bursary (Calgary),

NSERC Undergraduate Student Research Award (Calgary).

### Publications:

"A Non-Linear Gauge-Invariant Field Theory Of Elementary Particles", P. Sharman and F. I. Cooperstock, in press.

*Cosmic Ray Intensity Records*, D. Venkatesan, T. Mathews, H. Grauman, and P. Sharman, (University of Calgary: 1989).

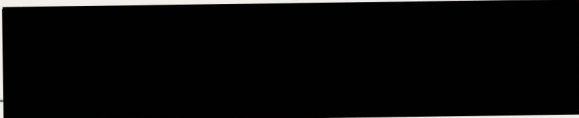
"The Trifair Cipher: The Trifair extends the Playfair into a ternary cipher system," *The Cryptogram*, March-April 1983.

"Trifair Variations and Extensions," *The Cryptogram*, May-June 1983.

## PARTIAL COPYRIGHT LICENSE

I hereby grant the right to lend my thesis entitled "Numerical Solutions of a Non-Linear Classical Field Theory of Leptons and Quarks" to users of the University of Victoria Library, and to make single copies only for such users or in response to a request from the library of any other university, or similar institution, on its behalf or for one of its users. I further agree that permission for extensive copying of this thesis for scholarly purposes may be granted by me or a member of the University designated by me. It is understood that copying or publication of this thesis for financial gain shall not be allowed without my written permission.

Author:

  
Philip Howard Sharman.

Date:

1990/Feb./6